

3.5. Wronski Construction for Periodic and Antiperiodic Green Functions

The Wronski Construction in §3.2.1 for

$$[-\partial_t^2 - \Omega^2(t)] G_{\Omega^2}(t, t') = \delta(t - t') \quad (3.158)$$

can be carried over to

$$[-\partial_t^2 - \Omega^2(t)] G_{\Omega^2}^{p,a}(t, t') = \delta^{p,a}(t - t') \quad (3.158a)$$

where

$$G_{\Omega^2}^{p,a}(t, t') = \sum_{n=-\infty}^{\infty} (\pm)^n G_{\Omega^2}(t - t' - n\beta\hbar) \quad (3.159a)$$

$$\delta^{p,a}(t - t') = \sum_{n=-\infty}^{\infty} (\pm)^n \delta(t - t' - n\beta\hbar) \quad (3.159)$$

with the B.C.

$$G_{\Omega^2}^{p,a}(t, t') = \pm G_{\Omega^2}^{p,a}(t + t_b - t_a, t') \quad (3.159b)$$

$$\partial_t G_{\Omega^2}^{p,a}(t, t') = \pm \partial_t G_{\Omega^2}^{p,a}(t + t_b - t_a, t') \quad (3.159c)$$

As in (3.53), we set

$$G_{\Omega^2}^{p,a}(t, t') = \overline{\Theta}(t - t') \Delta(t, t') + a(t') \xi(t) + b(t') \eta(t) \quad (3.159)$$

where ξ & η are independent solutions of the homogeneous solution and

$$[-\partial_t^2 - \Omega^2(t)] \Delta(t, t') = 0 \quad (3.161)$$

with I.C. [see (3.45-6)]

$$\Delta(t, t) = 0 \quad \partial_t \Delta(t, t') \Big|_{t=t'} = -1 \quad (3.161a)$$

Imposing the B.C. (3.159b) on (3.159) with

$$t = t_a \quad \& \quad t' \in [t_a, t_b)$$

gives

$$G_{\Omega^2}^{p,a}(t_a, t') = a(t') \xi(t_a) + b(t') \eta(t_a) \\ = \pm G_{\Omega^2}^{p,a}(t_b, t') = \pm [\Delta(t_b, t') + a(t') \xi(t_b) + b(t') \eta(t_b)]$$

$$\rightarrow [\xi(t_b) \mp \xi(t_a)] a(t') + [\eta(t_b) \mp \eta(t_a)] b(t') = -\Delta(t_b, t') \quad (3.162)$$

For the B.C. (3.159c),

$$\partial_t G_{\Omega^2}^{p,a}(t, t') \Big|_{t=t_a} = \delta(t_a - t') \Delta(t_a, t') + \overline{\Theta}(t_a - t') \partial_{t_a} \Delta(t_a, t') \\ + a(t') \dot{\xi}(t_a) + b(t') \dot{\eta}(t_a)$$

$$= a(t') \dot{\xi}(t_a) + b(t') \dot{\eta}(t_a)$$

$$= \pm \partial_t G_{\Omega^2}^{p,a}(t, t') \Big|_{t=t_b} = \pm [\delta(t_b - t') \Delta(t_b, t') + \partial_{t_b} \Delta(t_b, t') \\ + a(t') \dot{\xi}(t_b) + b(t') \dot{\eta}(t_b)]$$

$$= \pm [\partial_{t_b} \Delta(t_b, t') + a(t') \dot{\xi}(t_b) + b(t') \dot{\eta}(t_b)]$$

$$\rightarrow [\dot{\xi}(t_b) \mp \dot{\xi}(t_a)] a(t') + [\dot{\eta}(t_b) \mp \dot{\eta}(t_a)] b(t') = -\partial_{t_b} \Delta(t_b, t') \quad (3.162a)$$

(3.162-a) can be written in matrix form as

$$\begin{pmatrix} \xi(t_b) \mp \xi(t_a) & \eta(t_b) \mp \eta(t_a) \\ \dot{\xi}(t_b) \mp \dot{\xi}(t_a) & \dot{\eta}(t_b) \mp \dot{\eta}(t_a) \end{pmatrix} \begin{pmatrix} a(t') \\ b(t') \end{pmatrix} = \begin{pmatrix} -\Delta(t_b, t') \\ -\partial_{t_b} \Delta(t_b, t') \end{pmatrix} \quad (3.163a)$$

