

3.2. Schrodinger Field

Schrodinger eq
$$i \hbar \frac{\partial \phi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \phi + V \phi$$

Lagrangian

$$\mathcal{L} = i \hbar \phi^* \partial_t \phi - \frac{\hbar^2}{2m} \nabla \phi^* \cdot \nabla \phi - V \phi^* \phi$$

$$\frac{\partial \mathcal{L}}{\partial (\partial_t \phi^*)} = 0 \qquad \frac{\partial \mathcal{L}}{\partial (\partial_t \phi)} = 0$$

$$\frac{\partial \mathcal{L}}{\partial (\nabla \phi^*)} = -\frac{\hbar^2}{2m} \nabla \phi \qquad \nabla \cdot \frac{\partial \mathcal{L}}{\partial (\nabla \phi)} = -\frac{\hbar^2}{2m} \nabla^2 \phi$$

$$\frac{\partial \mathcal{L}}{\partial \phi^*} = i \hbar \partial_t \phi - V \phi$$

$$\rightarrow -\frac{\hbar^2}{2m} \nabla^2 \phi - i \hbar \partial_t \phi + V \phi = 0$$

$$\frac{\partial \mathcal{L}}{\partial (\partial_t \phi)} = i \hbar \phi^* \qquad \partial_t \frac{\partial \mathcal{L}}{\partial (\partial_t \phi)} = i \hbar \partial_t \phi^*$$

$$\frac{\partial \mathcal{L}}{\partial (\nabla \phi)} = -\frac{\hbar^2}{2m} \nabla \phi^* \qquad \nabla \cdot \frac{\partial \mathcal{L}}{\partial (\nabla \phi)} = -\frac{\hbar^2}{2m} \nabla^2 \phi^*$$

$$\frac{\partial \mathcal{L}}{\partial \phi} = -V \phi^*$$

$$\rightarrow i \hbar \partial_t \phi^* - \frac{\hbar^2}{2m} \nabla^2 \phi^* + V \phi^* = 0$$

$$\pi = \frac{\partial \mathcal{L}}{\partial (\partial_t \phi)} = i \hbar \phi^*$$

$$\pi^* = \frac{\partial \mathcal{L}}{\partial (\partial_t \phi^*)} = 0 \neq (\pi)^*$$

$$\mathcal{H} = i \hbar \phi^* \partial_t \phi - \mathcal{L} = \frac{\hbar^2}{2m} \nabla \phi^* \cdot \nabla \phi + V \phi^* \phi$$

2nd Quantization via equal-time commutation relations :

$$[\phi(t, \mathbf{r}), \pi(t, \mathbf{r}')]_{\mp} = i \hbar \delta(\mathbf{r} - \mathbf{r}') \qquad \text{for } \begin{matrix} \text{bosons} \\ \text{fermions} \end{matrix}$$

$$\rightarrow [\phi(t, \mathbf{r}), \phi^+(t, \mathbf{r}')]_{\mp} = \delta(\mathbf{r} - \mathbf{r}')$$

Hamiltonian

Let

$$H \varphi_j(\mathbf{r}) = \epsilon_j \varphi_j(\mathbf{r}) \qquad \text{with } \langle \varphi_j | \varphi_k \rangle = \delta_{jk}$$

$$\rightarrow \phi(t, \mathbf{r}) = \sum_j a_j \varphi_j(\mathbf{r}) e^{-i\epsilon_j t / \hbar} \qquad \phi^*(t, \mathbf{r}) = \sum_j a_j^* \varphi_j^*(\mathbf{r}) e^{i\epsilon_j t / \hbar}$$

