

3.4. Summing Spectral Representation of Green Function

Setting

$$t_2 = t_b - t \quad t_1 = t' - t_a \quad (3.140a)$$

we can write (3.70) as

$$\begin{aligned} G_{\omega^2}(t, t') &= \frac{2}{t_b - t_a} \sum_{n=1}^{\infty} (-)^{n+1} \frac{\sin v_n t_2 \sin v_n t_1}{v_n^2 - \omega^2} \\ &= \frac{2}{t_b - t_a} \sum_{n=1}^{\infty} (-)^{n+1} \frac{(e^{i v_n t_2} - e^{-i v_n t_2})(e^{i v_n t_1} - e^{-i v_n t_1})}{(2i)^2 (v_n^2 - \omega^2)} \\ &= \frac{1}{2} \frac{1}{t_b - t_a} \sum_{n=1}^{\infty} (-)^n \frac{1}{v_n^2 - \omega^2} (e^{-i v_n (t_2+t_1)} - e^{-i v_n (t_2-t_1)} + e^{i v_n (t_2+t_1)} - e^{i v_n (t_2-t_1)}) \\ &= \frac{1}{2} \frac{1}{t_b - t_a} \sum_{n=-\infty}^{\infty} (-)^n \frac{e^{-i v_n (t_2+t_1)} - e^{-i v_n (t_2-t_1)}}{v_n^2 - \omega^2} \end{aligned} \quad (3.140)$$

where we've used [see (3.64)]

$$v_n = \frac{n \pi}{t_b - t_a} = -v_{-n}$$

Setting [see (3.80) & (3.104)]

$$\omega_m = v_{2m} \quad \omega_m^f = v_{2m+1}$$

(3.140) becomes

$$G_{\omega^2}(t, t') = \frac{1}{2} \frac{1}{t_b - t_a} \sum_{m=-\infty}^{\infty} \left(\frac{e^{-i \omega_m (t_2+t_1)} - e^{-i \omega_m (t_2-t_1)}}{\omega_m^2 - \omega^2} - \frac{e^{-i \omega_m^f (t_2+t_1)} - e^{-i \omega_m^f (t_2-t_1)}}{\omega_m^{f2} - \omega^2} \right) \quad (3.141)$$

Using (3.97) & its fermionic counterpart, this becomes

$$G_{\omega^2}(t, t') = \frac{1}{2} \left[G_{\omega^2}^p(t_2 + t_1) - G_{\omega^2}^p(t_2 - t_1) - G_{\omega^2}^a(t_2 + t_1) + G_{\omega^2}^a(t_2 - t_1) \right] \quad (3.142)$$

With the help of (3.99), we have

$$\begin{aligned} G_{\omega^2}^p(t_2 + t_1) - G_{\omega^2}^p(t_2 - t_1) &= \\ &= \frac{1}{2 \omega} \left(\left(\cos \omega \left[t_2 + t_1 - \frac{1}{2} (t_b - t_a) \right] - \cos \omega \left[t_2 - t_1 - \frac{1}{2} (t_b - t_a) \right] \right) / \sin[\omega (t_b - t_a) / 2] \right) \\ &= \frac{\sin \omega t_1 \sin \omega \left[t_2 - \frac{1}{2} (t_b - t_a) \right]}{\omega \sin[\omega (t_b - t_a) / 2]} \end{aligned} \quad (3.143)$$

Using & (3.113), we get

$$\begin{aligned} G_{\omega^2}^a(t_2 + t_1) - G_{\omega^2}^a(t_2 - t_1) &= \\ &= \frac{1}{2 \omega} \left(\left(\sin \omega \left[t_2 + t_1 - \frac{1}{2} (t_b - t_a) \right] - \sin \omega \left[t_2 - t_1 - \frac{1}{2} (t_b - t_a) \right] \right) / \cos[\omega (t_b - t_a) / 2] \right) \\ &= - \frac{\sin \omega t_1 \cos \omega \left[t_2 - \frac{1}{2} (t_b - t_a) \right]}{\omega \cos[\omega (t_b - t_a) / 2]} \end{aligned} \quad (3.144)$$