Let

$$\overline{\Lambda}^{p,a}(t_a, t_b) = \begin{pmatrix} \xi(t_b) \mp \xi(t_a) & \eta(t_b) \mp \eta(t_a) \\ \dot{\xi}(t_b) \mp \dot{\xi}(t_a) & \dot{\eta}(t_b) \mp \dot{\eta}(t_a) \end{pmatrix} \quad (3.163)$$

then the condition for (3.163a) to have a unique solution is

$$\det \bar{\Lambda}^{p,a}(t_a, t_b) \neq 0 \quad (3.164a)$$

whereupon

$$a(t') = \frac{-[\dot{\eta}(t_b) \mp \dot{\eta}(t_a)] \Delta(t_b, t') + [\eta(t_b) \mp \eta(t_a)] \partial_{t_b} \Delta(t_b, t')}{\det \bar{\Lambda}^{p,a}(t_a, t_b)} \quad (3.163b)$$

$$b(t') = \frac{[\dot{\xi}(t_b) \mp \dot{\xi}(t_a)] \Delta(t_b, t') - [\xi(t_b) \mp \xi(t_a)] \partial_{t_b} \Delta(t_b, t')}{\det \bar{\Lambda}^{p,a}(t_a, t_b)} \quad (3.163c)$$

From (3.163), we have

$$\begin{aligned} \det \bar{\Lambda}^{p,a}(t_a, t_b) &= [\xi(t_b) \mp \xi(t_a)] [\dot{\eta}(t_b) \mp \dot{\eta}(t_a)] \\ &\quad - [\dot{\xi}(t_b) \mp \dot{\xi}(t_a)] [\eta(t_b) \mp \eta(t_a)] \\ &= \xi(t_b) \dot{\eta}(t_b) - \dot{\xi}(t_b) \eta(t_b) + \xi(t_a) \dot{\eta}(t_a) - \dot{\xi}(t_a) \eta(t_a) \\ &\quad \mp [\xi(t_b) \dot{\eta}(t_a) + \xi(t_a) \dot{\eta}(t_b) - \dot{\xi}(t_b) \eta(t_a) - \dot{\xi}(t_a) \eta(t_b)] \end{aligned} \quad (3.164b)$$

From (3.48a), we see that

$$W = \begin{vmatrix} \xi & \eta \\ \dot{\xi} & \dot{\eta} \end{vmatrix} = \xi \dot{\eta} - \dot{\xi} \eta \quad (3.164c)$$

is independent of time. By (3.49), we have

$$\Delta(t, t') = \frac{1}{W} [\xi(t) \eta(t') - \xi(t') \eta(t)] \quad (3.164d)$$

$$\rightarrow \partial_t \Delta(t, t') = \frac{1}{W} [\dot{\xi}(t) \eta(t') - \xi(t') \dot{\eta}(t)] \quad (3.164e)$$

$$\partial_{t'} \Delta(t, t') = \frac{1}{W} [\xi(t) \dot{\eta}(t') - \dot{\xi}(t') \eta(t)] \quad (3.164f)$$

(3.164b) thus simplifies to

$$\det \bar{\Lambda}^{p,a}(t_a, t_b) = 2 W \mp W [-\partial_{t_b} \Delta(t_b, t_a) + \partial_{t_a} \Delta(t_b, t_a)] \quad (3.164g)$$

Setting

$$\begin{aligned} \bar{\Delta}^{p,a}(t_a, t_b) &= 2 \pm [\partial_{t_b} \Delta(t_b, t_a) - \partial_{t_a} \Delta(t_b, t_a)] \\ &= 2 \pm [\partial_{t_b} \Delta(t_b, t_a) + \partial_{t_a} \Delta(t_a, t_b)] \quad [\Delta(t, t') = -\Delta(t', t)] \end{aligned} \quad (3.165)$$

(3.164g) & hence (3.164a) become

$$\det \bar{\Lambda}^{p,a}(t_a, t_b) = W \bar{\Delta}^{p,a}(t_a, t_b) \neq 0 \quad (3.164)$$

We now turn to the evaluation of (3.159). To begin, we write it as

$$G_{\Omega^2}^{p,a}(t, t') = \bar{\Theta}(t - t') \Delta(t, t') + \frac{\mathcal{I}}{W \bar{\Delta}^{p,a}(t_a, t_b)} \quad (a)$$

where [see (3.163b-c)]

$$\begin{aligned} \mathcal{I} &= [a(t') \xi(t) + b(t') \eta(t)] \det \bar{\Lambda}^{p,a}(t_a, t_b) \\ &= A \Delta(t_b, t') + B \partial_{t_b} \Delta(t_b, t') \end{aligned} \quad (b)$$