$$\begin{aligned}
\nabla \phi &= \sum_j a_j (\nabla \varphi_j) e^{-i \epsilon_j t / \hbar} & \nabla \phi^\dagger &= \sum_j a_j^\dagger (\nabla \varphi_j^*) e^{i \epsilon_j t / \hbar} \\
\mathcal{H} &= \sum_{j,k} \left[\frac{\hbar^2}{2m} (\nabla \varphi_j^* \cdot \nabla \varphi_k) + V \varphi_j^* \varphi_k \right] a_j^\dagger a_k e^{-i(\epsilon_k - \epsilon_j)t / \hbar} \\
H &= \int d^3 r \mathcal{H} \\
&= \sum_{j,k} a_j^\dagger a_k e^{-i(\epsilon_k - \epsilon_j)t / \hbar} \int d^3 r \left[\frac{\hbar^2}{2m} (\nabla \varphi_j^* \cdot \nabla \varphi_k) + V \varphi_j^* \varphi_k \right] \\
&= \sum_{j,k} a_j^\dagger a_k e^{-i(\epsilon_k - \epsilon_j)t / \hbar} \int d^3 r \varphi_j^* \left(-\frac{\hbar^2}{2m} \nabla^2 \varphi_k + V \varphi_k \right) \\
&= \sum_{j,k} a_j^\dagger a_k e^{-i(\epsilon_k - \epsilon_j)t / \hbar} \int d^3 r \varphi_j^* H \varphi_k \\
&= \sum_{j,k} a_j^\dagger a_k e^{-i(\epsilon_k - \epsilon_j)t / \hbar} \epsilon_j \delta_{jk} \\
&= \sum_j \epsilon_j a_j^\dagger a_j
\end{aligned}$$

Momentum Expansion

See FourierTransforms.pdf .

For interacting particles,

$$\begin{aligned}
\mathcal{L} &= i \hbar \phi^\dagger(x) \partial_t \phi(x) - \frac{\hbar^2}{2m} \nabla \phi^\dagger(x) \cdot \nabla \phi(x) - \\
&\quad V(x) \phi^\dagger(x) \phi(x) - \frac{1}{2} \int d^3 r' \phi^\dagger(x) \phi^\dagger(x') U(\mathbf{r} - \mathbf{r}') \phi(x') \phi(x) \Big|_{t=t'}
\end{aligned}$$

where $x = (t, \mathbf{r})$ $x' = (t', \mathbf{r}')$

$$\begin{aligned}
\rightarrow \mathcal{H} &= \frac{\hbar^2}{2m} \nabla \phi^\dagger(x) \cdot \nabla \phi(x) + V(\mathbf{r}) \phi^\dagger(x) \phi(x) \\
&\quad + \frac{1}{2} \int d^3 r' \phi^\dagger(x) \phi^\dagger(x') U(\mathbf{r} - \mathbf{r}') \phi(x') \phi(x) \Big|_{t=t'}
\end{aligned}$$

Momentum (plane wave) expansion:

$$\begin{aligned}
\phi(x) &= \int \frac{d^3 k}{(2\pi)^{3/2}} a_k(t) e^{i\mathbf{k} \cdot \mathbf{r}} \\
\int d^3 r e^{i\mathbf{k} \cdot \mathbf{r}} &= (2\pi)^3 \delta(\mathbf{k}) \\
\rightarrow a_k(t) &= \int \frac{d^3 r}{(2\pi)^{3/2}} \phi(x) e^{-i\mathbf{k} \cdot \mathbf{r}} \\
[\phi(t, \mathbf{r}), \phi^\dagger(t, \mathbf{r}')]_{\mp} &= \delta(\mathbf{r} - \mathbf{r}') \\
\rightarrow \int \frac{d^3 k}{(2\pi)^{3/2}} \int \frac{d^3 k'}{(2\pi)^{3/2}} [a_k(t), a_{k'}^\dagger(t)]_{\mp} e^{i\mathbf{k} \cdot \mathbf{r} - i\mathbf{k}' \cdot \mathbf{r}'} &= \int \frac{d^3 k}{(2\pi)^3} e^{i\mathbf{k} \cdot (\mathbf{r} - \mathbf{r}')} \\
\rightarrow [a_k(t), a_{k'}^\dagger(t)]_{\mp} &= \delta_{\mathbf{k}, \mathbf{k}'}
\end{aligned}$$

Similarly,

$$[\phi(t, \mathbf{r}), \phi(t, \mathbf{r}')]_{\mp} = [\phi^+(t, \mathbf{r}), \phi^+(t, \mathbf{r}')]_{\mp} = 0$$

$$\rightarrow [a_{\mathbf{k}}(t), a_{\mathbf{k}'}(t)]_{\mp} = [a_{\mathbf{k}}^+(t), a_{\mathbf{k}'}^+(t)]_{\mp} = 0$$