(3.142) thus becomes

$$G_{\omega^2}(t, t') = \frac{\sin \omega t_1}{2 \omega} \left(\frac{\sin \omega \left[t_2 - \frac{1}{2} (t_b - t_a) \right]}{\sin[\omega (t_b - t_a) / 2]} + \frac{\cos \omega \left[t_2 - \frac{1}{2} (t_b - t_a) \right]}{\cos[\omega (t_b - t_a) / 2]} \right)$$

$$\begin{aligned}
 &= \frac{\sin \omega t_1}{2 \omega \sin[\omega(t_b - t_a)/2] \cos[\omega(t_b - t_a)/2]} \\
 &\quad \times \left(\sin \omega \left[t_2 - \frac{1}{2}(t_b - t_a) \right] \cos[\omega(t_b - t_a)/2] \right. \\
 &\quad \left. + \cos \omega \left[t_2 - \frac{1}{2}(t_b - t_a) \right] \sin[\omega(t_b - t_a)/2] \right) \\
 &= \frac{\sin \omega t_1 \sin \omega t_2}{\omega \sin \omega(t_b - t_a)} \\
 &= \frac{\sin \omega(t' - t_a) \sin \omega(t_b - t)}{\omega \sin \omega(t_b - t_a)}
 \end{aligned} \tag{3.145}$$

in agreement with (3.36) for $t > t'$.

In conjunction with $G_{\omega^2}(t, t') = 0$ for $t < t'$, (3.145) then gives the retarded Green function.

Alternatively, by switching the role of t & t' for $t < t'$, we get the causal Green function of (3.36)

$$\begin{aligned}
 G_{\omega^2}(t, t') &= (\theta(t - t') \sin \omega(t_b - t) \sin \omega(t' - t_a) + \theta(t' - t) \sin \omega(t_b - t') \sin \omega(t - t_a)) / (\omega \sin \omega(t_b - t_a)) \\
 &= \frac{\sin \omega(t_b - t_>) \sin \omega(t_< - t_a)}{\omega \sin \omega(t_b - t_a)}
 \end{aligned} \tag{3.145a}$$

For the case

$$t_a \rightarrow -\infty \quad t_b \rightarrow \infty \tag{3.146}$$

there are only two independent solutions, $e^{\pm i\omega t}$, to the homogeneous eq. Using the Wronskian construction [see §3.2.1], we get

$$\begin{aligned}
 G_{\omega^2}(t, t') &= -\frac{i}{2\omega} \begin{cases} e^{-i\omega(t-t')} & \text{for } t > t' \\ e^{i\omega(t-t')} & \text{for } t < t' \end{cases} \\
 &= -\frac{i}{2\omega} \left[\Theta(t-t') e^{-i\omega(t-t')} + \Theta(t'-t) e^{i\omega(t-t')} \right] \\
 &= -\frac{i}{2\omega} e^{-i\omega|t-t'|} \\
 \rightarrow \partial_t G_{\omega^2}(t, t') &= -\frac{i}{2\omega} \left[\delta(t-t') - i\omega \Theta(t-t') e^{-i\omega(t-t')} - \delta(t'-t) + i\omega \Theta(t'-t) e^{i\omega(t-t')} \right] \\
 &= \frac{1}{2} \left[-\Theta(t-t') e^{-i\omega(t-t')} + \Theta(t'-t) e^{i\omega(t-t')} \right] \\
 \partial_t^2 G_{\omega^2}(t, t') &= \frac{1}{2} \left[-\delta(t-t') + i\omega \Theta(t-t') e^{-i\omega(t-t')} - \delta(t'-t) + i\omega \Theta(t'-t) e^{i\omega(t-t')} \right] \\
 &= -\delta(t-t') - \omega^2 G_{\omega^2}(t, t')
 \end{aligned} \tag{3.147}$$

