

3.3. Real Klein-Gordon Field

Ref: M.Kaku, "Quantum Field Theory", §3.2.

$$\begin{aligned}
 x^\mu &= (ct, \mathbf{r}) & x_\mu &= \eta_{\mu\nu} x^\nu = (ct, -\mathbf{r}) \\
 \partial_\mu &= \frac{\partial}{\partial x^\mu} = \left(\frac{\partial}{c \partial t}, \nabla \right) & \partial^\mu &= \frac{\partial}{\partial x_\mu} = \left(\frac{\partial}{c \partial t}, -\nabla \right) \\
 \eta_{\mu\nu} &= \eta^{\mu\nu} = \text{diag}(1, -1, -1, -1)
 \end{aligned}$$

Caution: Ezawa used $\eta_{\mu\nu} = \eta^{\mu\nu} = \text{diag}(-1, 1, 1, 1)$.

Lagrangian

$$\text{Action: } S = \int dt L = \int dt \int d^3 r \mathcal{L} = \frac{1}{c} \int d^4 x \mathcal{L}$$

$$E^2 = \mathbf{p}^2 c^2 + m^2 c^4$$

$$E \rightarrow i \hbar \frac{\partial}{\partial t} \quad \mathbf{p} \rightarrow \frac{\hbar}{i} \nabla$$

$$\text{i.e. } p^\mu = \left(\frac{E}{c}, \mathbf{p} \right) \rightarrow i \hbar \partial^\mu = \left(i \hbar \frac{\partial}{c \partial t}, \frac{\hbar}{i} \nabla \right)$$

$$p_\mu = \left(\frac{E}{c}, -\mathbf{p} \right) \rightarrow i \hbar \partial_\mu = \left(i \hbar \frac{\partial}{c \partial t}, -\frac{\hbar}{i} \nabla \right)$$

$$\rightarrow \frac{E^2}{c^2} - \mathbf{p}^2 = p^\mu p_\mu = -\hbar^2 \partial^\mu \partial_\mu = -\hbar^2 \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \right) = m^2 c^2$$

$$\text{i.e. } \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 + \frac{m^2 c^2}{\hbar^2} \right) \phi = 0 \quad (\text{Klein-Gordon eq.})$$

$$\text{Set } \mathcal{L} = \frac{1}{2} f \left(\partial_\mu \phi \cdot \partial^\mu \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right)$$

where f is a parameter that may involve m , c , & \hbar .

$$\begin{aligned}
 \mathcal{L} &= \frac{1}{2} f \left[(\partial_\mu \phi)^2 - \frac{m^2 c^2}{\hbar^2} \phi^2 \right] = \frac{1}{2} f \left[\frac{1}{c^2} \left(\frac{\partial \phi}{\partial t} \right)^2 - \nabla \phi \cdot \nabla \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right] \\
 &= \frac{1}{2} f \left(\eta^{\mu\nu} \partial_\mu \phi \cdot \partial_\nu \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right)
 \end{aligned}$$

Kaku's choice: $f = \hbar = c = 1$.

Ezawa's choice: $f = \hbar^2$.

Note: Changing f only changes the units of ϕ .

Units:

$$[\mathcal{L}] = \left[\frac{E}{L^3} \right] = \left[f \frac{1}{L^2} \phi^2 \right]$$

$$\rightarrow [\phi] = \left[\left(\frac{E}{f L} \right)^{1/2} \right]$$

$$\frac{\partial \mathcal{L}}{\partial \phi} = -f \frac{m^2 c^2}{\hbar^2} \phi$$

$$\begin{aligned}\frac{\partial \mathcal{L}}{\partial \partial_\mu \phi} &= \frac{1}{2} f \eta^{\sigma\nu} (\delta_\sigma^\mu \partial_\nu \phi + \delta_\nu^\mu \partial_\sigma \phi) \\ &= \frac{1}{2} f (\eta^{\mu\nu} \partial_\nu \phi + \eta^{\sigma\mu} \partial_\sigma \phi) \\ &= \frac{1}{2} f (\partial^\mu \phi + \partial^\mu \phi) = f \partial^\mu \phi\end{aligned}$$

$$\partial_\mu \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi} = f \partial_\mu \partial^\mu \phi$$

$$\rightarrow \left(\partial_\mu \partial^\mu + \frac{m^2 c^2}{\hbar^2} \right) \phi = 0 \quad \text{Klein-Gordon eq.}$$

$$\text{or } \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 + \frac{m^2 c^2}{\hbar^2} \right) \phi = 0$$

Set $p^\mu \rightarrow i \hbar \partial^\mu$, then the K-G eq becomes

$$\left(-\frac{1}{\hbar^2} p_\mu p^\mu + \frac{m^2 c^2}{\hbar^2} \right) \phi = \frac{1}{\hbar^2} (-E^2 c^2 + \mathbf{p}^2 + m^2 c^2) \phi = 0$$

$$\rightarrow E^2 = \mathbf{p}^2 c^2 + m^2 c^4 \quad \rightarrow \quad E = \pm \sqrt{\mathbf{p}^2 c^2 + m^2 c^4}$$

