

4.7. Goldstone Theorem

Bose condensation & spontaneous symmetry breaking

$$\rightarrow \langle 0 | \phi(x) | 0 \rangle = v \neq 0 \quad (\text{ground state})$$

$$\text{with } \langle \underline{0} | \phi(x) | \underline{0} \rangle = 0 \quad (\text{false vacuum})$$

Let H be invariant under a continuous global transformation.

Spontaneous symmetry breaking \rightarrow gapless Goldstone mode (boson).

See the original paper,

J. Goldstone, A. Salam, S. Weinberg, PR 127, 965 (62)

for a more general discussion.

Phase Transformation

Let Q be the generator of a phase transformation:

$$e^{-i\alpha Q} \phi(x) e^{i\alpha Q} = e^{i\alpha} \phi(x) \quad \alpha = \text{real}$$

For $|\alpha| \ll 1$,

$$(1 - i\alpha Q) \phi(x) (1 + i\alpha Q) = (1 + i\alpha) \phi(x)$$

$$\rightarrow [\phi(x), Q] = \phi(x)$$

If H is invariant under the phase transformation, then

$$Q = \frac{1}{c} \int d^3x j^0(x) \quad (j^\mu = \text{Noether current})$$

is conserved, i.e., $\frac{dQ}{dt} = 0$.

Ground state is not invariant under the phase transformation since

$$|\alpha\rangle \equiv e^{i\alpha Q} |0\rangle \neq |0\rangle \quad \text{is another ground state}$$

$$\begin{aligned} \langle \alpha | \phi(x) | \alpha \rangle &= \langle 0 | e^{-i\alpha Q} \phi(x) e^{i\alpha Q} | 0 \rangle \\ &= e^{i\alpha} \langle \alpha | \phi(x) | \alpha \rangle \\ &= e^{i\alpha} v \end{aligned}$$

$$e^{i\alpha Q} |0\rangle \neq |0\rangle$$

$$|\alpha| \ll 1 \quad \rightarrow \quad (1 + i\alpha Q) |0\rangle \neq |0\rangle$$

$$\therefore Q |0\rangle \neq 0$$

$$\phi(x) | \underline{0} \rangle = 0 \quad \rightarrow \quad Q | \underline{0} \rangle = 0$$

Proof

$$\begin{aligned} \langle 0 | [\phi(x), Q] | 0 \rangle &= \int d^3y \langle 0 | [\phi(x), j^0(y)] | 0 \rangle \\ &= \langle 0 | \phi(x) | 0 \rangle = v \neq 0 \end{aligned}$$

$$\rightarrow I_+ - I_- = v$$

$$\text{where } I_+ = \int d^3y \langle 0 | \phi(x) j^0(y) | 0 \rangle$$

$$I_- = \int d^3y \langle 0 | j^0(y) \phi(x) | 0 \rangle$$

Completeness:

$$\int d\xi \int dE \int d^3k |E, \mathbf{k}, \xi\rangle \langle E, \mathbf{k}, \xi| = I$$

Translational invariance

$$\rightarrow \phi(x) = e^{-\frac{i}{\hbar}p \cdot x} \phi(0) e^{\frac{i}{\hbar}p \cdot x}$$

$$j^0(y) = e^{-\frac{i}{\hbar}p \cdot y} j^0(0) e^{\frac{i}{\hbar}p \cdot y}$$

$$\begin{aligned} \therefore I_+ &= \int d\xi \int dE \int d^3k \int d^3y \langle 0 | \phi(x) | E, \mathbf{k}, \xi \rangle \langle E, \mathbf{k}, \xi | j^0(y) | 0 \rangle \\ &= \int d\xi \int dE \int d^3k \int d^3y e^{i\mathbf{k} \cdot (\mathbf{x}-\mathbf{y}) - \frac{i}{\hbar}E\Delta t} \langle 0 | \phi(0) | E, \mathbf{k}, \xi \rangle \langle E, \mathbf{k}, \xi | j^0(0) | 0 \rangle \\ &= \int dE \int d^3k \int d^3y e^{i\mathbf{k} \cdot (\mathbf{x}-\mathbf{y}) - \frac{i}{\hbar}E\Delta t} \sigma_+(\mathbf{k}, E) \end{aligned}$$

where

$$\Delta t = t_x - t_y$$

$$\sigma_+(\mathbf{k}, E) = \int d\xi \langle 0 | \phi(0) | E, \mathbf{k}, \xi \rangle \langle E, \mathbf{k}, \xi | j^0(0) | 0 \rangle \quad (\text{spectral function})$$

$$\begin{aligned} I_+ &= (2\pi)^3 \int dE \int d^3k \delta(\mathbf{k}) e^{-\frac{i}{\hbar}E\Delta t} \sigma_+(\mathbf{k}, E) \\ &= (2\pi)^3 \int dE e^{-\frac{i}{\hbar}E\Delta t} \sigma_+(\mathbf{0}, E) \end{aligned}$$

Similarly,

$$I_- = (2\pi)^3 \int dE e^{+\frac{i}{\hbar}E\Delta t} \sigma_-(\mathbf{0}, E)$$

$$\text{with } \sigma_-(\mathbf{k}, E) = \int d\xi \langle 0 | j^0(0) | E, \mathbf{k}, \xi \rangle \langle E, \mathbf{k}, \xi | \phi(0) | 0 \rangle$$

$$I_+ - I_- = v$$

$$\rightarrow (2\pi)^3 \int dE e^{-\frac{i}{\hbar}E\Delta t} [\sigma_+(\mathbf{0}, E) - \sigma_-(\mathbf{0}, E)] = v$$

$$(2\pi)^4 [\sigma_+(\mathbf{0}, E) - \sigma_-(\mathbf{0}, E)] = v \int \frac{d\Delta t}{\hbar} e^{\frac{i}{\hbar}E\Delta t}$$

$$\sigma_+(\mathbf{0}, E) - \sigma_-(\mathbf{0}, E) = \frac{v}{(2\pi)^3} \delta(E)$$

\rightarrow At least one of σ_+ & σ_- contains a $\delta(E)$ factor.

$\therefore \exists$ (gapless) mode with $E=0$, which occurs at $\mathbf{k}=0$.