Total Momentum & Hamiltonian

$$\begin{aligned} \mathbf{p} &= \int d^3 r \phi^+(x) \frac{\hbar}{i} \nabla \phi(x) \\ &= \int d^3 r \int \frac{d^3 k}{(2\pi)^{3/2}} \int \frac{d^3 k'}{(2\pi)^{3/2}} a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}(t) e^{-i\mathbf{k}\cdot\mathbf{r}} \frac{\hbar}{i} \nabla e^{i\mathbf{k}'\cdot\mathbf{r}} \\ &= \int d^3 r \int \frac{d^3 k}{(2\pi)^{3/2}} \int \frac{d^3 k'}{(2\pi)^{3/2}} a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}(t) \hbar \mathbf{k}' e^{i(\mathbf{k}'-\mathbf{k})\cdot\mathbf{r}} \\ &= \int d^3 k \int d^3 k' a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}(t) \hbar \mathbf{k}' \delta(\mathbf{k}'-\mathbf{k}) \\ &= \int d^3 k a_{\mathbf{k}}^+(t) a_{\mathbf{k}}(t) \hbar \mathbf{k} \end{aligned}$$

Similarly

$$\begin{aligned} \int d^3 r \frac{\hbar^2}{2m} \nabla \phi^+(x) \cdot \nabla \phi(x) &= -\frac{\hbar^2}{2m} \int d^3 x \phi^+(x) \nabla^2 \phi(x) \\ &= \int d^3 k a_{\mathbf{k}}^+(t) a_{\mathbf{k}}(t) \frac{\hbar^2 \mathbf{k}^2}{2m} \\ \int d^3 r V(x) \phi^+(x) \phi(x) &= \int d^3 r \int \frac{d^3 k}{(2\pi)^{3/2}} \int \frac{d^3 k'}{(2\pi)^{3/2}} a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}(t) V(\mathbf{r}) e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}} \\ &= \int d^3 k \int \frac{d^3 k'}{(2\pi)^3} a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}(t) V(\mathbf{k}-\mathbf{k}') \quad f(\mathbf{k}) = \int d^3 r f(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{r}} \\ &= \int d^3 k \int \frac{d^3 \mathbf{q}}{(2\pi)^3} a_{\mathbf{k}}^+(t) a_{\mathbf{k}-\mathbf{q}}(t) V(\mathbf{q}) \quad \mathbf{q} = \mathbf{k} - \mathbf{k}' \\ \frac{1}{2} \int d^3 r \int d^3 r' \phi^+(x) \phi^+(x') U(\mathbf{r}-\mathbf{r}') \phi(x') \phi(x) \Big|_{t=t'} & \\ &= \frac{1}{2} \int d^3 r \int d^3 r' \int \frac{d^3 k_1}{(2\pi)^{3/2}} \int \frac{d^3 k_2}{(2\pi)^{3/2}} \int \frac{d^3 k_3}{(2\pi)^{3/2}} \int \frac{d^3 k_4}{(2\pi)^{3/2}} \\ &\quad a_{\mathbf{k}_1}^+(t) a_{\mathbf{k}_2}^+(t) a_{\mathbf{k}_3}(t) a_{\mathbf{k}_4}(t) U(\mathbf{r}-\mathbf{r}') e^{-i(\mathbf{k}_1-\mathbf{k}_4)\cdot\mathbf{r}-i(\mathbf{k}_2-\mathbf{k}_3)\cdot\mathbf{r}'} \\ &= \frac{1}{2} \int \frac{d^3 k_1}{(2\pi)^{3/2}} \int \frac{d^3 k_2}{(2\pi)^{3/2}} \int \frac{d^3 k_3}{(2\pi)^{3/2}} \int \frac{d^3 k_4}{(2\pi)^{3/2}} \\ &\quad a_{\mathbf{k}_1}^+(t) a_{\mathbf{k}_2}^+(t) a_{\mathbf{k}_3}(t) a_{\mathbf{k}_4}(t) U(\mathbf{k}_1-\mathbf{k}_4-\mathbf{k}_2+\mathbf{k}_3) \\ &= \frac{1}{2} \int d^3 k \int d^3 k' \int \frac{d^3 \mathbf{q}}{(2\pi)^3} \int \frac{d^3 \mathbf{q}'}{(2\pi)^3} a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}^+(t) a_{\mathbf{k}'-\mathbf{q}'}(t) a_{\mathbf{k}-\mathbf{q}}(t) U(\mathbf{q}-\mathbf{q}') \end{aligned}$$