Hamiltonian

$$\pi = \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \frac{\partial \mathcal{L}}{c \partial_0 \phi} = f \frac{1}{c} \partial^0 \phi = f \frac{1}{c^2} \dot{\phi}$$

$$\begin{aligned}\mathcal{H} &= \pi \dot{\phi} - \mathcal{L} \\ &= f \frac{1}{c^2} \dot{\phi}^2 - \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 - \nabla \phi \cdot \nabla \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right) \\ &= \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 + \nabla \phi \cdot \nabla \phi + \frac{m^2 c^2}{\hbar^2} \phi^2 \right)\end{aligned}$$

Energy-Momentum Tensor

$$\mathcal{L} = \frac{1}{2} f \left(\partial_\mu \phi \cdot \partial^\mu \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right)$$

$$\begin{aligned}\Theta^{\mu\nu} &= \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi} \partial^\nu \phi - \eta^{\mu\nu} \mathcal{L} \\ &= \frac{1}{2} f \left[2 \partial^\mu \phi \partial^\nu \phi - \eta^{\mu\nu} \left(\partial_\tau \phi \cdot \partial^\tau \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right) \right]\end{aligned}$$

$$\begin{aligned}\therefore \Theta^{00} &= \frac{1}{2} f \left[\frac{2}{c^2} \dot{\phi}^2 - \left(\frac{1}{c^2} \dot{\phi}^2 - \nabla \phi \cdot \nabla \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right) \right] \\ &= \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 + \nabla \phi \cdot \nabla \phi + \frac{m^2 c^2}{\hbar^2} \phi^2 \right) = \mathcal{H}\end{aligned}$$

$$\Theta^{0j} = \Theta^{j0} = f \frac{1}{c} \dot{\phi} \partial^j \phi = c \mathcal{P}^j$$

$$\begin{aligned} \rightarrow \quad p^j &= \int d^3 r \mathcal{P}^j = \frac{f}{c^2} \int d^3 r \dot{\phi} \partial^j \phi \\ \mathbf{p} &= -\frac{f}{c^2} \int d^3 r \dot{\phi} \nabla \phi \end{aligned}$$

Plane Wave Solutions

Let $\phi_{\mathbf{k}}(x) = e^{-i \mathbf{k} \cdot \mathbf{x}}$ where $\mathbf{k} \cdot \mathbf{x} = \omega t - \mathbf{k} \cdot \mathbf{r}$

$$\text{KG eq.:} \quad \left(\partial_\mu \partial^\mu + \frac{m^2 c^2}{\hbar^2} \right) \phi_{\mathbf{k}} = \left(-k_\mu k^\mu + \frac{m^2 c^2}{\hbar^2} \right) \phi_{\mathbf{k}} = 0$$

$$\rightarrow \quad k_\mu k^\mu = k^2 = \left(\frac{\omega}{c} \right)^2 - \mathbf{k}^2 = \frac{m^2 c^2}{\hbar^2}$$

\therefore Plane wave solutions are of the form

$$\phi_{\mathbf{k}}^{(\pm)}(x) \propto e^{i(\mathbf{k} \cdot \mathbf{r} \mp \omega_{\mathbf{k}} t)}$$

$$\text{with} \quad \omega_{\mathbf{k}} = \sqrt{\mathbf{k}^2 c^2 + \frac{m^2 c^4}{\hbar^2}}$$

For \mathbf{k} continuous,

$$\int d^3 r e^{i(\mathbf{k} - \mathbf{k}') \cdot \mathbf{r}} = (2\pi)^3 \delta(\mathbf{k} - \mathbf{k}')$$

$$\int d^3 r \phi_{\mathbf{k}}^{(\pm)*}(x) \phi_{\mathbf{k}'}^{(\pm)}(x) = \delta(\mathbf{k} - \mathbf{k}') \rightarrow \phi_{\mathbf{k}}^{(\pm)}(x) = \frac{1}{(2\pi)^{3/2}} e^{i(\mathbf{k} \cdot \mathbf{r} \mp \omega_{\mathbf{k}} t)}$$

