# **Further Applications**

# Summation of Series

$$\sum_{n=-\infty}^{\infty} f(n) = -\pi \sum_{n} \text{Res} \{ f(z_n) \cot(\pi z_n) \}$$

$$\sum_{n=-\infty}^{\infty} (-)^n f(n) = -\pi \sum_{n} \text{Res} \{ f(z_n) \csc(\pi z_n) \}$$

where  $z_n$  are the poles of f(z).

### proof

The major step in the proof is to show that

$$g(z) = \pi \cot(\pi z)$$

has simple poles at z = n with residue 1.

Given the above, the 1st formula is obtained by considering the contour integral which encloses the entire z-plane.

Let 
$$g(z) = \pi \cot(\pi z) = \pi \frac{\cos(\pi z)}{\sin(\pi z)}$$

 $\longrightarrow$  poles of g(z) are simple & at z = n.

Res 
$$g(z) = \left(\pi \frac{\cos(\pi z)}{\frac{d}{dz}\sin(\pi z)}\right)_{z=n} = 1$$

To evaluate

$$S_N = \sum_{n=-N}^{N} f(n)$$

we use a contour  $C_N$  which is a square that is centered at z = 0 & intersects the real axis at  $x = \pm \left(N + \frac{1}{2}\right)$ :

$$\longrightarrow \oint_{c_N} dz \, g(z) \, f(z) = 2 \, \pi \, i \left\{ \sum_{n=-N}^N f(N) + \sum_k \text{Res} \left[ f(z_k) \, \pi \cot(\pi \, z_k) \right] \right\}$$

where k runs over all poles of f inside  $C_N$ .

As 
$$N \longrightarrow \infty$$
,  $\oint_{c_N} dz g(z) f(z) \longrightarrow 0$  since  $|f| \underset{|z| \to \infty}{\longrightarrow} 0$  if  $S$  converges. QED.

Proof for the 2nd formula is analogous. All we need is

Res 
$$\pi \csc(\pi z) = \left(\pi \frac{1}{\frac{d}{dz}\sin(\pi z)}\right)_{z=n} = \frac{1}{\cos(\pi n)} = (-)^n$$

# ■ Example

$$S = \sum_{n=-\infty}^{\infty} \frac{(-)^n}{(a+n)^2}$$

$$f(z) = \frac{1}{(a+z)^2}$$
 with 2nd order pole at  $z = -a$ 

$$S = -\pi \operatorname{Res}_{z=-a} \left\{ \frac{1}{(a+z)^2} \cot(\pi z) \right\}$$

$$= -\pi \left\{ \frac{d}{dz} \csc(\pi z) \right\}_{z=-a}$$

$$= \pi^2 \frac{\cos(\pi a)}{\sin^2(\pi a)}$$

#### Definition

Let

$$S_N = \sum_{n=0}^N \frac{A_n}{z^n} \qquad S = \sum_{n=0}^\infty \frac{A_n}{z^n}$$

If

$$\begin{split} f(z) &= \phi(z) \, S \\ \lim_{|z| \to \infty} \, \left\{ \, z^N \left[ \, \frac{f(z)}{\phi(z)} - \, S_N \right] \, \right\} &= 0 \end{split}$$

 $\implies$  S represents  $\frac{f(z)}{\phi(z)}$  asymptotically.

#### Note:

Usually, S diverges & there is an optimal N which gives the best approximation.

#### ■ Example

# **Exponential Function:**

$$Ei(x) = \int_{-\infty}^{x} dt \frac{e^{-t}}{t}$$

$$E_{1}(x) = \int_{x}^{\infty} dt \frac{e^{-t}}{t} = -Ei(-x)$$

$$= -\int_{x}^{\infty} de^{-t} \cdot \frac{1}{t} = -\left(\frac{e^{-t}}{t}\right)_{x}^{\infty} - \int_{x}^{\infty} dt \frac{e^{-t}}{t^{2}}$$

$$= \frac{e^{-x}}{x} - \int_{x}^{\infty} dt \frac{e^{-t}}{t^{2}}$$

$$= \frac{e^{-x}}{x} - \frac{e^{-x}}{x^{2}} + 2 \int_{x}^{\infty} dt \frac{e^{-t}}{t^{3}}$$

