

5.2.d. Spin of Photons

Consider the Coulomb gauge (see 5.2.c._CanonicalQuantization_LinearMediumCoulombGauge.pdf).

The EM field is a sum of 2 independent K-G fields specified by the polarization vectors $\boldsymbol{\epsilon}^\lambda(\mathbf{k})$ with $\mathbf{k} \cdot \boldsymbol{\epsilon}^\lambda(\mathbf{k}) = 0$.

The quanta of these fields are called photons.

One can rotate these vectors around \mathbf{k} by an angle θ to get

$$\begin{aligned}\boldsymbol{\epsilon}_\theta^1 &= e^{i\theta J_{\parallel}/\hbar} \boldsymbol{\epsilon}^1 = \boldsymbol{\epsilon}^1 \cos\theta + \boldsymbol{\epsilon}^2 \sin\theta \\ \boldsymbol{\epsilon}_\theta^2 &= e^{i\theta J_{\parallel}/\hbar} \boldsymbol{\epsilon}^2 = -\boldsymbol{\epsilon}^1 \sin\theta + \boldsymbol{\epsilon}^2 \cos\theta\end{aligned}$$

where $J_{\parallel} = \mathbf{J} \cdot \hat{\mathbf{k}}$ is the component of the internal angular momentum (spin) \mathbf{J} in the direction of \mathbf{k} .

Eigenstates of $e^{i\theta J_{\parallel}/\hbar}$ are

$$\boldsymbol{\epsilon}^\pm = \frac{1}{2} (\boldsymbol{\epsilon}^1 \mp i \boldsymbol{\epsilon}^2)$$

since

$$\begin{aligned}e^{i\theta J_{\parallel}/\hbar} \boldsymbol{\epsilon}^\pm &= \frac{1}{2} [\boldsymbol{\epsilon}^1 \cos\theta + \boldsymbol{\epsilon}^2 \sin\theta \mp i (-\boldsymbol{\epsilon}^1 \sin\theta + \boldsymbol{\epsilon}^2 \cos\theta)] \\ &= \frac{1}{2} [\boldsymbol{\epsilon}^1 (\cos\theta \pm i \sin\theta) \mp i \boldsymbol{\epsilon}^2 (\cos\theta \pm i \sin\theta)] \\ &= e^{\pm i\theta} \boldsymbol{\epsilon}^\pm\end{aligned}$$

For $\theta \ll 1$,

$$(1 + i\theta J_{\parallel}/\hbar + \dots) \boldsymbol{\epsilon}^\pm = (1 \pm i\theta + \dots) \boldsymbol{\epsilon}^\pm$$

$$\rightarrow J_{\parallel} \boldsymbol{\epsilon}^\pm = \pm \hbar \boldsymbol{\epsilon}^\pm$$

i.e., $\boldsymbol{\epsilon}^\pm$ are also the eigenstates of J_{\parallel} with \pm helicity.