

2.23. Heisenberg Operator Approach to Time Evolution Amplitude

2.23.1. Free Particle

Consider the matrix elements of the evolution operator

$$\langle \mathbf{x} t | \mathbf{x}' 0 \rangle = \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle \quad (2.755)$$

where

$$\hat{H} = H(\hat{\mathbf{p}}) = \frac{\hat{\mathbf{p}}^2}{2M} \quad (2.756)$$

Taking the time derivatives on both sides, (2.755) becomes

$$\begin{aligned} i\hbar \partial_t \langle \mathbf{x} t | \mathbf{x}' 0 \rangle &= \langle \mathbf{x} | \hat{H} e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle \\ &= \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} e^{i\hat{H}t/\hbar} \hat{H} e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle \\ &= \langle \mathbf{x} t | e^{i\hat{H}t/\hbar} \hat{H} e^{-i\hat{H}t/\hbar} | \mathbf{x}' 0 \rangle & | t \rangle &= e^{i\hat{H}t/\hbar} | 0 \rangle \\ &= \langle \mathbf{x} t | \hat{H}(t) | \mathbf{x}' 0 \rangle & \hat{O}(t) &= e^{i\hat{H}t/\hbar} \hat{O} e^{-i\hat{H}t/\hbar} \\ &= \langle \mathbf{x} t | H[\hat{\mathbf{p}}(t)] | \mathbf{x}' 0 \rangle \end{aligned} \quad (2.757)$$

Now,

$$\begin{aligned} \hat{\mathbf{x}}(t) | \mathbf{x} t \rangle &= e^{i\hat{H}t/\hbar} \hat{\mathbf{x}} e^{-i\hat{H}t/\hbar} | \mathbf{x} t \rangle = e^{i\hat{H}t/\hbar} \hat{\mathbf{x}} | \mathbf{x} \rangle = e^{i\hat{H}t/\hbar} | \mathbf{x} \rangle \mathbf{x} = | \mathbf{x} t \rangle \mathbf{x} \\ \rightarrow \langle \mathbf{x} t | \hat{\mathbf{x}}(t) &= \mathbf{x} \langle \mathbf{x} t | \end{aligned} \quad (2.759)$$

$$\hat{\mathbf{x}}(0) | \mathbf{x}' 0 \rangle = \hat{\mathbf{x}} | \mathbf{x}' \rangle = \mathbf{x}' | \mathbf{x}' \rangle = \mathbf{x}' | \mathbf{x}' 0 \rangle \quad (2.759a)$$

Therefore, if we can express $H[\hat{\mathbf{p}}(t)]$ in a time-ordered form

$$\hat{H} = H[\hat{\mathbf{x}}(t), \hat{\mathbf{x}}(0); t] \quad (2.758)$$

in which all $\hat{\mathbf{x}}(t)$'s stand to the left of all $\hat{\mathbf{x}}(0)$'s, then the R.H.S. of (2.757) can be evaluated using (2.759-a) as

$$\langle \mathbf{x} t | H[\hat{\mathbf{x}}(t), \hat{\mathbf{x}}(0); t] | \mathbf{x}' 0 \rangle = H(\mathbf{x}, \mathbf{x}'; t) \langle \mathbf{x} t | \mathbf{x}' 0 \rangle \quad (2.760)$$

This turns (2.757) into a differential eq.

$$i\hbar \partial_t \langle \mathbf{x} t | \mathbf{x}' 0 \rangle = H(\mathbf{x}, \mathbf{x}'; t) \langle \mathbf{x} t | \mathbf{x}' 0 \rangle \quad (2.761)$$

with solution

$$\begin{aligned} \langle \mathbf{x} t | \mathbf{x}' 0 \rangle &= C(\mathbf{x}, \mathbf{x}') E(\mathbf{x}, \mathbf{x}'; t) \\ &= C(\mathbf{x}, \mathbf{x}') \exp\left[-\frac{i}{\hbar} \int^t dt' H(\mathbf{x}, \mathbf{x}'; t')\right] \end{aligned} \quad (2.762)$$

where $C(\mathbf{x}, \mathbf{x}')$ is an integration constant.

The time-ordered form (2.758) can be obtained by solving the Heisenberg eqs. of motion

$$\frac{d\hat{\mathbf{x}}(t)}{dt} = \frac{1}{i\hbar} [\hat{\mathbf{x}}(t), \hat{H}] = \frac{1}{i\hbar} e^{i\hat{H}t/\hbar} \left[\hat{\mathbf{x}}, \frac{\hat{\mathbf{p}}^2}{2M} \right] e^{-i\hat{H}t/\hbar} = \frac{\hat{\mathbf{p}}(t)}{M} \quad (2.763)$$

$$\frac{d\hat{\mathbf{p}}(t)}{dt} = \frac{1}{i\hbar} [\hat{\mathbf{p}}(t), \hat{H}] = \frac{1}{i\hbar} e^{i\hat{H}t/\hbar} \left[\hat{\mathbf{p}}, \frac{\hat{\mathbf{p}}^2}{2M} \right] e^{-i\hat{H}t/\hbar} = 0 \quad (2.764)$$

(2.764) gives

$$\hat{\mathbf{p}}(t) = \hat{\mathbf{p}}(0) = \hat{\mathbf{p}} \quad (2.765)$$

so that the solution to (2.763) is simply

$$\hat{\mathbf{x}}(t) = \hat{\mathbf{x}}(0) + \frac{\hat{\mathbf{p}}}{M} t \quad (2.766)$$

(2.756) thus becomes

$$\hat{H} = \frac{1}{2} M \frac{[\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)]^2}{t^2} \quad (2.767)$$

$$\begin{aligned} &= \frac{M}{2t^2} [\hat{\mathbf{x}}(t)^2 - \hat{\mathbf{x}}(t) \cdot \hat{\mathbf{x}}(0) - \hat{\mathbf{x}}(t) \cdot \hat{\mathbf{x}}(0) + \hat{\mathbf{x}}(0)^2] \\ &= \frac{M}{2t^2} \left\{ \hat{\mathbf{x}}(t)^2 - 2 \hat{\mathbf{x}}(t) \cdot \hat{\mathbf{x}}(0) + \hat{\mathbf{x}}(0)^2 + \sum_{j=1}^D [\hat{x}_j(t), \hat{x}_j(0)] \right\} \end{aligned} \quad (2.768)$$

Using (2.766), we have

$$[\hat{x}_i(t), \hat{x}_j(0)] = \left[\hat{x}_i + \frac{\hat{p}_i}{M} t, \hat{x}_j \right] = -i \frac{\hbar t}{M} \delta_{ij} \quad (2.769a)$$

$$\rightarrow \sum_{j=1}^D [\hat{x}_j(t), \hat{x}_j(0)] = -i \frac{\hbar t}{M} D \quad (2.769)$$

(2.768) thus takes the time-ordered form

$$H[\hat{\mathbf{x}}(t), \hat{\mathbf{x}}(0); t] = \frac{M}{2t^2} \left\{ \hat{\mathbf{x}}(t)^2 - 2 \hat{\mathbf{x}}(t) \cdot \hat{\mathbf{x}}(0) + \hat{\mathbf{x}}(0)^2 \right\} - i \hbar \frac{D}{2t} \quad (2.770)$$

According to (2.760), we have

$$\begin{aligned} H(\mathbf{x}, \mathbf{x}'; t) &= H[\hat{\mathbf{x}}(t), \hat{\mathbf{x}}(0); t] \Big|_{\hat{\mathbf{x}}(t) \rightarrow \mathbf{x}, \hat{\mathbf{x}}(0) \rightarrow \mathbf{x}'} \\ &= \frac{M}{2t^2} (\mathbf{x}^2 - 2 \mathbf{x} \cdot \mathbf{x}' + \mathbf{x}'^2) - i \hbar \frac{D}{2t} \\ &= \frac{M}{2t^2} (\mathbf{x} - \mathbf{x}')^2 - i \hbar \frac{D}{2t} \end{aligned} \quad (2.771)$$