$$= \frac{e^{-x}}{x} - \frac{e^{-x}}{x^{2}} + 2 \frac{e^{-x}}{x^{3}} - 3! \int_{x}^{\infty} dt \frac{e^{-t}}{t^{4}}$$

$$= \frac{e^{-x}}{x} \left\{ 1 - \frac{1}{x} + \frac{2!}{x^{2}} - \frac{3!}{x^{3}} + \dots + (-)^{n} \frac{n!}{x^{n}} \right\} + (-)^{n+1} (n+1)! \int_{x}^{\infty} dt \frac{e^{-t}}{t^{n+2}}$$

$$= \frac{e^{-x}}{x} S_{n} + (-)^{n+1} (n+1)! \int_{x}^{\infty} dt \frac{e^{-t}}{t^{n+2}}$$
where  $S_{n} = 1 - \frac{1}{x} + \frac{2!}{x^{2}} - \frac{3!}{x^{3}} + \dots + (-)^{n} \frac{n!}{x^{n}} = \sum_{m=0}^{n} (-)^{m} \frac{m!}{x^{m}}$ 

Cauchy Test:

$$\lim_{m \to \infty} \ \left| \frac{\frac{(m+1)!}{x^{m+1}}}{\frac{m!}{x^m}} \right| = \lim_{m \to \infty} \frac{(m+1)}{|x|} \longrightarrow \infty$$

$$\therefore S = \lim_{n \to \infty} S_n \text{ diverges.}$$

On the other hand:

the other hand:  

$$\lim_{x \to \infty} (x^n [E_1(x) x e^x - S_n])$$

$$= \lim_{x \to \infty} \left( x^{n+1} e^x (-)^{n+1} (n+1)! \int_x^{\infty} dt \frac{e^{-t}}{t^{n+2}} \right)$$

$$< \lim_{x \to \infty} \left( x^{-1} e^x (-)^{n+1} (n+1)! \int_x^{\infty} dt e^{-t} \right)$$

$$= \lim_{x \to \infty} \left\{ x^{-1} (-)^{n+1} (n+1)! \right\}$$

$$= 0$$

 $E_1(x) = \frac{e^{-x}}{x} S$  is an asymptotic representation.

The error involved in using the  $S_n$  is

$$(-)^{n+1} (n+1)! \int_{x}^{\infty} dt \frac{e^{-t}}{t^{n+2}}$$

$$< (-)^{n+1} (n+1)! \frac{1}{x^{n+2}} \int_{x}^{\infty} dt e^{-t}$$

$$= (-)^{n+1} (n+1)! \frac{1}{x^{n+2}} e^{-x}$$

# Properties

1. Function represented by asymptotic expansion is not unique.

eg. 
$$\frac{f}{\phi}$$
 &  $\frac{f}{\phi} + e^{-z}$  have the same expansion.

- 2. Phase change in z often produces discontinuities ( **Stokes Phenomena** ).
- Asymptotic series can be added, multiplied & integrated. 3.
- 4. Differentiation of asymptotic series is valid only if the derivative function also has an asymptotic series expansion.

# **Method of Steepest Descent** (Saddle Point Method)

1st Term

### ■ Formula

$$\int_{c} dt g(t) e^{zf(t)} = \sum_{s} g(t_{s}) e^{zf(t_{s})} e^{i\tau_{s}} \sqrt{\frac{2\pi}{|zf''(t_{s})|}}$$
$$= \sum_{s} (\pm) g(t_{s}) e^{zf(t_{s})} \sqrt{\frac{2\pi}{zf''(t_{s})}}$$

where |z| >> 1,

 $g(t) \approx \text{const near the saddle points } t_s$ ,

$$e^{z f(t)} = 0$$
 at the end points of  $c$ .

& 
$$2\tau_s = \pm \pi - \operatorname{Arg} f''(t_s) - \operatorname{Arg} z$$

with the sign chosen to conform with the original contour.

We shall begin with the integral

$$I(z) = \int_{C} dt \, e^{z f(t)}$$

Consider a saddle points  $t_s$  of f

ie. 
$$f'(t_s) = 0$$
.