General Solutions

Plane wave expansion:

$$\begin{aligned} \phi(x) &= \int d^3 k \left[A(\mathbf{k}) \phi_{\mathbf{k}}^{(+)}(x) + B(\mathbf{k}) \phi_{\mathbf{k}}^{(-)}(x) \right] \\ &= \int \frac{d^3 k}{(2\pi)^{3/2}} \left[A(\mathbf{k}) e^{i(\mathbf{k} \cdot \mathbf{r} - \omega_{\mathbf{k}} t)} + B(\mathbf{k}) e^{i(\mathbf{k} \cdot \mathbf{r} + \omega_{\mathbf{k}} t)} \right] \\ &= \int \frac{d^3 k}{(2\pi)^{3/2}} \left[A(\mathbf{k}) e^{i(\mathbf{k} \cdot \mathbf{r} - \omega_{\mathbf{k}} t)} + B(-\mathbf{k}) e^{-i(\mathbf{k} \cdot \mathbf{r} - \omega_{\mathbf{k}} t)} \right] \\ &= \int \frac{d^3 k}{(2\pi)^{3/2}} \left[A(\mathbf{k}) e^{-i \mathbf{k} \cdot \mathbf{x}} + B(-\mathbf{k}) e^{i \mathbf{k} \cdot \mathbf{x}} \right]_{\omega=\omega_{\mathbf{k}}} \end{aligned}$$

$$\mathcal{H}^* = \mathcal{H} \quad \rightarrow \quad \phi^*(x) = \phi(x) \quad (\text{real field})$$

$$\phi^*(x) = \int \frac{d^3 k}{(2\pi)^{3/2}} \left[A^*(\mathbf{k}) e^{i \mathbf{k} \cdot \mathbf{x}} + B^*(-\mathbf{k}) e^{-i \mathbf{k} \cdot \mathbf{x}} \right]_{\omega=\omega_{\mathbf{k}}}$$

$$\rightarrow \quad B(-\mathbf{k}) = A^*(\mathbf{k})$$

$$\begin{aligned} \phi(x) &= \int \frac{d^3 k}{(2\pi)^{3/2}} \left[A(\mathbf{k}) e^{-i \mathbf{k} \cdot \mathbf{x}} + A^*(\mathbf{k}) e^{i \mathbf{k} \cdot \mathbf{x}} \right]_{\omega=\omega_{\mathbf{k}}} \\ &= \int \frac{d^3 k}{(2\pi)^{3/2}} \left[A(\mathbf{k}) e^{i(\mathbf{k} \cdot \mathbf{r} - \omega_{\mathbf{k}} t)} + A^*(\mathbf{k}) e^{-i(\mathbf{k} \cdot \mathbf{r} - \omega_{\mathbf{k}} t)} \right] \end{aligned}$$

$$\begin{aligned}\dot{\phi}(x) &= \int \frac{d^3 k}{(2\pi)^{3/2}} i \omega_k [-A(\mathbf{k}) e^{i(\mathbf{k}\cdot\mathbf{r} - \omega_k t)} + A^*(\mathbf{k}) e^{-i(\mathbf{k}\cdot\mathbf{r} - \omega_k t)}] \\ \pi(x) &= f \frac{1}{c^2} \dot{\phi}(x) = f \frac{1}{c^2} \int \frac{d^3 k}{(2\pi)^{3/2}} i \omega_k [-A(\mathbf{k}) e^{i(\mathbf{k}\cdot\mathbf{r} - \omega_k t)} + A^*(\mathbf{k}) e^{-i(\mathbf{k}\cdot\mathbf{r} - \omega_k t)}] \\ \int d^3 r e^{-i\mathbf{k}\cdot\mathbf{r}} \phi(x) &= \int \frac{d^3 k'}{(2\pi)^{3/2}} \int d^3 r [A(\mathbf{k}) e^{i(\mathbf{k}'-\mathbf{k})\cdot\mathbf{r} - i\omega'_k t} + A^*(\mathbf{k}') e^{-i(\mathbf{k}'+\mathbf{k})\cdot\mathbf{r} + i\omega'_k t}] \\ &= (2\pi)^{3/2} \int d^3 k' [A(\mathbf{k}) e^{-i\omega'_k t} \delta(\mathbf{k}'-\mathbf{k}) + A^*(\mathbf{k}') e^{i\omega'_k t} \delta(\mathbf{k}'+\mathbf{k})] \\ &= (2\pi)^{3/2} (A(\mathbf{k}) e^{-i\omega_k t} + A^*(-\mathbf{k}) e^{i\omega_k t})\end{aligned}$$