The exponential factor in (2.762) thus becomes

$$\begin{aligned} E(\mathbf{x}, \mathbf{x}'; t) &= \exp \left[-\frac{i}{\hbar} \int^t dt' H(\mathbf{x}, \mathbf{x}'; t') \right] \\ &= \exp \left\{ -\frac{i}{\hbar} \int^t dt' \left[\frac{M}{2t'^2} (\mathbf{x} - \mathbf{x}')^2 - i \hbar \frac{D}{2t'} \right] \right\} \\ &= \exp \left\{ -\frac{i}{\hbar} \left[-\frac{M}{2t} (\mathbf{x} - \mathbf{x}')^2 - i \hbar \frac{D}{2} \ln t \right] \right\} \\ &= \exp \left\{ \frac{i}{\hbar} \frac{M}{2t} (\mathbf{x} - \mathbf{x}')^2 - \frac{D}{2} \ln t \right\} \\ &= \frac{1}{t^{D/2}} \exp \left[\frac{i}{\hbar} \frac{M}{2t} (\mathbf{x} - \mathbf{x}')^2 \right] \end{aligned} \quad (2.772)$$

which turns (2.762) into

$$\langle \mathbf{x} t | \mathbf{x}' 0 \rangle = C(\mathbf{x}, \mathbf{x}') \frac{1}{t^{D/2}} \exp \left[\frac{i}{\hbar} \frac{M}{2t} (\mathbf{x} - \mathbf{x}')^2 \right] \quad (2.773)$$

Note that rigorously, there should be an integration constant from the t' -integral in (2.772). We shall simply merge it into $C(\mathbf{x}, \mathbf{x}')$.

$C(\mathbf{x}, \mathbf{x}')$ can be determined by considering the following differential eqs. obtained by spatial differentiation of (2.755):

$$\frac{\hbar}{i} \nabla \langle \mathbf{x} t | \mathbf{x}' 0 \rangle = \frac{\hbar}{i} \nabla \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle$$

$$\begin{aligned}
&= \langle \mathbf{x} | \hat{\mathbf{p}} e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle & \frac{\hbar}{i} \nabla \langle \mathbf{x} | \phi \rangle = \langle \mathbf{x} | \hat{\mathbf{p}} | \phi \rangle \\
&= \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} e^{i\hat{H}t/\hbar} \hat{\mathbf{p}} e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle \\
&= \langle \mathbf{x} t | \hat{\mathbf{p}}(t) | \mathbf{x}' 0 \rangle \\
\frac{\hbar}{i} \nabla' (\mathbf{x} t | \mathbf{x}' 0) &= \frac{\hbar}{i} \nabla' \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle = \left(-\frac{\hbar}{i} \nabla' \langle \mathbf{x}' | e^{i\hat{H}t/\hbar} | \mathbf{x} \rangle \right)^* \\
&= \left(-\langle \mathbf{x}' | \hat{\mathbf{p}} e^{i\hat{H}t/\hbar} | \mathbf{x} \rangle \right)^* \\
&= -\langle \mathbf{x} | e^{-i\hat{H}t/\hbar} \hat{\mathbf{p}} | \mathbf{x}' \rangle \\
&= -\langle \mathbf{x} t | \hat{\mathbf{p}} | \mathbf{x}' 0 \rangle
\end{aligned} \tag{2.774}$$

Replacing $\hat{\mathbf{p}}$ with (2.766), we have

$$\begin{aligned}
\frac{\hbar}{i} \nabla (\mathbf{x} t | \mathbf{x}' 0) &= \langle \mathbf{x} t | M \frac{\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}}{t} | \mathbf{x}' 0 \rangle \\
&= \frac{M}{t} \langle \mathbf{x} t | (e^{i\hat{H}t/\hbar} \hat{\mathbf{x}} e^{-i\hat{H}t/\hbar} - \hat{\mathbf{x}}) | \mathbf{x}' 0 \rangle \\
&= \frac{M}{t} (\mathbf{x} - \mathbf{x}') (\mathbf{x} t | \mathbf{x}' 0) \\
\frac{\hbar}{i} \nabla' (\mathbf{x} t | \mathbf{x}' 0) &= -\langle \mathbf{x} t | M \frac{\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}}{t} | \mathbf{x}' 0 \rangle \\
&= -\frac{M}{t} (\mathbf{x} - \mathbf{x}') (\mathbf{x} t | \mathbf{x}' 0)
\end{aligned} \tag{2.775}$$

Differentiating (2.773), we have

$$\begin{aligned}
\frac{\hbar}{i} \nabla (\mathbf{x} t | \mathbf{x}' 0) &= \left\{ \frac{\hbar}{i} \frac{\nabla C}{C} + \frac{M}{t} (\mathbf{x} - \mathbf{x}') \right\} (\mathbf{x} t | \mathbf{x}' 0) \\
\frac{\hbar}{i} \nabla' (\mathbf{x} t | \mathbf{x}' 0) &= \left\{ \frac{\hbar}{i} \frac{\nabla' C}{C} - \frac{M}{t} (\mathbf{x} - \mathbf{x}') \right\} (\mathbf{x} t | \mathbf{x}' 0)
\end{aligned}$$

Comparing with (2.775) then gives

$$\nabla C = 0 \quad \text{and} \quad \nabla' C = 0 \tag{2.776}$$

so that C is a true constant independent of \mathbf{x} and \mathbf{x}' . Its value is fixed by the I.C.

$$\lim_{t \rightarrow 0} (\mathbf{x} t | \mathbf{x}' 0) = \langle \mathbf{x} | \mathbf{x}' \rangle = \delta(\mathbf{x} - \mathbf{x}') \tag{2.777}$$

Using

$$\delta(\mathbf{x} - \mathbf{x}') = \lim_{a \rightarrow 0} \frac{1}{(\pi a)^{D/2}} e^{-(\mathbf{x} - \mathbf{x}')^2/a}$$

we can write (2.773) as

$$\begin{aligned}
\lim_{t \rightarrow 0} (\mathbf{x} t | \mathbf{x}' 0) &= \lim_{t \rightarrow 0} C \frac{1}{t^{D/2}} \exp \left[\frac{i}{\hbar} \frac{M}{2t} (\mathbf{x} - \mathbf{x}')^2 \right] \\
&= C \left(\frac{i 2 \pi \hbar}{M} \right)^{D/2} \delta(\mathbf{x} - \mathbf{x}')
\end{aligned}$$

Comparing with (2.777) then gives

$$C = \left(\frac{M}{i 2 \pi \hbar} \right)^{D/2} \tag{2.778}$$

(2.773) thus becomes

$$(\mathbf{x}t | \mathbf{x}'0) = \left(\frac{M}{i2\pi\hbar t} \right)^{D/2} \exp \left[\frac{i}{\hbar} \frac{M}{2t} (\mathbf{x} - \mathbf{x}')^2 \right] \quad (2.779)$$

in agreement with the previous result (2.130).

Note that in this approach, the fluctuation factor $t^{-D/2}$ arises from the commutator relation (2.769).

2.23.2. Harmonic Oscillator

The Hamiltonian

$$\hat{H} = H(\hat{\mathbf{x}}, \hat{\mathbf{p}}) = \frac{\hat{\mathbf{p}}^2}{2M} + \frac{1}{2} M \omega^2 \hat{\mathbf{x}}^2 \quad (2.780)$$

gives rise to eqs. of motion

$$\frac{d\hat{\mathbf{x}}(t)}{dt} = \frac{1}{i\hbar} \left[\hat{\mathbf{x}}(t), \frac{\hat{\mathbf{p}}(t)^2}{2M} \right] = \frac{\hat{\mathbf{p}}(t)}{M} \quad (2.781)$$

$$\frac{d\hat{\mathbf{p}}(t)}{dt} = \frac{1}{i\hbar} \left[\hat{\mathbf{p}}(t), \frac{1}{2} M \omega^2 \hat{\mathbf{x}}(t)^2 \right] = -M \omega^2 \hat{\mathbf{x}}(t) \quad (2.782)$$

The solution is [c.f. (2.156)]

$$\hat{\mathbf{p}}(t) = M \frac{d\hat{\mathbf{x}}(t)}{dt} = \frac{M\omega}{\sin \omega t} [\hat{\mathbf{x}}(t) \cos \omega t - \hat{\mathbf{x}}(0)] \quad (2.783)$$

To verify this claim, we differentiate (2.783) to get

$$\begin{aligned} \frac{d\hat{\mathbf{p}}(t)}{dt} &= -\frac{M\omega^2 \cos \omega t}{\sin^2 \omega t} [\hat{\mathbf{x}}(t) \cos \omega t - \hat{\mathbf{x}}(0)] + \frac{M\omega}{\sin \omega t} \left[\frac{d\hat{\mathbf{x}}(t)}{dt} \cos \omega t - \omega \hat{\mathbf{x}}(t) \sin \omega t \right] \\ &= -\frac{M\omega^2 \cos \omega t}{\sin^2 \omega t} [\hat{\mathbf{x}}(t) \cos \omega t - \hat{\mathbf{x}}(0)] \\ &\quad + \frac{M\omega}{\sin \omega t} \left\{ \frac{\omega}{\sin \omega t} [\hat{\mathbf{x}}(t) \cos \omega t - \hat{\mathbf{x}}(0)] \cos \omega t - \omega \hat{\mathbf{x}}(t) \sin \omega t \right\} \\ &= -M\omega^2 \hat{\mathbf{x}}(t) \end{aligned}$$

so that (2.783) is indeed a solution of (2.782).