Hence,

$$(-\partial_t^2 - \omega^2) G_{\omega^2}(t, t') = \delta(t - t') \tag{3.148}$$

On the other hand, the periodic & antiperiodic solutions can be constructed, as in (3.106), from the infinite ranged Green function (3.147) as

$$G_{\omega^2}^p(t, t') = \sum_{n=-\infty}^{\infty} G_{\omega^2}[t - t' - (t_b - t_a)n] \tag{3.149}$$

$$G_{\omega^2}^a(t, t') = \sum_{n=-\infty}^{\infty} (-1)^n G_{\omega^2}[t - t' - (t_b - t_a)n] \tag{3.149a}$$

For completeness, we consider the spectral representation with respect to the basis functions [c.f. (3.63)],

$$x_0(t) = \frac{1}{\sqrt{t_b - t_a}} \quad x_n(t) = \sqrt{\frac{2}{t_b - t_a}} \cos v_n(t - t_a) \quad (3.150)$$

with

$$v_n = \frac{n \pi}{t_b - t_a} \quad n = 0, 1, 2, 3, \dots \quad (3.150a)$$

that satisfy the **Neumann B.C.**

$$\left. \partial_t x_n(t) \right|_{t=t_a} = 0 \quad \left. \partial_t x_n(t) \right|_{t=t_b} = 0 \quad (3.150b)$$

(3.69a) then gives

$$\begin{aligned} G_{\omega^2}^N(t, t') &= \frac{2}{t_b - t_a} \sum_{n=0}^{\infty} \frac{x_n(t) x_n(t')}{v_n^2 - \omega^2} \\ &= \frac{2}{t_b - t_a} \left[-\frac{1}{2\omega^2} + \sum_{n=1}^{\infty} \frac{\cos v_n(t - t_a) \cos v_n(t' - t_a)}{v_n^2 - \omega^2} \right] \end{aligned} \quad (3.151)$$

Thus, $G_{\omega^2}^N(t, t')$ satisfies the Neumann B.C.

$$\begin{aligned} \left. \partial_t G_{\omega^2}^N(t, t') \right|_{t=t_b} &= -\frac{2}{t_b - t_a} \sum_{n=1}^{\infty} \frac{v_n \sin v_n(t_b - t_a) \cos v_n(t' - t_a)}{v_n^2 - \omega^2} = 0 \quad [(3.150a) \text{ used. }] \\ \left. \partial_{t'} G_{\omega^2}^N(t, t') \right|_{t'=t_a} &= -\frac{2}{t_b - t_a} \sum_{n=1}^{\infty} \frac{v_n \cos v_n(t - t_a) \sin v_n(t_a - t_a)}{v_n^2 - \omega^2} = 0 \end{aligned} \quad (3.152)$$

Following the derivations (3.140-2), we have

$$\begin{aligned} \cos v_n(t_b - t) &= \cos v_n \left[t - t_a - (t_b - t_a) \right] \\ &= \cos \left[v_n(t - t_a) - n\pi \right] \quad [(3.64) \text{ used. }] \\ &= (-1)^n \cos v_n(t - t_a) \\ G_{\omega^2}^N(t, t') &= \frac{2}{t_b - t_a} \left[-\frac{1}{2\omega^2} + \sum_{n=1}^{\infty} (-1)^n \frac{\cos v_n t_2 \cos v_n t_1}{v_n^2 - \omega^2} \right] \\ &= \frac{2}{t_b - t_a} \left[-\frac{1}{2\omega^2} + \sum_{n=1}^{\infty} (-1)^n \frac{(e^{i v_n t_2} + e^{-i v_n t_2})(e^{i v_n t_1} + e^{-i v_n t_1})}{4(v_n^2 - \omega^2)} \right] \\ &= \frac{1}{2} \frac{1}{t_b - t_a} \left[-\frac{2}{\omega^2} + \sum_{n=1}^{\infty} (-1)^n \frac{1}{v_n^2 - \omega^2} (e^{-i v_n(t_2+t_1)} + e^{-i v_n(t_2-t_1)} + e^{i v_n(t_2+t_1)} + e^{i v_n(t_2-t_1)}) \right] \\ &= \frac{1}{2} \frac{1}{t_b - t_a} \sum_{n=-\infty}^{\infty} (-1)^n \frac{e^{-i v_n(t_2+t_1)} + e^{-i v_n(t_2-t_1)}}{v_n^2 - \omega^2} \quad (3.152a) \\ &= \frac{1}{2} \left[G_{\omega^2}^p(t_2 + t_1) + G_{\omega^2}^p(t_2 - t_1) - G_{\omega^2}^a(t_2 + t_1) - G_{\omega^2}^a(t_2 - t_1) \right] \quad (3.153) \end{aligned}$$