Similarly,

$$\begin{aligned}\int d^3 r e^{-i\mathbf{k}\cdot\mathbf{x}} \dot{\phi}(x) &= (2\pi)^{3/2} i \omega_k [-A(\mathbf{k}) e^{-i\omega_k t} + A^*(-\mathbf{k}) e^{i\omega_k t}] \\ \therefore (2\pi)^{3/2} 2 \omega_k A(\mathbf{k}) e^{-i\omega_k t} &= \int d^3 r e^{-i\mathbf{k}\cdot\mathbf{r}} [\omega_k \phi(x) + i \dot{\phi}(x)] \\ \rightarrow A(\mathbf{k}) &= \frac{1}{2 \omega_k (2\pi)^{3/2}} \int d^3 r e^{i\mathbf{k}\cdot\mathbf{x}} [\omega_k \phi(x) + i \dot{\phi}(x)] \\ &\text{(Note: Despite the presence of } x = (\mathbf{r}, t) \text{ in the integrand, } A(\mathbf{k}) \text{ is independent of } t \text{ \& } f. \text{)} \\ (2\pi)^{3/2} 2 \omega_k A^*(-\mathbf{k}) e^{i\omega_k t} &= \int d^3 r e^{-i\mathbf{k}\cdot\mathbf{r}} [\omega_k \phi(x) - i \dot{\phi}(x)] \\ \rightarrow A^*(-\mathbf{k}) &= \frac{1}{2 \omega_k (2\pi)^{3/2}} \int d^3 r e^{-i\omega_k t - i\mathbf{k}\cdot\mathbf{r}} [\omega_k \phi(x) - i \dot{\phi}(x)] \\ A^*(\mathbf{k}) &= \frac{1}{2 \omega_k (2\pi)^{3/2}} \int d^3 r e^{-i\mathbf{k}\cdot\mathbf{x}} [\omega_k \phi(x) - i \dot{\phi}(x)] \\ &= [A(\mathbf{k})]^* \quad \text{since } \phi^* = \phi.\end{aligned}$$

Units:

$$\begin{aligned}[A] &= [L^3 \phi] \\ [\phi] &= \left[\left(\frac{E}{fL} \right)^{1/2} \right] \rightarrow [A] = \left[L^3 \left(\frac{E}{fL} \right)^{1/2} \right] = \left[\left(\frac{EL^5}{f} \right)^{1/2} \right]\end{aligned}$$

Quantization

Canonical quantization:

$$\begin{aligned}\phi \text{ \& } A \text{ become operators.} \quad A^* \rightarrow A^\dagger. \\ [\phi(\mathbf{r}, t), \pi(\mathbf{r}', t)] &= i \hbar \delta(\mathbf{r} - \mathbf{r}') \quad \text{(equal-time commutator)} \\ [\phi(\mathbf{r}, t), \phi(\mathbf{r}', t)] &= [\pi(\mathbf{r}, t), \pi(\mathbf{r}', t)] = 0 \\ \rightarrow [\phi(\mathbf{r}, t), \dot{\phi}(\mathbf{r}', t)] &= i \frac{\hbar c^2}{f} \delta(\mathbf{r} - \mathbf{r}') \\ [\phi(\mathbf{r}, t), \phi(\mathbf{r}', t)] &= [\dot{\phi}(\mathbf{r}, t), \dot{\phi}(\mathbf{r}', t)] = 0\end{aligned}$$

Note: For $t = t'$,

$$(x - x')^2 = -(\mathbf{r} - \mathbf{r}')^2 \leq 0 \text{ (space-like)}$$

$$A(\mathbf{k}) = \frac{1}{2 \omega_k (2\pi)^{3/2}} \int d^3 r e^{i\mathbf{k}\cdot\mathbf{x}} [\omega_k \phi(x) + i \dot{\phi}(x)]$$

$$\begin{aligned}
&= \frac{1}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \int d^3 r [-i c (\partial_0 e^{i \mathbf{k} \cdot \mathbf{x}}) \phi(x) + i e^{i \mathbf{k} \cdot \mathbf{x}} c \partial_0 \phi(x)] \\
&= i \frac{c}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \int d^3 r e^{i \mathbf{k} \cdot \mathbf{x}} \overleftrightarrow{\partial}_0 \phi(x)
\end{aligned}$$

where

$$A \overleftrightarrow{\partial} B = A \partial B - (\partial A) B$$

$$[A] = \left[\frac{L}{T T^{-1}} L^3 \frac{1}{L} \phi \right] = [L^3 \phi] \quad (\text{check})$$

Similarly,

$$\begin{aligned}
A^+(\mathbf{k}) &= \frac{1}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \int d^3 r e^{-i \mathbf{k} \cdot \mathbf{x}} [\omega_{\mathbf{k}} \phi(x) - i \dot{\phi}(x)] \\
&= \frac{c}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \int d^3 r [i (\partial_0 e^{-i \mathbf{k} \cdot \mathbf{x}}) \phi(x) - i e^{-i \mathbf{k} \cdot \mathbf{x}} \partial_0 \phi(x)] \\
&= -i \frac{c}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \int d^3 r e^{-i \mathbf{k} \cdot \mathbf{x}} \overleftrightarrow{\partial}_0 \phi(x) \\
&= [A(\mathbf{k})]^+ \quad \text{since } \phi^+ = \phi
\end{aligned}$$