(2.780) then becomes

$$\hat{H} = \frac{M\omega^2}{2\sin^2 \omega t} \left\{ [\hat{\mathbf{x}}(t) \cos \omega t - \hat{\mathbf{x}}(0)]^2 + \sin^2 \omega t \hat{\mathbf{x}}(t)^2 \right\} \quad (2.784)$$

$$= \frac{M\omega^2}{2\sin^2 \omega t} \left\{ \hat{\mathbf{x}}(t)^2 - 2\hat{\mathbf{x}}(t) \cdot \hat{\mathbf{x}}(0) \cos \omega t + \hat{\mathbf{x}}(0)^2 + \cos \omega t \sum_{j=1}^D [\hat{x}_j(t), \hat{x}_j(0)] \right\} \quad (2.785)$$

From (2.783), we have

$$\begin{aligned} [\hat{x}_i(t), \hat{p}_j(t)] &= i\hbar \delta_{ij} \\ &= \left[\hat{x}_i(t), \frac{M\omega}{\sin \omega t} \{ \hat{x}_j(t) \cos \omega t - \hat{x}_j(0) \} \right] \\ &= -\frac{M\omega}{\sin \omega t} [\hat{x}_i(t), \hat{x}_j(0)] \end{aligned} \quad (2.786a)$$

so that

$$\sum_{j=1}^D [\hat{x}_j(t), \hat{x}_j(0)] = -\frac{i\hbar \sin \omega t}{M\omega} D \quad (2.786)$$

(2.785) thus implies [c.f.(2.771)]

$$\begin{aligned}
H(\mathbf{x}, \mathbf{x}'; t) &= \frac{M\omega^2}{2\sin^2\omega t} \left(\mathbf{x}^2 - 2\mathbf{x}\cdot\mathbf{x}'\cos\omega t + \mathbf{x}'^2 - \frac{i\hbar\sin\omega t\cos\omega t}{M\omega} D \right) \\
&= \frac{M\omega^2}{2\sin^2\omega t} \left(\mathbf{x}^2 + \mathbf{x}'^2 - 2\mathbf{x}\cdot\mathbf{x}'\cos\omega t \right) - \frac{i\hbar\omega\cot\omega t}{2} D
\end{aligned} \quad (2.787)$$

Using

$$\begin{aligned}
\int dt \frac{1}{\sin^2\omega t} &= -\frac{\cot\omega t}{\omega} & \int dt \frac{\cos\omega t}{\sin^2\omega t} &= -\frac{1}{\omega\sin\omega t} \\
\int dt \cot\omega t &= \frac{1}{\omega} \ln(\sin\omega t)
\end{aligned}$$

we get

$$\int dt H(\mathbf{x}, \mathbf{x}'; t) = -\frac{M\omega}{2\sin\omega t} \left[(\mathbf{x}^2 + \mathbf{x}'^2) \cos\omega t - 2\mathbf{x}\cdot\mathbf{x}' \right] - \frac{i\hbar}{2} D \ln(\sin\omega t) \quad (2.788)$$

Inserting this into (2.762) gives a $\langle \mathbf{x} t | \mathbf{x}' 0 \rangle$ that agrees with (2.175) aside from the prefactor $C(\mathbf{x}, \mathbf{x}')$, which can be determined from differential eqs similar to (2.774).

2.23.3. Charged Particle in Magnetic Field

The Hamiltonian is

$$\hat{H} = \frac{1}{2M} \hat{\mathbf{P}}^2 \quad (2.790)$$

where

$$\hat{\mathbf{P}} = \hat{\mathbf{p}} - \frac{e}{c} \mathbf{A}(\hat{\mathbf{x}}) \quad (2.789)$$

is the **covariant momentum** with non-commuting components:

$$\begin{aligned}
[P_i, P_j] &= \left[p_i, -\frac{e}{c} A_j \right] + \left[-\frac{e}{c} A_i, p_j \right] \\
&= i \frac{e\hbar}{c} (\partial_i A_j - \partial_j A_i) \\
&= i \frac{e\hbar}{c} B_{ij}
\end{aligned} \quad (2.791)$$

where

$$B_{ij} = \epsilon_{ijk} B_k = \begin{pmatrix} 0 & B_3 & -B_2 \\ -B_3 & 0 & B_1 \\ B_2 & -B_1 & 0 \end{pmatrix} = \mathbb{B} \quad (2.791a)$$

is the anti-symmetric tensor representation of the magnetic field $\mathbf{B} = \nabla \times \mathbf{A}$.

The eqs. of motion are

$$\frac{d\hat{\mathbf{x}}(t)}{dt} = \frac{1}{i\hbar} \left[\hat{\mathbf{x}}(t), \frac{\hat{\mathbf{P}}(t)^2}{2M} \right] = \frac{\hat{\mathbf{P}}(t)}{M} \quad (2.792)$$

$$\begin{aligned}
\frac{d\hat{P}_i(t)}{dt} &= \frac{1}{i\hbar} \left[\hat{P}_i(t), \frac{\hat{P}_j(t)\hat{P}_j(t)}{2M} \right] \\
&= \frac{1}{2i\hbar M} \left([\hat{P}_i(t), \hat{P}_j(t)] \hat{P}_j(t) + \hat{P}_j(t) [\hat{P}_i(t), \hat{P}_j(t)] \right) \\
&= \frac{e}{2Mc} \left(B_{ij} \hat{P}_j(t) + \hat{P}_j(t) B_{ij} \right)
\end{aligned} \quad (2.793a)$$

Using

$$\hat{P}_j(t) B_{ij} \psi = \frac{\hbar}{i} \partial_j (B_{ij} \psi) - \frac{e}{c} A_j B_{ij} \psi = B_{ij} \left(\frac{\hbar}{i} \partial_j - \frac{e}{c} A_j \right) \psi + \frac{\hbar}{i} (\partial_j B_{ij}) \psi$$

$$= B_{ij} \hat{P}_j(t) \psi + \frac{\hbar}{i} (\partial_j B_{ij}) \psi$$

(2.793a) becomes

$$\begin{aligned} \frac{d \hat{P}_i(t)}{dt} &= \frac{e}{Mc} B_{ij} \hat{P}_j(t) + \frac{e \hbar}{2 M c i} (\partial_j B_{ij}) \\ &= \frac{e}{Mc} B_{ij} \hat{P}_j(t) + i \frac{e \hbar}{2 M c} (\partial_j B_{ji}) \quad (B_{ij} = -B_{ji}) \end{aligned} \quad (2.793)$$

or in vector form

$$\frac{d \hat{\mathbf{P}}(t)}{dt} = \frac{e}{Mc} \mathbf{B} \cdot \hat{\mathbf{P}}(t) + i \frac{e \hbar}{2 M c} (\nabla \cdot \mathbf{B}) \quad (2.793b)$$

In the following, we shall restrict ourselves to a constant field so that $\nabla \cdot \mathbf{B} = 0$ and (2.793) reduces to

$$\frac{d \hat{P}_i(t)}{dt} = \frac{e}{Mc} B_{ij} \hat{P}_j(t)$$

with solution

$$\hat{P}_i(t) = \exp\left(\frac{e}{Mc} B_{ij} t\right) \hat{P}_j(0) = e^{W_{Lij} t} \hat{P}_j(0) \quad (2.794a)$$

where

$$W_{Lij} = \frac{e}{Mc} B_{ij} \quad (2.795)$$

are components of the **Landau frequency matrix**

$$\mathbb{W}_L = \frac{e}{Mc} \mathbf{B} \quad (2.795a)$$

Note: we've used W_{Lij} instead of Kleinert's Ω_{Lij} because there is no double-stroke font for Ω .