With the help of (3.99), we have [c.f. (3.143)]

$$\begin{aligned} G_{\omega^2}^p(t_2 + t_1) + G_{\omega^2}^p(t_2 - t_1) &= \\ &= -\frac{1}{2\omega} \left(\left(\cos \omega \left[t_2 + t_1 - \frac{1}{2}(t_b - t_a) \right] + \cos \omega \left[t_2 - t_1 - \frac{1}{2}(t_b - t_a) \right] \right) / \sin[\omega(t_b - t_a)/2] \right) \\ &= -\frac{\cos \omega t_1 \cos \omega \left[t_2 - \frac{1}{2}(t_b - t_a) \right]}{\omega \sin[\omega(t_b - t_a)/2]} \end{aligned} \quad (3.154)$$

Using & (3.113), we get [c.f. (3.144)]

$$\begin{aligned}
 G_{\omega^2}^a(t_2 + t_1) + G_{\omega^2}^a(t_2 - t_1) &= \\
 -\frac{1}{2\omega} \left(\left(\sin \omega \left[t_2 + t_1 - \frac{1}{2}(t_b - t_a) \right] + \sin \omega \left[t_2 - t_1 - \frac{1}{2}(t_b - t_a) \right] \right) / \cos[\omega(t_b - t_a)/2] \right) \\
 &= -\frac{\cos \omega t_1 \sin \omega \left[t_2 - \frac{1}{2}(t_b - t_a) \right]}{\omega \cos[\omega(t_b - t_a)/2]} \tag{3.155}
 \end{aligned}$$

(3.153) thus becomes [c.f. (3.145)]

$$\begin{aligned}
 G_{\omega^2}^N(t, t') &= -\frac{\cos \omega t_1}{2\omega} \left(\frac{\cos \omega \left[t_2 - \frac{1}{2}(t_b - t_a) \right]}{\sin[\omega(t_b - t_a)/2]} - \frac{\sin \omega \left[t_2 - \frac{1}{2}(t_b - t_a) \right]}{\cos[\omega(t_b - t_a)/2]} \right) \\
 &= -\frac{\cos \omega t_1 \cos \omega t_2}{\omega \sin[\omega(t_b - t_a)]} \\
 &= -\frac{\cos \omega(t' - t_a) \cos \omega(t_b - t)}{\omega \sin[\omega(t_b - t_a)]} \tag{3.156a}
 \end{aligned}$$

which is easily verified to satisfy the Neumann B.C. (3.152). As with (3.145), this is valid for $t > t'$.

The causal Green function is therefore [see (3.145a)]

$$G_{\omega^2}^N(t, t') = -\frac{\cos \omega(t_b - t_>) \cos \omega(t_< - t_a)}{\omega \sin[\omega(t_b - t_a)]} \tag{3.156}$$

The power expansion for $\omega \rightarrow 0$ can be obtained from the following *Mathematica* code

$$\begin{aligned}
 &-\frac{\text{Cos}[\omega(t_p - t_a)] \text{Cos}[\omega(t_b - t)]}{\omega \text{Sin}[\omega(t_b - t_a)]} + O[\omega]^2 \\
 &\frac{1}{(t_a - t_b) \omega^2} + \left(\frac{1}{6} (t_a - t_b) - \frac{-\frac{1}{2}(t_p - t_a)^2 - \frac{1}{2}(-t + t_b)^2}{-t_a + t_b} \right) + O[\omega]^2
 \end{aligned}$$

Thus,

$$G_{\omega^2}^N(t, t') = -\frac{1}{\omega^2(t_b - t_a)} - \frac{1}{6}(t_b - t_a) - \frac{1}{2(t_b - t_a)}[(t' - t_a)^2 + (t_b - t)^2] + O[\omega^2] \tag{3.157}$$