$$\begin{aligned}
[A(\mathbf{k}), A^+(\mathbf{k}')] &= \left(\frac{1}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \right)^2 \int d^3 r \int d^3 r' e^{i \mathbf{k} \cdot \mathbf{x} - i \mathbf{k}' \cdot \mathbf{x}'} \left[\omega_{\mathbf{k}} \phi(x) + i \frac{c^2}{f} \pi(x), \omega_{\mathbf{k}'} \phi(x') - i \frac{c^2}{f} \pi(x') \right]_{t=t'} \\
&= \left(\frac{1}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \right)^2 \int d^3 r \int d^3 r' e^{i \mathbf{k} \cdot \mathbf{x} - i \mathbf{k}' \cdot \mathbf{x}'} \hbar \frac{c^2}{f} [\omega_{\mathbf{k}} \delta(\mathbf{r} - \mathbf{r}') + \omega_{\mathbf{k}'} \delta(\mathbf{r} - \mathbf{r}')]_{t=t'} \\
&= \frac{1}{f} \left(\frac{c}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \right)^2 \hbar \int d^3 r e^{i (\mathbf{k} - \mathbf{k}') \cdot \mathbf{x}} [\omega_{\mathbf{k}} + \omega_{\mathbf{k}'}] \\
&= \frac{1}{f} \left(\frac{c}{2 \omega_{\mathbf{k}} (2 \pi)^{3/2}} \right)^2 (2 \pi)^3 \delta(\mathbf{k} - \mathbf{k}') 2 \hbar \omega_{\mathbf{k}}
\end{aligned}$$

$$\text{Similarly, } [A(\mathbf{k}), A(\mathbf{k}')] = [A^+(\mathbf{k}), A^+(\mathbf{k}')] = 0$$

Setting

$$[a_{\mathbf{k}}, a_{\mathbf{k}'}] = [a_{\mathbf{k}}^+, a_{\mathbf{k}'}^+] = 0$$

$$[a_{\mathbf{k}}, a_{\mathbf{k}'}^+] = \delta(\mathbf{k} - \mathbf{k}')$$

$$\rightarrow a_{\mathbf{k}} = \frac{2 \omega_{\mathbf{k}}}{c} \sqrt{f} \frac{A(\mathbf{k})}{\sqrt{2 \hbar \omega_{\mathbf{k}}}}$$

$$A(\mathbf{k}) = \frac{c}{2 \omega_{\mathbf{k}} \sqrt{f}} \sqrt{2 \hbar \omega_{\mathbf{k}}} a_{\mathbf{k}} = \frac{\hbar c}{\sqrt{f}} \sqrt{\frac{1}{2 \hbar \omega_{\mathbf{k}}}} a_{\mathbf{k}}$$

$$\phi(x) = \int \frac{d^3 k}{(2 \pi)^{3/2}} [A(\mathbf{k}) e^{-i \mathbf{k} \cdot \mathbf{x}} + A^+(\mathbf{k}) e^{i \mathbf{k} \cdot \mathbf{x}}]$$

$$\therefore \phi(x) = \frac{\hbar c}{\sqrt{f}} \int \frac{d^3 k}{\sqrt{(2 \pi)^3 2 \hbar \omega_{\mathbf{k}}}} [a_{\mathbf{k}} e^{-i \mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}}^+ e^{i \mathbf{k} \cdot \mathbf{x}}]$$

$$\dot{\phi}(x) = \frac{\hbar c}{\sqrt{f}} \int \frac{d^3 k}{\sqrt{(2 \pi)^3 2 \hbar \omega_{\mathbf{k}}}} i \omega_{\mathbf{k}} [-a_{\mathbf{k}} e^{-i \mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}}^+ e^{i \mathbf{k} \cdot \mathbf{x}}]$$

$$\begin{aligned}
\pi(x) &= f \frac{1}{c^2} \dot{\phi}(x) = \sqrt{f} \frac{\hbar}{c} \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_k}} i\omega_k [-a_k e^{-ik \cdot x} + a_k^+ e^{ik \cdot x}] \\
A(\mathbf{k}) &= i \frac{c}{2\omega_k (2\pi)^{3/2}} \int d^3 r e^{ik \cdot x} \overleftrightarrow{\partial}_0 \phi(x) \\
\rightarrow a_k &= \frac{2\omega_k}{c} \sqrt{f} \frac{1}{\sqrt{2\hbar\omega_k}} i \frac{c}{2\omega_k (2\pi)^{3/2}} \int d^3 r e^{ik \cdot x} \overleftrightarrow{\partial}_0 \phi(x) \\
&= \sqrt{f} \int \frac{d^3 r}{\sqrt{(2\pi)^3 2\hbar\omega_k}} i e^{ik \cdot x} \overleftrightarrow{\partial}_0 \phi(x) \\
&= \sqrt{f} \int \frac{d^3 r}{\sqrt{(2\pi)^3 2\hbar\omega_k}} \frac{1}{c} e^{ik \cdot x} [\omega_k \phi(x) + i \dot{\phi}(x)] \\
a_k^+ &= \sqrt{f} \int \frac{d^3 r}{\sqrt{(2\pi)^3 2\hbar\omega_k}} (-i) e^{-ik \cdot x} \overleftrightarrow{\partial}_0 \phi(x) \\
&= \sqrt{f} \int \frac{d^3 r}{\sqrt{(2\pi)^3 2\hbar\omega_k}} \frac{1}{c} e^{-ik \cdot x} [\omega_k \phi(x) - i \dot{\phi}(x)]
\end{aligned}$$