where

$$\begin{aligned} \mathbf{k} &= \mathbf{k}_1 & \mathbf{q} &= \mathbf{k}_1 - \mathbf{k}_4 \\ \mathbf{k}' &= \mathbf{k}_2 & \mathbf{q}' &= \mathbf{k}_2 - \mathbf{k}_3 \end{aligned}$$

$$\rightarrow H = \int d^3 r \mathcal{H}$$

$$\begin{aligned}
&= \int d^3 k \frac{\hbar^2 \mathbf{k}^2}{2m} a_{\mathbf{k}}^+(t) a_{\mathbf{k}}(t) \\
&\quad + \int d^3 k \int \frac{d^3 q}{(2\pi)^3} V(\mathbf{q}) a_{\mathbf{k}}^+(t) a_{\mathbf{k}-\mathbf{q}}(t) \\
&\quad + \frac{1}{2} \int d^3 k \int d^3 k' \int \frac{d^3 q}{(2\pi)^3} \int \frac{d^3 q'}{(2\pi)^3} U(\mathbf{q}-\mathbf{q}') a_{\mathbf{k}}^+(t) a_{\mathbf{k}'}^+(t) a_{\mathbf{k}'-\mathbf{q}'}(t) a_{\mathbf{k}-\mathbf{q}}(t)
\end{aligned}$$

Systems with Finite Volume

Momentum (plane wave) expansion:

$$\phi(x) = \frac{1}{\sqrt{V}} \sum_{\mathbf{k}} \tilde{a}_{\mathbf{k}}(t) e^{i\mathbf{k}\cdot\mathbf{r}} \quad \mathbf{k} = \frac{2\pi}{V^{1/3}} (n_1, n_2, n_3)$$

$$\int d^3 r e^{i\mathbf{k}\cdot\mathbf{r}} = V \delta_{\mathbf{k},0}$$

$$\tilde{a}_{\mathbf{k}}(t) = \frac{1}{\sqrt{V}} \int d^3 r \phi(x) e^{-i\mathbf{k}\cdot\mathbf{r}} \rightarrow \sum_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{r}} = V \delta(\mathbf{r})$$

$$\phi(x) \approx \int \frac{d^3 k}{(2\pi)^{3/2}} a_{\mathbf{k}}(t) e^{i\mathbf{k}\cdot\mathbf{r}} \rightarrow \tilde{a}_{\mathbf{k}}(t) \approx \sqrt{\frac{(2\pi)^3}{V}} a_{\mathbf{k}}(t)$$

$$\begin{aligned}
[\tilde{a}_{\mathbf{k}}(t), \tilde{a}_{\mathbf{k}'}^+(t)]_{\mp} &= \frac{1}{V} \int d^3 r \int d^3 r' [\phi(t, \mathbf{r}) \phi^+(t, \mathbf{r}')]_{\mp} e^{-i\mathbf{k}\cdot\mathbf{r} + i\mathbf{k}'\cdot\mathbf{r}'} \\
&= \frac{1}{V} \int d^3 r \int d^3 r' \delta(\mathbf{r}-\mathbf{r}') e^{-i\mathbf{k}\cdot\mathbf{r} + i\mathbf{k}'\cdot\mathbf{r}'} \\
&= \frac{1}{V} \int d^3 r e^{i(\mathbf{k}'-\mathbf{k})\cdot\mathbf{r}} \\
&= \delta_{\mathbf{k}\mathbf{k}'}
\end{aligned}$$

$$\begin{aligned}
\mathbf{p} &= \int d^3 r \phi^+(x) \frac{\hbar}{i} \nabla \phi(x) \\
&= \frac{1}{V} \int d^3 r \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}(t) e^{-i\mathbf{k}\cdot\mathbf{r}} \frac{\hbar}{i} \nabla e^{i\mathbf{k}'\cdot\mathbf{r}} \\
&= \frac{1}{V} \int d^3 r \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}(t) \hbar \mathbf{k}' e^{i(\mathbf{k}'-\mathbf{k})\cdot\mathbf{r}} \\
&= \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}(t) \hbar \mathbf{k}' \delta_{\mathbf{k}\mathbf{k}'} \\
&= \sum_{\mathbf{k}} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}}(t) \hbar \mathbf{k}
\end{aligned}$$