(2.794a) can be written in vector form as

$$\hat{\mathbf{P}}(t) = e^{\mathbb{W}_L t} \cdot \hat{\mathbf{P}}(0) \quad (2.794)$$

Using the **Landau frequency vector**

$$\boldsymbol{\omega}_L = \frac{e}{Mc} \mathbf{B} \quad (2.796)$$

and the 3-D generators \mathbb{L}_k of the rotation group with components

$$(\mathbb{L}_k)_{ij} = -i \epsilon_{kij} = \mathbb{L}_k \quad (2.797)$$

(2.795) becomes

$$W_{Lij} = \frac{e}{Mc} \epsilon_{ijk} B_k = \frac{e}{Mc} i (\mathbb{L}_k)_{ij} B_k = i (\mathbb{L}_k)_{ij} \omega_{Lk} \quad (2.797a)$$

or, in matrix form,

$$\mathbb{W}_L = i \mathbb{L}_k \omega_{Lk} = i \mathbb{L} \cdot \boldsymbol{\omega}_L \quad (2.798)$$

where \mathbb{L} is a vector with components \mathbb{L}_k that are themselves matrices with components $(\mathbb{L}_k)_{ij}$.

Inserting (2.794) into (2.792) gives

$$\frac{d \hat{\mathbf{x}}(t)}{dt} = \frac{1}{M} e^{\mathbb{W}_L t} \cdot \hat{\mathbf{P}}(0) \quad (2.798a)$$

Using

$$e^{\mathbb{W}_L t} = \sum_{n=0}^{\infty} \frac{t^n}{n!} \mathbb{W}_L^n \quad (2.798b)$$

we have

$$\frac{d e^{\mathbf{W}_L t}}{d t} = \sum_{n=1}^{\infty} \frac{t^{n-1}}{(n-1)!} \mathbf{W}_L^n = \sum_{n=0}^{\infty} \frac{t^n}{n!} \mathbf{W}_L^{n+1} = \mathbf{W}_L \cdot e^{\mathbf{W}_L t} \quad (2.798c)$$

The solution of (2.798a) is then easily proved to be

$$\begin{aligned} \hat{\mathbf{x}}(t) &= \hat{\mathbf{x}}(0) + \frac{1}{M} \mathbf{W}_L^{-1} \cdot (e^{\mathbf{W}_L t} - \mathbb{I}) \cdot \hat{\mathbf{P}}(0) \\ &= \hat{\mathbf{x}}(0) + \frac{1}{M} (e^{\mathbf{W}_L t} - \mathbb{I}) \cdot \mathbf{W}_L^{-1} \cdot \hat{\mathbf{P}}(0) \\ &= \hat{\mathbf{x}}(0) + \frac{1}{M} \frac{e^{\mathbf{W}_L t} - \mathbb{I}}{\mathbf{W}_L} \cdot \hat{\mathbf{P}}(0) \end{aligned} \quad (2.799)$$

where \mathbb{I} is the 3×3 unit matrix.

Using (2.798b), we have

$$\begin{aligned} \mathbf{W}_L^{-1} \cdot (e^{\mathbf{W}_L t} - 1) &= \mathbf{W}_L^{-1} \cdot \left(\sum_{n=1}^{\infty} \frac{t^n}{n!} \mathbf{W}_L^n \right) = \sum_{n=1}^{\infty} \frac{t^n}{n!} \mathbf{W}_L^{n-1} \\ &= t \mathbb{I} + \frac{t^2}{2!} \mathbf{W}_L + \frac{t^3}{3!} \mathbf{W}_L^2 + \dots \end{aligned} \quad (2.800)$$

Since the terms involving \mathbf{W}_L all commute among themselves, (2.799) can be inverted to give

$$\begin{aligned} \frac{\hat{\mathbf{P}}(0)}{M} &= \frac{\mathbf{W}_L}{e^{\mathbf{W}_L t} - \mathbb{I}} \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \\ &= \frac{\mathbf{W}_L}{e^{\mathbf{W}_L t/2} (e^{\mathbf{W}_L t/2} - e^{-\mathbf{W}_L t/2})} \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \\ &= \frac{\mathbf{W}_L \cdot e^{-\mathbf{W}_L t/2}}{2 \sinh \frac{\mathbf{W}_L t}{2}} \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \end{aligned} \quad (2.801)$$

(2.794) then becomes

$$\begin{aligned} \hat{\mathbf{P}}(t) &= M \frac{\mathbf{W}_L \cdot e^{\mathbf{W}_L t/2}}{2 \sinh \frac{\mathbf{W}_L t}{2}} \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \\ &= M \mathbb{N}(\mathbf{W}_L t) \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \end{aligned} \quad (2.802)$$

where

$$\mathbb{N}(\mathbf{W}_L t) = \frac{\mathbf{W}_L \cdot e^{\mathbf{W}_L t/2}}{2 \sinh \frac{\mathbf{W}_L t}{2}} \quad (2.803)$$

(2.802) implies

$$\begin{aligned} \hat{H} &= \frac{\hat{\mathbf{P}}(t)^2}{2M} = \frac{\hat{\mathbf{P}}(t)^T \cdot \hat{\mathbf{P}}(t)}{2M} \\ &= \frac{1}{2} M [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)]^T \cdot \mathbb{N}^T(\mathbf{W}_L t) \cdot \mathbb{N}(\mathbf{W}_L t) \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \\ &= \frac{1}{2} M [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)]^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] \end{aligned} \quad (2.804)$$

where

$$\mathbb{K}(\mathbf{W}_L t) = \mathbb{N}^T(\mathbf{W}_L t) \cdot \mathbb{N}(\mathbf{W}_L t) \quad (2.805)$$

From (2.795a), we see that

$$\mathbb{B}^T = -\mathbb{B} \quad \rightarrow \quad \mathbf{W}_L^T = -\mathbf{W}_L \quad (2.805a)$$

Therefore, (2.803) gives

$$\mathbb{N}^T(\mathbf{W}_L t) = \mathbb{N}(-\mathbf{W}_L t) = \frac{\mathbf{W}_L \cdot e^{-\mathbf{W}_L t/2}}{2 \sinh \frac{\mathbf{W}_L t}{2}}$$

$$\rightarrow \mathbb{K}(\mathbf{W}_L t) = \frac{\mathbf{W}_L^2}{4 \sinh^2 \frac{\mathbf{W}_L t}{2}} \tag{2.806}$$

Using

$$[\hat{x}_i, \hat{p}_j] = \left[\hat{x}_i, \hat{p}_j - \frac{e}{c} A_j(\hat{\mathbf{x}}) \right] = [\hat{x}_i, \hat{p}_j] = i \hbar \delta_{ij} \tag{2.806a}$$

(2.799) implies

$$[\hat{x}_i(t), \hat{x}_j(0)] = \left[\frac{1}{M} \left(\frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right)_{ik} \hat{p}_k, \hat{x}_j \right]$$

$$= -i \hbar \frac{1}{M} \left(\frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right)_{ij} \tag{2.807}$$

so that

$$[\hat{x}_i(t), \hat{x}_j(0)] + [\hat{x}_j(t), \hat{x}_i(0)] = -i \hbar \frac{1}{M} \left\{ \left(\frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right)_{ij} + \left(\frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right)_{ji} \right\}$$

$$= -i \hbar \frac{1}{M} \left\{ \frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} + \left(\frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right)^T \right\}_{ij}$$

$$= -i \hbar \frac{1}{M} \left(\frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} - \frac{e^{-\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right)_{ij}$$

$$= -i \hbar \frac{1}{M} \left(\frac{e^{\mathbf{W}_L t} - e^{-\mathbf{W}_L t}}{\mathbf{W}_L} \right)_{ij}$$

$$= -i \hbar \frac{2}{M} \left(\frac{\sinh \mathbf{W}_L t}{\mathbf{W}_L} \right)_{ij} \tag{2.808}$$

