

Dispersion Relations

■ Linear Response

■ Definition

The **response** r of a system subject to an external time dependent influence f is defined to be:

$$r(t) = \int_{-\infty}^{\infty} dt' g(t, t') f(t')$$

When g is independent of f , the response is called **linear**.

■ Impulsive Forces

Consider an influence of the **impulse** type:

$$f_I(t'; t_0) = f_0 \delta(t' - t_0) \quad \rightarrow \quad r_I(t; t_0) = f_0 g_I(t, t_0)$$

If the system is in equilibrium before the application of f , the response should have the form

$$\begin{aligned} r_I(t; t_0) &= \theta(t - t_0) r_I(t - t_0) \\ \rightarrow \quad g_I(t, t_0) &= \theta(t - t_0) g(t - t_0) \\ g(t - t_0) &= r_I(t - t_0) / f_0 & r_I(t - t_0) &= g(t - t_0) f_0 \end{aligned}$$

■ General Case

Using $f(t') = \int_{-\infty}^{\infty} dt'' f(t'') \delta(t' - t'')$, we have

$$\begin{aligned} r(t) &= \int_{-\infty}^{\infty} dt' g(t, t') \int_{-\infty}^{\infty} dt'' f(t'') \delta(t' - t'') \\ &= \int_{-\infty}^{\infty} dt'' \int_{-\infty}^{\infty} dt' g(t, t') f(t'') \delta(t' - t'') \\ &= \int_{-\infty}^{\infty} dt'' \int_{-\infty}^{\infty} dt' g(t, t') f_I(t''; t') \\ &= \int_{-\infty}^{\infty} dt'' r_I(t; t'') \\ &= \int_{-\infty}^{\infty} dt'' \theta(t - t'') r_I(t - t'') \\ &= \int_{-\infty}^{\infty} dt'' \theta(t - t'') g(t - t'') f(t'') \end{aligned}$$

$$\rightarrow \quad g(t, t') = \theta(t - t') g(t - t') \quad (\text{causality})$$

$$\begin{aligned} \therefore r(t) &= \int_{-\infty}^{\infty} dt' \theta(t - t') g(t - t') f(t') = \int_{-\infty}^t dt' g(t - t') f(t') \\ &= \int_{\tau=t-t'}^{\infty} d\tau g(\tau) f(t - \tau) \end{aligned}$$

■ Dispersion

■ Fourier Transform $G(\omega)$

Let the fourier transform of a function $h(t)$ be denoted by the capitalized symbol $H(\omega)$:

$$h(t) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} H(\omega)$$

$$H(\omega) = \int_{-\infty}^{\infty} dt e^{i\omega t} h(t)$$

Convolution Theorem applied to the linear response

$$\rightarrow R(\omega) = G(\omega) F(\omega)$$

$$\text{where } G(\omega) = \int_{-\infty}^{\infty} d\tau e^{i\omega\tau} \theta(\tau) g(\tau) = \int_0^{\infty} d\tau e^{i\omega\tau} g(\tau)$$

$$g(\tau) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega\tau} G(\omega)$$

$$\text{Now: } \int_0^{\infty} d\tau e^{i\omega\tau} g(\tau) \leq \int_0^{\infty} d\tau \left| e^{i\omega\tau} g(\tau) \right| = \int_0^{\infty} d\tau \left| g(\tau) \right|$$

$\therefore G(\omega)$ exists

$$\Rightarrow \int_0^{\infty} d\tau \left| g(\tau) \right| \text{ exists}$$

$$\Rightarrow \left| g(\tau) \right| \xrightarrow{\tau \rightarrow \infty} 0$$

■ Analytic Continuation $G(z)$

We now analytic continue $G(\omega)$ into the entire complex plane $z = \omega + i\omega'$.

$$G(z) = \int_0^{\infty} d\tau e^{iz\tau} g(\tau) = \int_0^{\infty} d\tau e^{i\omega\tau - \omega'\tau} g(\tau)$$

$$\left| G(z) \right| \xrightarrow{|z| \rightarrow \infty} 0 \quad \text{in the upper half plane.}$$

Using

$$\left| e^{iz\tau} \right|^2 = e^{iz\tau - iz^*\tau} = e^{-2\tau \text{Im}z} = e^{-2\tau |z| \sin\theta} \quad \text{where } \theta = \arg z$$

we have:

$$\left| G(z) \right| \leq \int_0^{\infty} d\tau \left| e^{iz\tau} \right| \left| g(\tau) \right|$$

$$= \int_0^{\infty} d\tau e^{-\tau |z| \sin\theta} \left| g(\tau) \right|$$

$$\leq M_G \int_0^{\infty} d\tau e^{-\tau |z| \sin \theta}$$

$$= \frac{M_G}{|z| \sin \theta} \quad \text{if } 0 < \theta < \pi \text{ (upper half } z \text{- plane)}$$

where $M_G = \max |g|$

Hence for $\theta \neq 0, \pi$,

$$\left| G(z) \right| \xrightarrow{|z| \rightarrow \infty} 0$$

For $\theta = 0$ or π , $G(z) = G(\omega)$.

Similar arguments applied to the fact that

$$g(\tau) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega\tau} G(\omega) \text{ exists}$$

gives: $\left| G(\omega) \right| \xrightarrow{|\omega| \rightarrow \infty} 0$

Hence $\left| G(z) \right| \xrightarrow{|z| \rightarrow \infty} 0$ in the upper half plane.

$G(z)$ is analytic in the half plane above the real axis