Eq. of Motion

$$\begin{aligned}
i\hbar \partial_t \dot{\phi} &= [\dot{\phi}, H] = \int d^3 r [\dot{\phi}, \mathcal{H}] \\
&= \frac{1}{2} f \int d^3 r \left[\dot{\phi}(x), \nabla' \phi(x') \cdot \nabla' \phi(x') + \frac{m^2 c^2}{\hbar^2} \phi(x')^2 \right]_{t=t'} \\
[\dot{\phi}(x), \phi(x')^2]_{t=t'} &= -2i \frac{\hbar c^2}{f} \delta(\mathbf{r} - \mathbf{r}') \phi(x') \\
[\dot{\phi}(x), \nabla' \phi(x') \cdot \nabla' \phi(x')]_{t=t'} &= [\dot{\phi}(x), \nabla' \phi(x')] \cdot \nabla' \phi(x') + \nabla' \phi(x') \cdot [\dot{\phi}(x), \nabla' \phi(x')] \\
&= -2i \frac{\hbar c^2}{f} [\nabla' \delta(\mathbf{r} - \mathbf{r}')] \cdot \nabla' \phi(x') \\
&= -2i \frac{\hbar c^2}{f} \{ \nabla' \cdot [\delta(\mathbf{r} - \mathbf{r}') \nabla' \phi(x')] - \delta(\mathbf{r} - \mathbf{r}') \nabla'^2 \phi(x') \} \\
\rightarrow i\hbar \partial_t \dot{\phi} &= i\hbar c^2 \left[\nabla^2 \phi(x) - \frac{m^2 c^2}{\hbar^2} \phi(x) \right] \\
\text{or } \frac{\partial^2 \phi}{c^2 \partial t^2} - \nabla^2 \phi + \frac{m^2 c^2}{\hbar^2} \phi &= 0 \quad (\text{K-G eq.})
\end{aligned}$$

I-Particle State

Wave function of a particle with momentum $\hbar \mathbf{k}$:

$$\begin{aligned}
\Phi_{\mathbf{k}}(x) &= \langle 0 | \phi(x) | \mathbf{k} \rangle = \langle 0 | \phi(x) a_{\mathbf{k}}^+ | 0 \rangle \\
&= \frac{\hbar c}{\sqrt{f}} \int \frac{d^3 k'}{\sqrt{(2\pi)^3 2\hbar\omega_{k'}}} [\langle 0 | a_{\mathbf{k}'} | \mathbf{k} \rangle e^{-ik' \cdot x} + \langle 0 | a_{\mathbf{k}'}^+ | \mathbf{k} \rangle e^{ik' \cdot x}]
\end{aligned}$$

$$\langle 0 | a_{\mathbf{k}'} | \mathbf{k} \rangle = \langle \mathbf{k}' | \mathbf{k} \rangle = \delta(\mathbf{k} - \mathbf{k}') \quad \langle 0 | a_{\mathbf{k}'}^\dagger | \mathbf{k} \rangle = 0$$

$$\rightarrow \Phi_{\mathbf{k}}(x) = \frac{\hbar c}{\sqrt{f}} \frac{1}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}'}}} e^{-i\mathbf{k}\cdot\mathbf{x}} = \frac{c}{\sqrt{f}} \sqrt{\frac{\hbar}{(2\pi)^3 2\omega_{\mathbf{k}'}}} e^{i\mathbf{k}\cdot\mathbf{r} - i\omega_{\mathbf{k}}t}$$

$$\therefore \mathbf{p}\Phi_{\mathbf{k}}(x) = \frac{\hbar}{i} \nabla \Phi_{\mathbf{k}}(x) = \hbar \mathbf{k} \Phi_{\mathbf{k}}(x)$$

$$H\Phi_{\mathbf{k}}(x) = i\hbar \frac{\partial}{\partial t} \Phi_{\mathbf{k}}(x) = \hbar \omega_{\mathbf{k}} \Phi_{\mathbf{k}}(x)$$

$$E_{\mathbf{k}} = \hbar \omega_{\mathbf{k}} = \sqrt{\hbar^2 \mathbf{k}^2 c^2 + m^2 c^4}$$