Expanding (2.804), we have

$$\hat{H} = \frac{1}{2} M \left\{ \hat{\mathbf{x}}^T(t) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(t) - \hat{\mathbf{x}}^T(t) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(0) \right.$$

$$\left. - \hat{\mathbf{x}}^T(0) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(t) + \hat{\mathbf{x}}^T(0) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(0) \right\} \tag{2.808a}$$

Using

$$\hat{\mathbf{x}}^T(t) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(0) - \hat{\mathbf{x}}^T(0) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(t)$$

$$= \left(\hat{\mathbf{x}}^T - \frac{1}{M} \hat{\mathbf{p}}^T \cdot \frac{e^{-\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \right) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}} - \hat{\mathbf{x}}^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \left(\hat{\mathbf{x}} + \frac{1}{M} \frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \cdot \hat{\mathbf{p}} \right)$$

$$= -\frac{1}{M} \hat{\mathbf{p}}^T \cdot \frac{e^{-\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}} - \frac{1}{M} \hat{\mathbf{x}}^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \cdot \hat{\mathbf{p}}$$

$$= -\frac{1}{M} \hat{\mathbf{p}}^T \cdot \frac{e^{-\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \cdot \frac{\mathbf{W}_L^2}{4 \sinh^2 \frac{\mathbf{W}_L t}{2}} \cdot \hat{\mathbf{x}} - \frac{1}{M} \hat{\mathbf{x}}^T \cdot \frac{\mathbf{W}_L^2}{4 \sinh^2 \frac{\mathbf{W}_L t}{2}} \cdot \frac{e^{\mathbf{W}_L t} - \mathbb{1}}{\mathbf{W}_L} \cdot \hat{\mathbf{p}}$$

$$= \frac{1}{M} \hat{\mathbf{p}}^T \cdot \frac{e^{-\mathbf{W}_L t/2} \cdot \mathbf{W}_L}{2 \sinh \frac{\mathbf{W}_L t}{2}} \cdot \hat{\mathbf{x}} - \frac{1}{M} \hat{\mathbf{x}}^T \cdot \frac{\mathbf{W}_L \cdot e^{\mathbf{W}_L t/2}}{2 \sinh \frac{\mathbf{W}_L t}{2}} \cdot \hat{\mathbf{p}}$$

$$= \frac{1}{M} \left(\frac{e^{-\mathbf{W}_L t/2} \cdot \mathbf{W}_L}{2 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ij} \hat{p}_i \hat{x}_j - \frac{1}{M} \left(\frac{\mathbf{W}_L \cdot e^{\mathbf{W}_L t/2}}{2 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ji} \hat{x}_j \hat{p}_i$$

$$\begin{aligned}
&= \frac{1}{M} \left(\frac{e^{-\mathbf{W}_L t/2} \cdot \mathbf{W}_L}{2 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ij} \hat{P}_i \hat{x}_j - \frac{1}{M} \left(\frac{e^{-\mathbf{W}_L t/2} \cdot \mathbf{W}_L}{2 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ij} \hat{x}_i \hat{P}_j \quad [\mathbf{W}_L^T = -\mathbf{W}_L \text{ used. }] \\
&= \frac{1}{M} \left(\frac{e^{-\mathbf{W}_L t/2} \cdot \mathbf{W}_L}{2 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ij} (-i \hbar \delta_{ij}) \quad [(2.806a) \text{ used.}] \\
&= -\frac{i \hbar}{M} \left(\frac{e^{-\mathbf{W}_L t/2} \cdot \mathbf{W}_L}{2 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ii} \\
&= -\frac{i \hbar}{M} \left(\frac{(e^{-\mathbf{W}_L t/2} + e^{\mathbf{W}_L t/2}) \cdot \mathbf{W}_L}{4 \sinh \frac{\mathbf{W}_L t}{2}} \right)_{ii} \quad [\mathbf{A}_{ij} = \mathbf{A}_{ij}^T \text{ used. }] \\
&= -\frac{i \hbar}{M} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right)_{ii} \\
&= -\frac{i \hbar}{M} \text{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) \quad (2.808b)
\end{aligned}$$

(2.808a) becomes

$$\begin{aligned}
\hat{H} = \frac{1}{2} M \left\{ \hat{\mathbf{x}}^T(t) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(t) - 2 \hat{\mathbf{x}}^T(t) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(0) \right. \\
\left. + \hat{\mathbf{x}}^T(0) \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \hat{\mathbf{x}}(0) - \frac{i \hbar}{M} \text{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) \right\} \quad (2.809)
\end{aligned}$$

Using the recipe prescribed in (2.771), we get

$$\begin{aligned}
H(\mathbf{x}, \mathbf{x}'; t) &= \frac{1}{2} M \left\{ \mathbf{x}^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \mathbf{x} - 2 \mathbf{x}^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \mathbf{x}' + \mathbf{x}'^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \mathbf{x}' \right\} \\
&\quad - \frac{i \hbar}{2} \text{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) \\
&= \frac{1}{2} M (\mathbf{x} - \mathbf{x}')^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot (\mathbf{x} - \mathbf{x}') - \frac{i \hbar}{2} \text{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) \quad (2.809a)
\end{aligned}$$

where we've used

$$\mathbf{x}^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \mathbf{x}' = \mathbf{x}'^T \cdot \mathbb{K}(\mathbf{W}_L t) \cdot \mathbf{x}$$

since (2.806) gives

$$\mathbb{K}^T(\mathbf{W}_L t) = \mathbb{K}(\mathbf{W}_L t)$$

The next step is to calculate

$$E(\mathbf{x}, \mathbf{x}'; t) = \exp \left[-\frac{i}{\hbar} \int_0^t dt' H(\mathbf{x}, \mathbf{x}'; t') \right]$$

From (2.806), we have

$$\begin{aligned}
\int dt \mathbb{K}(\mathbf{W}_L t) &= \int dt \frac{\mathbf{W}_L^2}{4 \sinh^2 \frac{\mathbf{W}_L t}{2}} \quad [\int dt \frac{1}{\sinh^2(at)} = -\frac{1}{a} \coth(at)] \\
&= -\frac{1}{2} \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \quad (2.810)
\end{aligned}$$

Let w_n be the eigenvalues of \mathbf{W}_L . Using

$$\int dt \coth(at) = \frac{1}{a} \ln[\sinh(at)]$$

we have

$$\int dt \operatorname{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) = \int dt \sum_n \frac{1}{2} w_n \coth \frac{w_n t}{2} = \sum_n \ln \left[\sinh \left(\frac{w_n t}{2} \right) \right] + C \quad (2.810a)$$

where C is a constant.

Now

$$\mathbf{W}_L^T = -\mathbf{W}_L \quad \rightarrow \quad \operatorname{tr} \mathbf{W}_L^T = \operatorname{tr} \mathbf{W}_L = -\operatorname{tr} \mathbf{W}_L = 0 = \sum_n w_n \quad (2.810b)$$

Hence, for some n , $w_n < 0$, thus making $\ln \left[\sinh \left(\frac{w_n t}{2} \right) \right]$ imaginary. However, the diagonal elements $E(\mathbf{x}, \mathbf{x}; t)$ should be real. We therefore guarantee the argument of each logarithm to be positive by setting

$$C = -\sum_n \ln \left(\frac{w_n}{2} \right)$$

so that

$$\begin{aligned} \int dt \operatorname{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) &= \sum_n \ln \left[\frac{\sinh(w_n t / 2)}{w_n / 2} \right] \\ &= \sum_n \ln \left[\frac{\sinh(w_n t / 2)}{w_n t / 2} \right] + \sum_{n=1}^D \ln t \\ &= \operatorname{tr} \ln \left[\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right] + D \ln t \end{aligned} \quad (2.811)$$

where D is the dimension of the system.

