

I.4. Grassmann Algebras, Differential Forms

Grassmann Algebra

A Grassmann (exterior) algebra \mathfrak{U} over \mathbb{R} or \mathbb{C} is an associative algebra constructed from $\{1, \theta_1, \dots, \theta_n\}$, where the generators θ_i are anti-symmetric, i.e.,

$$\theta_i \theta_j = -\theta_j \theta_i \quad \forall i, j \quad (1.31)$$

$$\rightarrow \theta_i^2 = -\theta_i^2 = 0$$

Thus,

1. All elements of \mathfrak{U} are polynomials of at most 1st degree in all θ_i .

2. The basis of \mathfrak{U} is the set

$$\{ \theta_1^{j_1} \dots \theta_n^{j_n} \mid j_k \in \{0, 1\} \}$$

where $\theta_j^0 = 1$ is understood.

3. Dimension of \mathfrak{U} is therefore 2^n .

For example, $n = 2$ gives $\{1, \theta_1, \theta_2\}$ so that the basis is

$$\{ \theta_1^0 \theta_2^0, \theta_1^1 \theta_2^0, \theta_1^0 \theta_2^1, \theta_1^1 \theta_2^1 \} = \{ 1, \theta_1, \theta_2, \theta_1 \theta_2 \}$$

The dimension of \mathfrak{U} is indeed $4 = 2^2$.

The degree of a monomial $\theta_1^{j_1} \dots \theta_m^{j_m}$ is $\sum_{k=1}^m j_k$.

Monomials with the highest degree are therefore multiples of $\theta_1 \dots \theta_n$.

Polynomials formed by linear combinations of monomials of the same degree are called homogeneous.

Obviously, all homogeneous polynomials of the same degree form a linear vector space. The vector space \mathfrak{U} is therefore partitioned into subspaces of different degrees (or grades). In this sense, \mathfrak{U} is a graded algebra.

Note that an element of \mathfrak{U} is invertible if and only if it contains the basis 1.

For example, element $1 + \theta$ is invertible.

Its inverse is $1 - \theta$ since

$$(1 + \theta)(1 - \theta) = 1 = (1 - \theta)(1 + \theta)$$

Grassmannian Parity

Parity operator P is defined by

$$P(\theta_i) = -\theta_i \quad \forall i \quad (1.32)$$

$$\rightarrow P^2(\theta_i) = P(-\theta_i) = \theta_i$$

$$\therefore P^2 = 1 \quad (\text{idempotent})$$

$$\begin{aligned} P(\theta_{i_1} \dots \theta_{i_p}) &= P(\theta_{i_1}) \dots P(\theta_{i_p}) \\ &= (-\theta_{i_1}) \dots (-\theta_{i_p}) \\ &= (-1)^p \theta_{i_1} \dots \theta_{i_p} \end{aligned} \quad (1.33)$$

i.e., monomials of even(odd) degree in θ_i has even(odd) parity.

Let $A^{(m)}$ & $B^{(n)}$ be homogeneous polynomials of degrees m & n , respectively, then

$$A^{(m)} B^{(n)} = (-1)^{mn} B^{(n)} A^{(m)}$$

$$= \begin{cases} -B^{(n)} A^{(m)} & \text{if both } n \text{ \& } m \text{ are odd} \\ B^{(n)} A^{(m)} & \text{otherwise} \end{cases}$$

For example,

$$\begin{aligned} (\theta_1 \theta_2) (\theta_3 \theta_4 \theta_5) &= (-)^3 \theta_1 (\theta_3 \theta_4 \theta_5) \theta_2 \\ &= (-)^6 (\theta_3 \theta_4 \theta_5) (\theta_1 \theta_2) \\ &= (\theta_3 \theta_4 \theta_5) (\theta_1 \theta_2) \end{aligned}$$

Let

$$P(\mathfrak{U}^\pm) = \pm \mathfrak{U}^\pm \tag{1.34}$$

→ \mathfrak{U}^+ is a (commuting) sub-algebra of \mathfrak{U} & is spanned by $\theta_{i_1} \dots \theta_{i_p}$ of even p .
 \mathfrak{U}^- does not form an algebra since 1 is not in it.

Since every element of the basis $\{ \theta_1^{j_1} \dots \theta_n^{j_n} \mid j_k \in \{0, 1\} \}$ must belong either to \mathfrak{U}^+ or \mathfrak{U}^- , any $A \in \mathfrak{U}$ can be written as

$$A = A^+ + A^- \quad \text{with} \quad A^\pm \in \mathfrak{U}^\pm$$

This means any element of \mathfrak{U}^+ commutes with every element of \mathfrak{U} .
 Elements of \mathfrak{U}^- anti-commute with each other but commute with elements of \mathfrak{U}^+ .

Differential Forms

Let $\Omega_{\mu_1 \dots \mu_m}(x)$ be a totally antisymmetric tensor. Then

$$\Omega = \Omega_{\mu_1 \dots \mu_m}(x) \theta^{\mu_1} \dots \theta^{\mu_m} \tag{1.35}$$

is an m -form (with $\theta^\mu \theta^\nu$ taking the place of $d x^\mu \wedge d x^\nu$).

For \mathfrak{U} of n generators, any m -form vanishes if $m > n$.

The differential operator is defined by

$$d \equiv \theta^\mu \frac{\partial}{\partial x^\mu} \tag{1.36}$$

Since $\theta^\mu \theta^\nu = -\theta^\nu \theta^\mu$ & $\frac{\partial}{\partial x^\mu} \frac{\partial}{\partial x^\nu} = \frac{\partial}{\partial x^\nu} \frac{\partial}{\partial x^\mu}$, we have

$$\begin{aligned} d^2 &= \theta^\mu \frac{\partial}{\partial x^\mu} \theta^\nu \frac{\partial}{\partial x^\nu} \\ &= \theta^\mu \theta^\nu \frac{\partial}{\partial x^\mu} \frac{\partial}{\partial x^\nu} = -\theta^\nu \theta^\mu \frac{\partial}{\partial x^\mu} \frac{\partial}{\partial x^\nu} \\ &= \theta^\nu \theta^\mu \frac{\partial}{\partial x^\nu} \frac{\partial}{\partial x^\mu} = \theta^\nu \theta^\mu \frac{\partial}{\partial x^\mu} \frac{\partial}{\partial x^\nu} \\ &= 0 \end{aligned} \tag{1.37}$$

In differential geometry,

$$\begin{aligned} d \Omega = 0 &\Leftrightarrow \Omega \text{ is closed} \\ \Omega = d \Omega' &\Leftrightarrow \Omega \text{ is exact} \end{aligned}$$

Since

$$d \Omega = d^2 \Omega' = 0$$

∴ All exact forms are closed.