where

$$\begin{aligned} A &= -[\dot{\eta}(t_b) \mp \dot{\eta}(t_a)] \xi(t) + [\dot{\xi}(t_b) \mp \dot{\xi}(t_a)] \eta(t) \\ &= W \left[\partial_{t_b} \Delta(t_b, t) \pm \partial_{t_a} \Delta(t, t_a) \right] \end{aligned} \quad (c)$$

$$\begin{aligned} B &= [\eta(t_b) \mp \eta(t_a)] \xi(t) - [\xi(t_b) \mp \xi(t_a)] \eta(t) \\ &= W \left[-\Delta(t_b, t) \mp \Delta(t, t_a) \right] \end{aligned} \quad (d)$$

$$\rightarrow \mathcal{I} = W \left\{ \left[\partial_{t_b} \Delta(t_b, t) \pm \partial_{t_a} \Delta(t, t_a) \right] \Delta(t_b, t') + \left[-\Delta(t_b, t) \mp \Delta(t, t_a) \right] \partial_{t_b} \Delta(t_b, t') \right\}$$

Using [see (3.51-2)]

$$\partial_{t_b} \Delta(t_b, t') = \frac{\Delta(t_b, t') \partial_{t_b} \Delta(t_b, t_a) - \Delta(t', t_a)}{\Delta(t_b, t_a)}$$

$$\partial_{t_a} \Delta(t, t_a) = \frac{\partial_{t_a} \Delta(t_b, t_a) \Delta(t, t_a) + \Delta(t_b, t)}{\Delta(t_b, t_a)}$$

we have

$$\begin{aligned} \mathcal{I} &= \frac{W}{\Delta(t_b, t_a)} \left\{ \left[\Delta(t_b, t) \partial_{t_b} \Delta(t_b, t_a) - \Delta(t, t_a) \pm \partial_{t_a} \Delta(t_b, t_a) \Delta(t, t_a) \pm \Delta(t_b, t) \right] \Delta(t_b, t') \right. \\ &\quad \left. + \left[-\Delta(t_b, t) \mp \Delta(t, t_a) \right] \left[\Delta(t_b, t') \partial_{t_b} \Delta(t_b, t_a) - \Delta(t', t_a) \right] \right\} \\ &= \frac{W}{\Delta(t_b, t_a)} \left\{ \left[-\Delta(t, t_a) \pm \Delta(t_b, t) \right] \Delta(t_b, t') - \left[-\Delta(t_b, t) \mp \Delta(t, t_a) \right] \Delta(t', t_a) \right. \\ &\quad \left. \pm \Delta(t, t_a) \Delta(t_b, t') \left[\partial_{t_a} \Delta(t_b, t_a) - \partial_{t_b} \Delta(t_b, t_a) \right] \right\} \\ &= \frac{W}{\Delta(t_b, t_a)} \left\{ \left[\Delta(t, t_a) \pm \Delta(t_b, t) \right] \left[\Delta(t_b, t') \pm \Delta(t', t_a) \right] \right. \\ &\quad \left. - 2 \Delta(t, t_a) \Delta(t_b, t') \pm \Delta(t, t_a) \Delta(t_b, t') \left[\partial_{t_a} \Delta(t_b, t_a) - \partial_{t_b} \Delta(t_b, t_a) \right] \right\} \\ &= \frac{W}{\Delta(t_b, t_a)} \left\{ \left[\Delta(t, t_a) \pm \Delta(t_b, t) \right] \left[\Delta(t_b, t') \pm \Delta(t', t_a) \right] \quad [(3.165) \text{ used.}] \right. \\ &\quad \left. - \Delta(t, t_a) \Delta(t_b, t') \overline{\Delta}^{p,a}(t_a, t_b) \right\} \end{aligned}$$

(a) thus becomes

$$\begin{aligned} G_{\Omega^2}^{p,a}(t, t') &= \overline{\Theta}(t-t') \Delta(t, t') - \frac{\Delta(t, t_a) \Delta(t_b, t')}{\Delta(t_b, t_a)} + \frac{\left[\Delta(t, t_a) \pm \Delta(t_b, t) \right] \left[\Delta(t_b, t') \pm \Delta(t', t_a) \right]}{\Delta(t_b, t_a) \overline{\Delta}^{p,a}(t_a, t_b)} \\ &= G_{\Omega^2}(t, t') + \frac{\left[\Delta(t, t_a) \pm \Delta(t_b, t) \right] \left[\Delta(t_b, t') \pm \Delta(t', t_a) \right]}{\Delta(t_b, t_a) \overline{\Delta}^{p,a}(t_a, t_b)} \end{aligned} \quad (3.166)$$

where the equation for $G_{\Omega^2}(t, t')$ just below (3.58a) in §3.2 was used.