For $m=0$,

$$E_{\mathbf{k}} = \hbar c |\mathbf{k}|$$

H & p

$$\mathcal{H} = \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 + \nabla \phi \cdot \nabla \phi + \frac{m^2 c^2}{\hbar^2} \phi^2 \right)$$

$$H = \int d^3 r \mathcal{H} = \int d^3 r \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 + \nabla \phi \cdot \nabla \phi + \frac{m^2 c^2}{\hbar^2} \phi^2 \right)$$

$$\phi(x) = \frac{\hbar c}{\sqrt{f}} \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}}}} (a_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{x}} + a_{\mathbf{k}}^\dagger e^{i\mathbf{k}\cdot\mathbf{x}})$$

$$\dot{\phi}(x) = \frac{\hbar c}{\sqrt{f}} \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}}}} i\omega_{\mathbf{k}} (-a_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{x}} + a_{\mathbf{k}}^\dagger e^{i\mathbf{k}\cdot\mathbf{x}})$$

$$\begin{aligned} \rightarrow \int d^3 r \dot{\phi}^2 &= \frac{\hbar^2 c^2}{f} \int d^3 r \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}}}} \int \frac{d^3 k'}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}'}}} \\ &\quad i\omega_{\mathbf{k}} (-a_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{x}} + a_{\mathbf{k}}^\dagger e^{i\mathbf{k}\cdot\mathbf{x}}) i\omega_{\mathbf{k}'} (-a_{\mathbf{k}'} e^{-i\mathbf{k}'\cdot\mathbf{x}} + a_{\mathbf{k}'}^\dagger e^{i\mathbf{k}'\cdot\mathbf{x}}) \\ &= \frac{\hbar^2 c^2}{f} \int \frac{d^3 k}{2\hbar\omega_{\mathbf{k}}} \omega_{\mathbf{k}}^2 (-a_{\mathbf{k}} a_{-\mathbf{k}} e^{-2i\omega_{\mathbf{k}}t} - a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger e^{2i\omega_{\mathbf{k}}t} + a_{\mathbf{k}} a_{\mathbf{k}}^\dagger + a_{\mathbf{k}}^\dagger a_{\mathbf{k}}) \end{aligned}$$

$$\therefore \frac{1}{2} f \int d^3 r \frac{1}{c^2} \dot{\phi}^2 = \frac{1}{4} \int d^3 k \hbar \omega_{\mathbf{k}} (-a_{\mathbf{k}} a_{-\mathbf{k}} e^{-2i\omega_{\mathbf{k}}t} - a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger e^{2i\omega_{\mathbf{k}}t} + a_{\mathbf{k}} a_{\mathbf{k}}^\dagger + a_{\mathbf{k}}^\dagger a_{\mathbf{k}})$$

$$\begin{aligned} \int d^3 r \nabla \phi \cdot \nabla \phi &= \frac{\hbar^2 c^2}{f} \int d^3 r \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}}}} \int \frac{d^3 k'}{\sqrt{(2\pi)^3 2\hbar\omega_{\mathbf{k}'}}} \\ &\quad (a_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{x}} - a_{\mathbf{k}}^\dagger e^{i\mathbf{k}\cdot\mathbf{x}}) (i\mathbf{k}) \cdot (i\mathbf{k}') (a_{\mathbf{k}'} e^{-i\mathbf{k}'\cdot\mathbf{x}} - a_{\mathbf{k}'}^\dagger e^{i\mathbf{k}'\cdot\mathbf{x}}) \\ &= \frac{\hbar^2 c^2}{f} \int \frac{d^3 k}{2\hbar\omega_{\mathbf{k}}} \mathbf{k}^2 (a_{\mathbf{k}} a_{-\mathbf{k}} e^{-2i\omega_{\mathbf{k}}t} + a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger e^{2i\omega_{\mathbf{k}}t} + a_{\mathbf{k}} a_{\mathbf{k}}^\dagger + a_{\mathbf{k}}^\dagger a_{\mathbf{k}}) \\ &= \frac{1}{2f} \int d^3 k \frac{\hbar \mathbf{k}^2 c^2}{\omega_{\mathbf{k}}} (a_{\mathbf{k}} a_{-\mathbf{k}} e^{-2i\omega_{\mathbf{k}}t} + a_{\mathbf{k}}^\dagger a_{-\mathbf{k}}^\dagger e^{2i\omega_{\mathbf{k}}t} + a_{\mathbf{k}} a_{\mathbf{k}}^\dagger + a_{\mathbf{k}}^\dagger a_{\mathbf{k}}) \end{aligned}$$

$$\begin{aligned}
\therefore \quad \frac{1}{2} f \int d^3 r \nabla \phi \cdot \nabla \phi &= \frac{1}{4} \int d^3 k \frac{\hbar}{\omega_k} \mathbf{k}^2 c^2 (a_k a_{-k} e^{-2i\omega_k t} + a_k^+ a_{-k}^+ e^{2i\omega_k t} + a_k a_k^+ + a_k^+ a_k) \\
\int d^3 r \phi^2 &= \frac{\hbar^2 c^2}{f} \int d^3 r \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_k}} \int \frac{d^3 k'}{\sqrt{(2\pi)^3 2\hbar\omega_{k'}}} \\
&\quad (a_k e^{-ik \cdot x} + a_k^+ e^{ik \cdot x}) (a_{k'} e^{-ik' \cdot x} + a_{k'}^+ e^{ik' \cdot x}) \\
&= \frac{\hbar^2 c^2}{f} \int \frac{d^3 k}{2\hbar\omega_k} (a_k a_{-k} e^{-2i\omega_k t} + a_k^+ a_{-k}^+ e^{2i\omega_k t} + a_k a_k^+ + a_k^+ a_k) \\
&= \frac{1}{2f} \int d^3 k \frac{\hbar c^2}{\omega_k} (a_k a_{-k} e^{-2i\omega_k t} + a_k^+ a_{-k}^+ e^{2i\omega_k t} + a_k a_k^+ + a_k^+ a_k) \\
\frac{1}{2} f \int d^3 r \frac{m^2 c^2}{\hbar^2} \phi^2 &= \frac{1}{4} \int d^3 k \frac{\hbar}{\omega_k} \frac{m^2 c^4}{\hbar^2} (a_k a_{-k} e^{-2i\omega_k t} + a_k^+ a_{-k}^+ e^{2i\omega_k t} + a_k a_k^+ + a_k^+ a_k)
\end{aligned}$$