Let a_n be the eigenvalues of \mathbf{A} . Then

$$\operatorname{tr} \ln \mathbf{A} = \sum_n \ln a_n = \ln \left(\prod_n a_n \right) = \ln \det \mathbf{A}$$

Hence,

$$\begin{aligned} \int dt \operatorname{tr} \left(\frac{1}{2} \mathbf{W}_L \cdot \coth \frac{\mathbf{W}_L t}{2} \right) &= \ln \left[\det \left(\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right) \right] + D \ln t \\ &= \ln \left[t^D \det \left(\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right) \right] \end{aligned} \quad (2.811a)$$

(2.809a) then gives

$$\int dt H(\mathbf{x}, \mathbf{x}'; t) = -\frac{1}{2} M(\mathbf{x} - \mathbf{x}')^T \cdot \frac{1}{2} \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') \quad (2.811b)$$

$$\begin{aligned} &\quad -\frac{i\hbar}{2} \left\{ \operatorname{tr} \ln \left[\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right] + D \ln t \right\} \\ &= -\frac{1}{2} M(\mathbf{x} - \mathbf{x}')^T \cdot \frac{1}{2} \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') \quad (2.811c) \\ &\quad -\frac{i\hbar}{2} \ln \left[t^D \det \left(\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right) \right] \end{aligned}$$

and

$$E(\mathbf{x}, \mathbf{x}'; t) = t^{-D/2} \exp \left\{ -\frac{1}{2} \operatorname{tr} \ln \left[\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right] \right\} \quad (2.812)$$

$$\begin{aligned} & \times \exp \left[\frac{i}{\hbar} \frac{1}{2} M (\mathbf{x} - \mathbf{x}')^T \cdot \frac{1}{2} \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') \right] \\ & = t^{-D/2} \left[\det \left(\frac{\sinh(\mathbf{W}_L t / 2)}{\mathbf{W}_L t / 2} \right) \right]^{-1/2} \\ & \times \exp \left[\frac{i}{\hbar} \frac{1}{2} M (\mathbf{x} - \mathbf{x}')^T \cdot \frac{1}{2} \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') \right] \end{aligned} \quad (2.813)$$

Next, we need to fix $C(\mathbf{x}, \mathbf{x}')$ of (2.762) by solving the differential eqs. obtained by spatial differentiation of the evolution amplitude

$$(\mathbf{x} t | \mathbf{x}' 0) = \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle$$

where

$$\hat{H} = \frac{1}{2M} \hat{\mathbf{P}}^2 \quad \hat{\mathbf{P}} = \hat{\mathbf{p}} - \frac{e}{c} \mathbf{A}(\hat{\mathbf{x}})$$

Owing to the form of \hat{H} , it is easier to work with the covariant derivatives so that

$$\begin{aligned} \left(\frac{\hbar}{i} \nabla - \frac{e}{c} \mathbf{A}(\mathbf{x}) \right) (\mathbf{x} t | \mathbf{x}' 0) &= \langle \mathbf{x} | \hat{\mathbf{P}} e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle \\ &= \langle \mathbf{x} | e^{-i\hat{H}t/\hbar} e^{i\hat{H}t/\hbar} \hat{\mathbf{P}} e^{-i\hat{H}t/\hbar} | \mathbf{x}' \rangle \\ &= \langle \mathbf{x} t | \hat{\mathbf{P}}(t) | \mathbf{x}' \rangle \\ &= M \mathbf{N}(\mathbf{W}_L t) \cdot \langle \mathbf{x} t | [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] | \mathbf{x}' \rangle \quad [(2.802) \text{ used.}] \\ &= M \mathbf{N}(\mathbf{W}_L t) \cdot (\mathbf{x} - \mathbf{x}') (\mathbf{x} t | \mathbf{x}' 0) \end{aligned} \quad (2.814)$$

Similarly,

$$\begin{aligned} \left(\frac{\hbar}{i} \nabla' + \frac{e}{c} \mathbf{A}(\mathbf{x}') \right) (\mathbf{x} t | \mathbf{x}' 0) &= -\langle \mathbf{x} | e^{-i\hat{H}t/\hbar} \hat{\mathbf{P}} | \mathbf{x}' \rangle \\ &= -\langle \mathbf{x} t | \hat{\mathbf{P}}(0) | \mathbf{x}' \rangle \\ &= -M \mathbf{N}^T(\mathbf{W}_L t) \cdot \langle \mathbf{x} t | [\hat{\mathbf{x}}(t) - \hat{\mathbf{x}}(0)] | \mathbf{x}' \rangle \quad [(2.801) \text{ used.}] \\ &= -M \mathbf{N}^T(\mathbf{W}_L t) \cdot (\mathbf{x} - \mathbf{x}') (\mathbf{x} t | \mathbf{x}' 0) \end{aligned} \quad (2.815)$$

Also,

$$\begin{aligned} (\mathbf{x} - \mathbf{x}')^T \cdot \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) &= (x_i - x'_i) W_{Lij} \coth \left(\frac{\mathbf{W}_L t}{2} \right)_{jk} \\ \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot \mathbf{W}_L \cdot (\mathbf{x} - \mathbf{x}') &= \coth \left(\frac{\mathbf{W}_L t}{2} \right)_{kj} W_{Lji} (x_i - x'_i) \\ &= \left[-\coth \left(\frac{\mathbf{W}_L t}{2} \right)_{jk} \right] (-W_{Lij}) (x_i - x'_i) \\ \rightarrow (\mathbf{x} - \mathbf{x}')^T \cdot \mathbf{W}_L \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) &= \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot \mathbf{W}_L \cdot (\mathbf{x} - \mathbf{x}') \end{aligned}$$

(2.813) then gives

$$\frac{\hbar}{i} \nabla E = M \frac{\mathbf{W}_L}{2} \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') E \quad (2.815a)$$

$$\frac{\hbar}{i} \nabla' E = -M \frac{\mathbf{W}_L}{2} \cdot \coth \left(\frac{\mathbf{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') E \quad (2.815b)$$

From (2.762)

$$(\mathbf{x}t | \mathbf{x}'0) = C(\mathbf{x}, \mathbf{x}') E(\mathbf{x}, \mathbf{x}'; t)$$

we get

$$\begin{aligned} \frac{\hbar}{i} \nabla(\mathbf{x}t | \mathbf{x}'0) &= \left(\frac{\hbar}{i} \nabla C \right) E + C \frac{\hbar}{i} \nabla E \\ &= \left(\frac{\hbar}{i} \nabla C \right) E + CM \frac{\mathbb{W}_L}{2} \cdot \coth\left(\frac{\mathbb{W}_L t}{2}\right) \cdot (\mathbf{x} - \mathbf{x}') E \end{aligned} \quad (2.815c)$$

$$\begin{aligned} \frac{\hbar}{i} \nabla'(\mathbf{x}t | \mathbf{x}'0) &= \left(\frac{\hbar}{i} \nabla' C \right) E + C \frac{\hbar}{i} \nabla' E \\ &= \left(\frac{\hbar}{i} \nabla' C \right) E - CM \frac{\mathbb{W}_L}{2} \cdot \coth\left(\frac{\mathbb{W}_L t}{2}\right) \cdot (\mathbf{x} - \mathbf{x}') E \end{aligned} \quad (2.815d)$$

(2.814-5) thus become

$$\left\{ \frac{\hbar}{i} \nabla - \frac{e}{c} \mathbf{A}(\mathbf{x}) + M \frac{\mathbb{W}_L}{2} \cdot \left[\coth\left(\frac{\mathbb{W}_L t}{2}\right) - \mathbb{N}(\mathbb{W}_L t) \right] \cdot (\mathbf{x} - \mathbf{x}') \right\} C = 0 \quad (2.816a)$$

$$\left\{ \frac{\hbar}{i} \nabla' + \frac{e}{c} \mathbf{A}(\mathbf{x}') + M \frac{\mathbb{W}_L}{2} \cdot \left[\coth\left(\frac{\mathbb{W}_L t}{2}\right) - \mathbb{N}^T(\mathbb{W}_L t) \right] \cdot (\mathbf{x}' - \mathbf{x}) \right\} C = 0 \quad (2.816b)$$