Let 
$$f(t) = f(t_s) + \frac{1}{2} f''(t_s) (t - t_s)^2 + \dots$$
$$f''(t_s) = \left| f''(t_s) \right| e^{i\phi}$$
$$t - t_s = T e^{i\tau}$$
$$z = \left| z \right| e^{i\zeta}$$

Keeping only terms up to the 2nd order in *T*:

$$z f(t) = z f(t_s) + \frac{1}{2} |z f''(t_s)| e^{i(\zeta + \phi + 2\tau)} T^2$$
  
=  $z f(t_s) + \frac{1}{2} |z f''(t_s)| T^2 \{\cos(\zeta + \phi + 2\tau) + i\sin(\zeta + \phi + 2\tau)\}$ 

Now, c or  $\tau$  is chosen such that along c,

1. Re  $\{z[f(t) - f(t_s)]\}$  is a maximum

2. 
$$\operatorname{Im}\left\{z\left[f(t)-f(t_{s})\right]\right\} = \operatorname{const}\operatorname{near}t_{s}.$$

This can be accomplished by setting

$$\zeta + \phi + 2\tau = \pm \pi$$

so that

$$z\{f(t) - f(t_s)\} = \text{Re}\{z[f(t) - f(t_s)]\}\$$

$$= -\frac{1}{2}|zf''(t_s)| T^2$$

$$< 0$$

on c near  $t_s$ :

$$dt = e^{i\tau} dT$$
  $(\tau = const)$ 

Thus, near  $t_s$ , the contibution to I(z) is:

$$I_s(z) = e^{z f(t_s)} e^{i \tau} \int_{C_s} dT e^{-\frac{1}{2} |z f''(t_s)|} T^2$$

where  $c_s$  is the portion ( $\tau = \text{const}$ ) of the deformed c which goes through  $t_s$ .

Since |z| >> 1, we can replace  $\int_{C_s} dT$  by  $2 \int_0^{\infty} dT$ .

$$\longrightarrow \int_{c_s} dT e^{-\frac{1}{2} \left| z f''(t_0) \right| T^2} \simeq \sqrt{\frac{2\pi}{\left| z f''(t_s) \right|}}$$

$$I_s(z) = e^{z f(t_s)} e^{i\tau} \sqrt{\frac{2\pi}{\left| z f''(t_s) \right|}}$$

where

$$2\tau = \pm \pi - \phi - \zeta$$
  
= \pm \pi - \text{Arg } f \cdot' \((t\_s) - \text{Arg } z\)

with the sign chosen to conform with the original contour.

Away from the saddle points,  $\text{Im} \{z[f(t) - f(t_s)]\}$  varies rapidly so that contributions from different parts of the contour cancels each other. It is therefore a good approximation to write:

$$I(z) = \sum_{s} I_{s}(z)$$

$$= \sum_{s} e^{z f(t_{s})} e^{i \tau_{s}} \sqrt{\frac{2\pi}{|z f''(t_{s})|}}$$

$$2 \tau_{s} = \pm \pi - \phi_{s} - \zeta$$

where

Substituting  $\tau$  into the last expression gives

$$I_s(z) = \pm e^{z f(t_s)} \sqrt{\frac{-2\pi}{z f''(t_s)}}$$

with the sign is again chosen to conform with the original contour.

Actually, this result can be directly arrived at following the previous derivation but without introducing all the phase angles.

It's easy to see that

$$\int_{c} dt g(t) e^{z f(t)} = \sum_{s} g(t_{s}) e^{z f(t_{s})} e^{i \tau_{s}} \sqrt{\frac{2\pi}{|z f''(t_{s})|}}$$

$$= \sum_{s} (\pm) g(t_{s}) e^{z f(t_{s})} \sqrt{\frac{2\pi}{z f''(t_{s})}}$$

provided  $g \simeq \text{const}$  near each saddle point.

More precisely, we require  $|\Delta g| \ll |\Delta e^{zf}|$  near each  $t_s$ .

This is satisfied for any rational function *g*.

#### $H_{..}^{(1)}(s)$ Example

$$H_{\nu}^{(1)}(s) = \frac{1}{\pi i} \int_{c} dz \, \frac{e^{\frac{s}{2} \left(z - \frac{1}{c}\right)}}{z^{\nu+1}}$$
 (s real & >> 1)

where c goes from z = 0 with initial zero slope via the upper plane to  $z = (-\infty, 0)$  asymptotically along the negative x-axis. ( See Arfken, 3rd ed., eg 7.4.1.)