Similarly

$$\begin{aligned}
\int d^3 r \frac{\hbar^2}{2m} \nabla \phi^+(x) \cdot \nabla \phi(x) &= -\frac{\hbar^2}{2m} \int d^3 r \phi^+(x) \nabla^2 \phi(x) \\
&= \sum_{\mathbf{k}} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}}(t) \frac{\hbar^2 \mathbf{k}^2}{2m}
\end{aligned}$$

$$\begin{aligned}
\int d^3 r V(\mathbf{x}) \phi^+(\mathbf{x}) \phi(\mathbf{x}) &= \frac{1}{V} \int d^3 r \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}(t) V(\mathbf{r}) e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}} \\
&= \frac{1}{V} \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}(t) V(\mathbf{k}-\mathbf{k}') & f(\mathbf{k}) &= \int d^3 r f(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{r}} \\
&= \frac{1}{V} \sum_{\mathbf{k}} \sum_{\mathbf{q}} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}-\mathbf{q}}(t) V(\mathbf{q}) & \mathbf{q} &= \mathbf{k}-\mathbf{k}'
\end{aligned}$$

$$\begin{aligned}
&\frac{1}{2} \int d^3 r \int d^3 r' \phi^+(\mathbf{x}) \phi^+(\mathbf{x}') U(\mathbf{r}-\mathbf{r}') \phi(\mathbf{x}') \phi(\mathbf{x}) \Big|_{t=t'} \\
&= \frac{1}{2V^2} \int d^3 r \int d^3 r' \sum_{\mathbf{k}_1} \sum_{\mathbf{k}_2} \sum_{\mathbf{k}_3} \sum_{\mathbf{k}_4} \tilde{a}_{\mathbf{k}_1}^+(t) \tilde{a}_{\mathbf{k}_2}^+(t) \tilde{a}_{\mathbf{k}_3}(t) \tilde{a}_{\mathbf{k}_4}(t) \\
&\quad U(\mathbf{r}-\mathbf{r}') e^{-i(\mathbf{k}_1-\mathbf{k}_4)\cdot\mathbf{r}-i(\mathbf{k}_2-\mathbf{k}_3)\cdot\mathbf{r}'} \\
&= \frac{1}{2V^2} \sum_{\mathbf{k}_1} \sum_{\mathbf{k}_2} \sum_{\mathbf{k}_3} \sum_{\mathbf{k}_4} \tilde{a}_{\mathbf{k}_1}^+(t) \tilde{a}_{\mathbf{k}_2}^+(t) \tilde{a}_{\mathbf{k}_3}(t) \tilde{a}_{\mathbf{k}_4}(t) U(\mathbf{k}_1-\mathbf{k}_4-\mathbf{k}_2+\mathbf{k}_3) \\
&= \frac{1}{2V^2} \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \sum_{\mathbf{q}} \sum_{\mathbf{q}'} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}^+(t) \tilde{a}_{\mathbf{k}'-\mathbf{q}'}(t) \tilde{a}_{\mathbf{k}-\mathbf{q}}(t) U(\mathbf{q}-\mathbf{q}')
\end{aligned}$$

where $\mathbf{k} = \mathbf{k}_1$ $\mathbf{q} = \mathbf{k}_1 - \mathbf{k}_4$
 $\mathbf{k}' = \mathbf{k}_2$ $\mathbf{q}' = \mathbf{k}_2 - \mathbf{k}_3$