$$\omega_k = \sqrt{\mathbf{k}^2 c^2 + \frac{m^2 c^4}{\hbar^2}}$$

$$\begin{aligned}
\therefore \quad H &= \frac{1}{4} \int d^3 k \left(\hbar\omega_k + \frac{\hbar}{\omega_k} \mathbf{k}^2 c^2 + \frac{\hbar}{\omega_k} \frac{m^2 c^4}{\hbar^2} \right) (a_k a_k^+ + a_k^+ a_k) \\
&\quad + \frac{1}{4} \int d^3 k \left(-\hbar\omega_k + \frac{\hbar}{\omega_k} \mathbf{k}^2 c^2 + \frac{\hbar}{\omega_k} \frac{m^2 c^4}{\hbar^2} \right) (a_k a_{-k} e^{-2i\omega_k t} + a_k^+ a_{-k}^+ e^{2i\omega_k t}) \\
&= \frac{1}{2} \int d^3 k \hbar\omega_k (a_k a_k^+ + a_k^+ a_k) \\
[a_k, a_{k'}^+] &= \delta(\mathbf{k} - \mathbf{k}') \quad \rightarrow \quad a_k a_k^+ = a_k^+ a_k + \delta(\mathbf{0})
\end{aligned}$$

$$\begin{aligned}
\therefore \quad H &= \int d^3 k \hbar\omega_k \left[a_k^+ a_k + \frac{1}{2} \delta(\mathbf{0}) \right] \\
&= \int d^3 k \hbar\omega_k a_k^+ a_k + \frac{1}{2} \hbar\omega_0
\end{aligned}$$

$$\begin{aligned}
\mathbf{p} &= -\frac{f}{c^2} \int d^3 r \dot{\phi} \nabla \phi \\
&= -\hbar^2 \int d^3 r \int \frac{d^3 k}{\sqrt{(2\pi)^3 2\hbar\omega_k}} \int \frac{d^3 k'}{\sqrt{(2\pi)^3 2\hbar\omega_{k'}}} \\
&\quad \times i\omega_k (-a_k e^{-ik \cdot x} + a_k^+ e^{ik \cdot x}) i (a_{k'} \mathbf{k}' e^{-ik' \cdot x} - a_{k'}^+ \mathbf{k}' e^{ik' \cdot x}) \\
&= \hbar^2 \int \frac{d^3 k}{2\hbar\omega_k} \omega_k \mathbf{k} (a_k a_{-k} e^{-2i\omega_k t} + a_k^+ a_{-k}^+ e^{2i\omega_k t} + a_k a_k^+ + a_k^+ a_k) \\
&= \frac{1}{2} \int d^3 k \hbar \mathbf{k} (a_k a_k^+ + a_k^+ a_k) \\
&= \int d^3 k \hbar \mathbf{k} \left(a_k^+ a_k + \frac{1}{2} \delta(\mathbf{0}) \right) \\
&= \int d^3 k \hbar \mathbf{k} a_k^+ a_k
\end{aligned}$$

Zero Point Energy

$$\begin{aligned}
 \langle 0 | H | 0 \rangle &= \int d^3 k \, \hbar \omega_k \left\langle 0 \left| a_k^\dagger a_k + \frac{1}{2} \delta(\mathbf{0}) \right| 0 \right\rangle \\
 &= \delta(\mathbf{0}) \int d^3 k \, \frac{1}{2} \hbar \omega_k \\
 &= \frac{V}{(2\pi)^3} \int d^3 k \, \frac{1}{2} \hbar \omega_k \quad (\text{zero-point energy})
 \end{aligned}$$

Normal ordering:

$$\begin{aligned}
 :H: &= \int d^3 r \, :\mathcal{H}: \\
 &= \frac{1}{2} \int d^3 k \, \hbar k \, :(a_k a_k^\dagger + a_k^\dagger a_k): \\
 &= \int d^3 k \, \hbar k \, a_k^\dagger a_k \\
 &= H - \langle 0 | H | 0 \rangle
 \end{aligned}$$