Let 
$$f(z) = \frac{1}{2} \left( z - \frac{1}{z} \right)$$
 
$$g(z) = \frac{1}{\pi i z^{\nu+1}}$$

$$f' = \frac{1}{2} \left( 1 + \frac{1}{z^2} \right)$$

The saddle point in the upper plane is  $z_0 = i$ .

$$f(z_0) = i g(z_0) = -\frac{1}{\pi i'}$$

$$f'' = -\frac{1}{z^3} f''(z_0) = -i = e^{i\frac{3\pi}{2}}$$

$$\to \phi = \frac{3\pi}{2} \tau = \pm \frac{\pi}{2} - \frac{3\pi}{4} = -\frac{\pi}{4} \text{ or } -\frac{5\pi}{4}$$

To conform with the original c, we must have  $\tau = -\frac{5\pi}{4}$ 

If we are to use the other formula, all the phase calculations can be omitted & we arrive directly at:

$$H_{\nu}^{(1)}(s) = -\left(-\frac{1}{\pi i^{\nu}}\right)e^{is} \sqrt{\frac{2\pi}{si}}$$

where the negative overall sign is used because the upper limit of the original contour is  $-\infty$ .

#### Example

$$\Gamma(z+1) = z! = \int_{0}^{\infty} dx \ x^{z} e^{-x}$$
 (|z|>>1)

The original contour is along the positive real axis.

To put the integral into a form to which the saddle point method is applicable, let

$$x = zt$$

$$dx = z dt x^{z} = z^{z} t^{z} t \in [0, \infty)$$

$$z! = z^{z+1} \int_{0}^{\infty} dt \ t^{z} e^{-zt}$$

$$= z^{z+1} \int_{0}^{\infty} dt \ e^{z(\ln t - t)}$$

Let 
$$f(t) = \ln t - t$$

$$\longrightarrow f' = \frac{1}{t} - 1$$

$$\therefore$$
 saddle point at  $t_0 = 1$ 

saudie point at 
$$t_0 = 1$$
  

$$f(1) = -1$$

$$f'' = -\frac{1}{t^2} \longrightarrow f''(1) = -1 = e^{i\pi} \qquad (\phi = \pi)$$

$$z = |z| e^{i\zeta}$$

$$2\tau = \pm \pi - \pi - \zeta = -\zeta \text{ or } -2\pi - \zeta$$

$$\tau = -\frac{\zeta}{2} \text{ or } -\pi - \frac{\zeta}{2}$$

We must use  $\tau = -\frac{\zeta}{2}$  to conform with the original contour. (This is obvious for the case of z real, ie.  $\zeta = 0$ ).

$$z != z^{z+1} e^{-z} e^{-i\frac{\zeta}{2}} \sqrt{\frac{2\pi}{|z|}}$$

$$= z^{z+1} e^{-z} \sqrt{\frac{2\pi}{z}}$$

$$= z^{z+\frac{1}{2}} e^{-z} \sqrt{2\pi}$$

If we are to use the other formula, all the phase calculations can be omitted & we arrive directly at:

$$z!=z^{z+1}e^{-z}\sqrt{\frac{2\pi}{z}}$$

#### Asymptotic Series

# ■ Theory

Keeping the higher order terms in the series

$$f(t) = f(t_s) + \frac{1}{2} f''(t_s) (t - t_s)^2 + \dots$$

gives us an asymptotic series of the integral.

For simplicity, we'll treat only the case

$$I(z) = \int_{C} dt \, e^{z f(t)}$$

with a single saddle point. Generalization to more complicated situations should be straightforward.

To begin, let

$$f(t) = f(t_s) - w^2$$

where w is real.

$$\longrightarrow I(z) = e^{z f(t_s)} \int_C dt e^{-z w^2}$$
$$= e^{z f(t_s)} \int_C dw \frac{dt}{dw} e^{-z w^2}$$

Expanding  $\frac{dt}{dw}$  as a power series:

$$\frac{dt}{dw} = \sum_{n=0}^{\infty} a_n w^n$$

where only even power need be retained.

$$\longrightarrow I(z) = \sum_{n=0}^{\infty} a_n e^{z f(t_s)} \int_{c} dw \ w^n e^{-z w^2} 
\simeq \pm \sum_{n=0}^{\infty} a_n e^{z f(t_s)} \int_{-\infty}^{\infty} dw \ w^n e^{-z w^2} 
= \pm \sum_{n=0}^{\infty} a_{2n} e^{z f(t_s)} \Gamma(n + \frac{1}{2}) \frac{1}{a^{n+\frac{1}{2}}}$$

with the sign chosen to conform with the original contour.