$$\begin{aligned}
\rightarrow H &= \int d^3 r \mathcal{H} \\
&= \sum_{\mathbf{k}} \frac{\hbar^2 \mathbf{k}^2}{2m} \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}}(t) \\
&\quad + \frac{1}{V} \sum_{\mathbf{k}} \sum_{\mathbf{q}} V(\mathbf{q}) \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}-\mathbf{q}}(t) \\
&\quad + \frac{1}{2V^2} \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \sum_{\mathbf{q}} \sum_{\mathbf{q}'} U(\mathbf{q}-\mathbf{q}') \tilde{a}_{\mathbf{k}}^+(t) \tilde{a}_{\mathbf{k}'}^+(t) \tilde{a}_{\mathbf{k}'-\mathbf{q}'}(t) \tilde{a}_{\mathbf{k}-\mathbf{q}}(t)
\end{aligned}$$

Alternative Lagrangian

$$\begin{aligned}
\mathcal{L} &= \frac{1}{2} (i \hbar \varphi^* \partial_t \varphi - i \hbar \varphi \partial_t \varphi^*) - \frac{\hbar^2}{2m} \nabla \varphi^* \cdot \nabla \varphi - V \varphi^* \varphi \\
\frac{\partial \mathcal{L}}{\partial (\partial_t \varphi^*)} &= -\frac{1}{2} i \hbar \varphi & \frac{\partial \mathcal{L}}{\partial (\partial_t \varphi)} &= -\frac{1}{2} i \hbar \partial_t \varphi \\
\frac{\partial \mathcal{L}}{\partial (\nabla \varphi^*)} &= -\frac{\hbar^2}{2m} \nabla \varphi & \nabla \cdot \frac{\partial \mathcal{L}}{\partial (\nabla \varphi^*)} &= -\frac{\hbar^2}{2m} \nabla^2 \varphi \\
\frac{\partial \mathcal{L}}{\partial \varphi^*} &= \frac{1}{2} i \hbar \partial_t \varphi - V \varphi \\
\rightarrow -\frac{\hbar^2}{2m} \nabla^2 \varphi - i \hbar \partial_t \varphi + V \varphi &= 0 \\
\frac{\partial \mathcal{L}}{\partial (\partial_t \varphi)} &= \frac{1}{2} i \hbar \varphi^* & \frac{\partial \mathcal{L}}{\partial (\partial_t \varphi^*)} &= \frac{1}{2} i \hbar \partial_t \varphi^*
\end{aligned}$$

$$\frac{\partial \mathcal{L}}{\partial (\nabla \varphi)} = -\frac{\hbar^2}{2m} \nabla \varphi^* \qquad \nabla \cdot \frac{\partial \mathcal{L}}{\partial (\nabla \varphi)} = -\frac{\hbar^2}{2m} \nabla^2 \varphi^*$$

$$\frac{\partial \mathcal{L}}{\partial \varphi} = -\frac{1}{2} i \hbar \partial_t \varphi^* - V \varphi^*$$

$$\rightarrow -\frac{\hbar^2}{2m} \nabla^2 \varphi^* + i \hbar \partial_t \varphi^* + V \varphi^* = 0$$

$$\pi = \frac{\partial \mathcal{L}}{\partial (\partial_t \varphi)} = \frac{1}{2} i \hbar \varphi^*$$

$$\pi^* = \frac{\partial \mathcal{L}}{\partial (\partial_t \varphi^*)} = -\frac{1}{2} i \hbar \varphi = (\pi)^*$$

$$\mathcal{H} = \frac{1}{2} (i \hbar \varphi^* \partial_t \varphi - i \hbar \varphi \partial_t \varphi^*) - \mathcal{L} = \frac{\hbar^2}{2m} \nabla \varphi^* \cdot \nabla \varphi + V \varphi^* \varphi$$

2nd Quantization via equal-time commutation relations :

$$[\varphi(\mathbf{x}, t), \pi(\mathbf{x}', t)]_{\mp} = i \hbar \delta(\mathbf{x} - \mathbf{x}') \qquad \text{for } \begin{array}{l} \text{bosons} \\ \text{fermions} \end{array}$$

$$\rightarrow [\varphi(\mathbf{x}, t), \varphi^+(\mathbf{x}', t)]_{\mp} = 2 \delta(\mathbf{x} - \mathbf{x}')$$