Using (2.803),

$$\mathbb{N}(\mathbb{W}_L t) = \frac{\mathbb{W}_L \cdot e^{\mathbb{W}_L t/2}}{2 \sinh \frac{\mathbb{W}_L t}{2}} \quad \mathbb{N}^T(\mathbb{W}_L t) = \frac{\mathbb{W}_L \cdot e^{-\mathbb{W}_L t/2}}{2 \sinh \frac{\mathbb{W}_L t}{2}}$$

we have

$$\begin{aligned} \frac{\mathbb{W}_L}{2} \cdot \coth\left(\frac{\mathbb{W}_L t}{2}\right) - \mathbb{N}(\mathbb{W}_L t) &= \frac{\mathbb{W}_L \cdot (e^{\mathbb{W}_L t/2} + e^{-\mathbb{W}_L t/2})}{4 \sinh \frac{\mathbb{W}_L t}{2}} - \frac{\mathbb{W}_L \cdot e^{\mathbb{W}_L t/2}}{2 \sinh \frac{\mathbb{W}_L t}{2}} \\ &= \frac{\mathbb{W}_L \cdot (-e^{\mathbb{W}_L t/2} + e^{-\mathbb{W}_L t/2})}{4 \sinh \frac{\mathbb{W}_L t}{2}} = -\frac{\mathbb{W}_L}{2} \end{aligned}$$

$$\begin{aligned} \frac{\mathbb{W}_L}{2} \cdot \coth\left(\frac{\mathbb{W}_L t}{2}\right) - \mathbb{N}^T(\mathbb{W}_L t) &= \frac{\mathbb{W}_L \cdot (e^{\mathbb{W}_L t/2} + e^{-\mathbb{W}_L t/2})}{4 \sinh \frac{\mathbb{W}_L t}{2}} - \frac{\mathbb{W}_L \cdot e^{-\mathbb{W}_L t/2}}{2 \sinh \frac{\mathbb{W}_L t}{2}} \\ &= \frac{\mathbb{W}_L \cdot (e^{\mathbb{W}_L t/2} - e^{-\mathbb{W}_L t/2})}{4 \sinh \frac{\mathbb{W}_L t}{2}} = \frac{\mathbb{W}_L}{2} \end{aligned}$$

Hence, (2.816a-b) become

$$\left[\frac{\hbar}{i} \nabla - \frac{e}{c} \mathbf{A}(\mathbf{x}) - M \frac{\mathbb{W}_L}{2} \cdot (\mathbf{x} - \mathbf{x}') \right] C = 0 \quad (2.817)$$

$$\left[\frac{\hbar}{i} \nabla' + \frac{e}{c} \mathbf{A}(\mathbf{x}') + M \frac{\mathbb{W}_L}{2} \cdot (\mathbf{x}' - \mathbf{x}) \right] C = 0 \quad (2.818)$$

Remembering that we are dealing with a constant field \mathbf{B} so that \mathbb{W}_L is also constant, the solution is

$$C(\mathbf{x}, \mathbf{x}') = C \exp \left\{ \frac{i}{\hbar} \int_{\mathbf{x}'}^{\mathbf{x}} d\xi \cdot \left[\frac{e}{c} \mathbf{A}(\xi) + \frac{1}{2} M \mathbb{W}_L \cdot (\xi - \mathbf{x}') \right] \right\} \quad (2.819)$$

where C is some constant since

$$\begin{aligned} \frac{\hbar}{i} \nabla C &= \left\{ \nabla \int_{\mathbf{x}'}^{\mathbf{x}} d\xi \cdot \left[\frac{e}{c} \mathbf{A}(\xi) + \frac{1}{2} M \mathbb{W}_L \cdot (\xi - \mathbf{x}') \right] \right\} C \\ &= \left[\frac{e}{c} \mathbf{A}(\xi) + \frac{1}{2} M \mathbb{W}_L \cdot (\xi - \mathbf{x}') \right]_{\xi=\mathbf{x}} C \end{aligned}$$

$$\begin{aligned}
&= \left[\frac{e}{c} \mathbf{A}(\mathbf{x}) + \frac{1}{2} M \mathbb{W}_L \cdot (\mathbf{x} - \mathbf{x}') \right] C \\
\frac{\hbar}{i} \nabla' C &= \left\{ \nabla' \int_{\mathbf{x}'}^{\mathbf{x}} d \boldsymbol{\xi} \cdot \left[\frac{e}{c} \mathbf{A}(\boldsymbol{\xi}) + \frac{1}{2} M \mathbb{W}_L \cdot (\boldsymbol{\xi} - \mathbf{x}') \right] \right\} C \\
&= \left\{ - \left[\frac{e}{c} \mathbf{A}(\boldsymbol{\xi}) + \frac{1}{2} M \mathbb{W}_L \cdot (\boldsymbol{\xi} - \mathbf{x}') \right] \Big|_{\boldsymbol{\xi}=\mathbf{x}'} - \int_{\mathbf{x}'}^{\mathbf{x}} d \xi_j \frac{1}{2} M \mathbb{W}_{Ljk} (\nabla' x_k') \right\} C \\
&= \left\{ - \frac{e}{c} \mathbf{A}(\mathbf{x}') - \frac{1}{2} [M(\mathbf{x} - \mathbf{x}')^T \cdot \mathbb{W}_L]^T \right\} C \\
&= \left\{ - \frac{e}{c} \mathbf{A}(\mathbf{x}') - \frac{1}{2} M \mathbb{W}_L^T \cdot (\mathbf{x} - \mathbf{x}') \right\} C \\
&= \left\{ - \frac{e}{c} \mathbf{A}(\mathbf{x}') - \frac{1}{2} M \mathbb{W}_L \cdot (\mathbf{x}' - \mathbf{x}) \right\} C
\end{aligned}$$

Let

$$\mathbf{A}'(\boldsymbol{\xi}) = \mathbf{A}(\boldsymbol{\xi}) + \frac{1}{2} \frac{c}{e} M \mathbb{W}_L \cdot (\boldsymbol{\xi} - \mathbf{x}') \quad (2.820a)$$

Using (2.797a),

$$\mathbb{W}_{Lij} = \frac{e}{Mc} \epsilon_{ijk} B_k$$

we have

$$A_i' = A_i + \frac{1}{2} \epsilon_{ijk} B_k (\xi_j - x_j') \quad (2.820b)$$

or

$$\mathbf{A}'(\boldsymbol{\xi}) = \mathbf{A}(\boldsymbol{\xi}) - \frac{1}{2} \mathbf{B} \times (\boldsymbol{\xi} - \mathbf{x}') \quad (2.820)$$

Since \mathbf{B} is constant, we have

$$\nabla_{\boldsymbol{\xi}} \times [\mathbf{B} \times (\boldsymbol{\xi} - \mathbf{x}')] = \mathbf{B} [\nabla_{\boldsymbol{\xi}} \cdot (\boldsymbol{\xi} - \mathbf{x}')] - (\mathbf{B} \cdot \nabla_{\boldsymbol{\xi}}) (\boldsymbol{\xi} - \mathbf{x}') = 3\mathbf{B} - \mathbf{B} = 2\mathbf{B}$$

so that

$$\nabla_{\boldsymbol{\xi}} \times \mathbf{A}'(\boldsymbol{\xi}) = \nabla_{\boldsymbol{\xi}} \times \mathbf{A}(\boldsymbol{\xi}) - \mathbf{B} = 0$$

Therefore, the line integral in (2.819) is path-independent. Choosing the straight line connecting \mathbf{x}' to \mathbf{x} as the integration path, we have

$$\begin{aligned}
&d \boldsymbol{\xi} \parallel (\boldsymbol{\xi} - \mathbf{x}') \\
\rightarrow d \boldsymbol{\xi} \cdot \mathbb{W}_L \cdot (\boldsymbol{\xi} - \mathbf{x}') &\propto d \boldsymbol{\xi} \cdot [\mathbf{B} \times (\boldsymbol{\xi} - \mathbf{x}')] = 0
\end{aligned}$$

Therefore, (2.819) reduces to

$$C(\mathbf{x}, \mathbf{x}') = C \exp \left\{ \frac{i}{\hbar} \frac{e}{c} \int_{\mathbf{x}'}^{\mathbf{x}} d \boldsymbol{\xi} \cdot \mathbf{A}(\boldsymbol{\xi}) \right\} \quad (2.821)$$