The main job is to find  $a_{2n}$ .

By integration from  $t_s$  to t, we see that

$$t - t_s = \sum_{n=0}^{\infty} a_n \frac{w^{n+1}}{n+1} \qquad (w(t_s) = 0)$$
$$= \sum_{n=1}^{\infty} a_{n-1} \frac{w^n}{n}$$

On the other hand,

$$w^2 = f(t) - f(t_s) = \sum_{n=2}^{\infty} A_n (t - t_s)^n$$

where

$$A_n = -\frac{1}{n!} f^{(n)}(t_s)$$

The calculation of  $a_{2n}$  is then just the inversion of the power series of  $w^2$ .

There are many ways to do so, one of which is by contour integration. ( see Morse & Feshbach ).

A more elementary way is to substitute the series of  $t - t_s$  into that of  $w^2$  & collect coefficients. Thus, putting  $T = t - t_s$ , we have

$$w^{2} = \sum_{n=2}^{\infty} A_{n} \left\{ \sum_{m=1}^{\infty} a_{m-1} \frac{w^{m}}{m} \right\}^{n}$$

$$= \sum_{n=2}^{\infty} A_{n} \left\{ a_{0} w + \frac{a_{1}}{2} w^{2} + \frac{a_{2}}{3} w^{3} + \dots \right\}^{n}$$

$$= \sum_{n=2}^{\infty} A_{n} a_{0}^{n} w^{n} \left\{ 1 + \frac{a_{1}}{2a_{0}} w + \frac{a_{2}}{3a_{0}} w^{2} + \dots \right\}^{n}$$

$$= A_{2} a_{0}^{2} w^{2} \left\{ 1 + \frac{a_{1}}{2a_{0}} w + \frac{a_{2}}{3a_{0}} w^{2} + \dots \right\}^{2}$$

$$+ A_3 a_0^3 w^3 \left\{ 1 + \frac{a_1}{2a_0} w + \frac{a_2}{3a_0} w^2 + \dots \right\}^3$$
  
+  $A_4 a_0^4 w^4 \left\{ 1 + \frac{a_1}{2a_0} w + \frac{a_2}{3a_0} w^2 + \dots \right\}^4 + \dots$ 

$$\begin{array}{ll} \longrightarrow & \\ w^2: & 1 = A_2 \, a_0^2 \\ w^3: & 0 = A_2 \, a_0^2 \, \frac{a_1}{a_0} + A_3 \, a_0^3 \\ w^4: & 0 = A_2 \, a_0^2 \left\{ \left( \frac{a_1}{2 \, a_0} \right)^2 + 2 \, \frac{a_2}{3 \, a_0} \right\} + A_3 \, a_0^3 \cdot \frac{3 \, a_1}{2 \, a_0} + A_4 \, a_0^4 \end{array}$$

Hence:

$$a_0 = \frac{1}{\sqrt{A_2}}$$

$$a_1 = -\frac{A_3}{A_2} a_0^2 = -\frac{A_3}{A_2^2}$$

$$0 = A_2 \frac{a_1^2}{4} + \frac{2}{3} A_2 a_0 a_2 + \frac{3}{2} A_3 a_0^2 a_1 + A_4 a_0^4$$

$$= \frac{1}{4} \frac{A_3^2}{A_2^3} + \frac{2}{3} \sqrt{A_2} a_2 - \frac{3}{2} \frac{A_3^2}{A_2^3} + \frac{A_4}{A_2^2}$$

$$\Rightarrow a_2 = \frac{3}{2\sqrt{A_2}} \left\{ \frac{5}{4} \frac{A_3^2}{A_2^3} - \frac{A_4}{A_2^2} \right\}$$

It is clear the process is tedious & error prone.

Anyone who intend to further develope the proceedings is strongly advised to use one of the symbolic manipulation programs such as mathematica.

# ■ Example $\Gamma(z)$