C is fixed by the I.C. [see (2.777)]

$$\lim_{t \rightarrow 0} \langle \mathbf{x} t | \mathbf{x}' 0 \rangle = \langle \mathbf{x} | \mathbf{x}' \rangle = \delta(\mathbf{x} - \mathbf{x}')$$

and take the value [see (2.778)]

$$C = \left(\frac{M}{i 2 \pi \hbar} \right)^{D/2}$$

Putting (2.813) & (2.821) together, we have

$$\begin{aligned}
 (\mathbf{x}t | \mathbf{x}'0) &= \left(\frac{M}{i2\pi\hbar t} \right)^{D/2} \left[\det \left(\frac{\sinh(\mathbb{W}_L t / 2)}{\mathbb{W}_L t / 2} \right) \right]^{-1/2} \exp \frac{i}{\hbar} \frac{e}{c} \int_{\mathbf{x}'}^{\mathbf{x}} d\xi \cdot \mathbf{A}(\xi) \\
 &\quad \times \exp \left[\frac{i}{\hbar} \frac{1}{2} M (\mathbf{x} - \mathbf{x}')^T \cdot \frac{1}{2} \mathbb{W}_L \cdot \coth \left(\frac{\mathbb{W}_L t}{2} \right) \cdot (\mathbf{x} - \mathbf{x}') \right]
 \end{aligned} \tag{2.822}$$

These expressions can be simplified by assuming $\mathbf{B} = B\hat{\mathbf{z}}$. (2.791a) & (2.795a) thus become

$$\mathbb{B} = B \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \mathbb{W}_L = \omega_L \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \tag{2.823}$$

$$\rightarrow \mathbb{W}_L^2 = -\omega_L^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\begin{aligned}
 \cos \frac{\mathbb{W}_L t}{2} &= \mathbb{I} + \sum_{n=1}^{\infty} \frac{(-1)^n}{(2n)!} \left(\frac{\mathbb{W}_L t}{2} \right)^{2n} = \mathbb{I} + \sum_{n=1}^{\infty} \frac{(-\mathbb{W}_L^2)^n}{(2n)!} \left(\frac{t}{2} \right)^{2n} \\
 &= \mathbb{I} + \sum_{n=1}^{\infty} \frac{1}{(2n)!} \left(\frac{\omega_L t}{2} \right)^{2n} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \\
 &= \mathbb{I} + \left[\cosh \left(\frac{\omega_L t}{2} \right) - 1 \right] \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \\
 &= \begin{pmatrix} \cosh \left(\frac{\omega_L t}{2} \right) & 0 & 0 \\ 0 & \cosh \left(\frac{\omega_L t}{2} \right) & 0 \\ 0 & 0 & 1 \end{pmatrix}
 \end{aligned} \tag{2.824}$$

$$\begin{aligned}
 \sinh \frac{\mathbb{W}_L t}{2} &= \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \left(\frac{\mathbb{W}_L t}{2} \right)^{2n+1} \\
 &= \frac{\mathbb{W}_L t}{2} + \frac{\mathbb{W}_L t}{2} \cdot \sum_{n=1}^{\infty} \frac{1}{(2n+1)!} \left(\frac{\mathbb{W}_L t}{2} \right)^{2n} \\
 &= \frac{\mathbb{W}_L t}{2} + \frac{\mathbb{W}_L t}{2} \cdot \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \sum_{n=1}^{\infty} \frac{(-1)^n}{(2n+1)!} \left(\frac{\omega_L t}{2} \right)^{2n} \\
 &= \frac{\mathbb{W}_L t}{2} + \frac{\mathbb{W}_L t}{2} \cdot \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \left[\frac{\sin(\omega_L t / 2)}{\omega_L t / 2} - 1 \right]
 \end{aligned}$$

Although \mathbb{W}_L^{-1} does not actually exist, it is legitimate to write

$$\begin{aligned}
 \frac{\sinh(\mathbb{W}_L t / 2)}{\mathbb{W}_L t / 2} &= \mathbb{I} + \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \left[\frac{\sin(\omega_L t / 2)}{\omega_L t / 2} - 1 \right] \\
 &= \begin{pmatrix} \frac{\sin(\omega_L t / 2)}{\omega_L t / 2} & 0 & 0 \\ 0 & \frac{\sin(\omega_L t / 2)}{\omega_L t / 2} & 0 \\ 0 & 0 & 1 \end{pmatrix}
 \end{aligned} \tag{2.825}$$

$$\rightarrow \det \frac{\sinh(\mathbb{W}_L t / 2)}{\mathbb{W}_L t / 2} = \left(\frac{\sin(\omega_L t / 2)}{\omega_L t / 2} \right)^2 \tag{2.826}$$

Alternatively, the above results can be obtained easily using the fact \mathbb{W}_L^2 is diagonal so that any function $F(\mathbb{W}_L^2)$ must be diagonal. The eigenvalues of \mathbb{W}_L are $\pm i\omega_L$ and 0, with respective eigenvectors

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm i \\ 0 \end{pmatrix} \quad \& \quad \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

Any function $f(W_L)$ that can be put into the form $F(W_L^2)$ can therefore be written as

$$f(W_L) = F(W_L^2) = \begin{pmatrix} f(\pm i \omega_L) & 0 & 0 \\ 0 & f(\mp i \omega_L) & 0 \\ 0 & 0 & f(0) \end{pmatrix} \quad (2.826a)$$

where the result is independent of the signs we choose since $(\pm i \omega_L)^2 = -\omega_L^2$.

Consider now the line integral

$$\mathcal{I} = \int_{\mathbf{x}'}^{\mathbf{x}} d\xi \cdot \mathbf{A}(\xi)$$

for the gauge with

$$\mathbf{A}(\mathbf{x}) = B x \hat{\mathbf{y}}$$

along the straight-line path

$$\xi = \mathbf{x}' + s(\mathbf{x} - \mathbf{x}') \quad s \in [0, 1] \quad (2.827)$$

Hence,

$$\begin{aligned} \mathcal{I} &= \int_0^1 ds (\mathbf{x} - \mathbf{x}') \cdot (B \xi_x \hat{\mathbf{y}}) \\ &= B(y - y') \int_0^1 ds [x' + s(x - x')] \\ &= B(y - y') \left[x' + \frac{1}{2}(x - x') \right] \\ &= \frac{1}{2} B(y - y')(x + x') \\ &= \frac{1}{2} B(xy - x'y') + \frac{1}{2} B(x'y - x'y') \end{aligned} \quad (2.828)$$

Thus, $\frac{e}{c} \mathcal{I}$ provides \mathcal{A}_{sf} in (2.668) as well as the last term of \mathcal{A}_{cl} in (2.667).

Using (2.826a), we have

$$\begin{aligned} \mathbb{W}_L \cdot \coth\left(\frac{\mathbb{W}_L t}{2}\right) &= \begin{pmatrix} (\pm i \omega_L) \coth\left(\frac{\pm i \omega_L t}{2}\right) & 0 & 0 \\ 0 & (\mp i \omega_L) \coth\left(\frac{\mp i \omega_L t}{2}\right) & 0 \\ 0 & 0 & \lim_{w \rightarrow 0} w \coth \frac{wt}{2} \end{pmatrix} \\ &= \begin{pmatrix} \omega_L \cot\left(\frac{\omega_L t}{2}\right) & 0 & 0 \\ 0 & \omega_L \cot\left(\frac{\omega_L t}{2}\right) & 0 \\ 0 & 0 & \frac{2}{t} \end{pmatrix} \end{aligned}$$

Hence, the other term in (2.822) that reads

$$\begin{aligned} &\frac{1}{2} M(\mathbf{x} - \mathbf{x}')^T \cdot \frac{1}{2} \mathbb{W}_L \cdot \coth\left(\frac{\mathbb{W}_L t}{2}\right) \cdot (\mathbf{x} - \mathbf{x}') \\ &= \frac{1}{2} M \left\{ \frac{1}{2} \omega_L \cot\left(\frac{\omega_L t}{2}\right) [(x - x')^2 + (y - y')^2] + \frac{(z - z')^2}{t} \right\} \end{aligned}$$

gives the rest of \mathcal{A}_{cl} in (2.667).

Therefore, (2.822) agrees with (2.666), as it should.