

6.3.a. Fermi Sea

The interpretation of a negative energy particle state as an anti-particle state of positive energy was already discussed in Chap. 3 for the K-G (bosonic) fields. The same procedure is also applicable to the Dirac field.

Historically, before the concept of anti-particles, Dirac uses the Fermi sea to justify such an interpretation. The hole concept was subsequently abandoned by particle physicists as being unnecessarily complicated while the anti-particles themselves were experimentally observed. However, in semi-conductor physics, the hole concept is the preferred one since it is closer to the actual physics.

Retracing the steps followed in 3.3._RealKlein-GordonField.pdf, we begin with

$$\psi(x) = \sum_{\sigma} \int \frac{d^3 k}{(2\pi)^{3/2}} \alpha_k \left[c_{\sigma}^{(+)}(\mathbf{k}) u_{\sigma}^{(+)}(\mathbf{k}) e^{-i\omega_k t + i\mathbf{k}\cdot\mathbf{r}} + c_{\sigma}^{(-)}(\mathbf{k}) u_{\sigma}^{(-)}(\mathbf{k}) e^{i\omega_k t + i\mathbf{k}\cdot\mathbf{r}} \right]$$

where
$$\alpha_k = \sqrt{\frac{m^* c}{\omega_k}} = \frac{m c^2}{\hbar \omega_k}$$

& the superscript +/ - denotes the sign of the energy of the plane wave.

Note that both $c_{\sigma}^{(\pm)}(\mathbf{k})$ are annihilation operators of plane waves of energies $\pm \hbar \omega_k$ but momenta $+\hbar \mathbf{k}$.

Thus, after the usual rigmarole (c.f. 6.3._CanonicalQuantization.pdf), we should get

$$H = \sum_{\sigma} \int d^3 k \hbar \omega_k \left[c_{\sigma}^{(+)}(\mathbf{k}) c_{\sigma}^{(+)}(\mathbf{k}) - c_{\sigma}^{(-)}(\mathbf{k}) c_{\sigma}^{(-)}(\mathbf{k}) \right]$$

$$P = \sum_{\sigma} \int d^3 k \hbar \mathbf{k} \left[c_{\sigma}^{(+)}(\mathbf{k}) c_{\sigma}^{(+)}(\mathbf{k}) + c_{\sigma}^{(-)}(\mathbf{k}) c_{\sigma}^{(-)}(\mathbf{k}) \right]$$

Taking $\mathbf{k} \rightarrow -\mathbf{k}$, the (-) term in ψ becomes

$$\sum_{\sigma} \int \frac{d^3 k}{(2\pi)^{3/2}} \alpha_k c_{\sigma}^{(-)}(-\mathbf{k}) u_{\sigma}^{(-)}(-\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}}$$

Thus, setting

$$d_{\sigma}^{+}(\mathbf{k}) = c_{-\sigma}^{(-)}(-\mathbf{k}) \quad v_{\sigma}(\mathbf{k}) = u_{-\sigma}^{(-)}(-\mathbf{k})$$

where $\sigma = \pm 1$ is used to denote spin \uparrow / \downarrow , turns it into

$$\sum_{\sigma} \int \frac{d^3 k}{(2\pi)^{3/2}} \alpha_k d_{-\sigma}^{+}(\mathbf{k}) v_{-\sigma}(\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}} = \sum_{\sigma} \int \frac{d^3 k}{(2\pi)^{3/2}} \alpha_k d_{\sigma}^{+}(\mathbf{k}) v_{\sigma}(\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}}$$

the integrand of which describes creating a plane wave of positive energy $\hbar \omega_k$ & momentum $\hbar \mathbf{k}$.

The Fermi (or Dirac) sea is a Dirac field with all the negative energy states filled.

This is possible since according to the Pauli exclusion principle, each state can accommodate just 1 particle. Since for each state with momentum \mathbf{p} , there is a state with $-\mathbf{p}$, & likewise for the spins, we have

$$\mathbf{P} = 0 \quad \& \quad \mathbf{S} = 0$$

On the other hand,

$$E = E_0 = -\infty$$

Thus, when $c_{-\sigma}^{(-)}(-\mathbf{k})$ annihilates a particle in the state of momentum $-\hbar \mathbf{k}$, spin $-\sigma$ & energy $-\hbar \omega_k$ in the Fermi sea, a hole is created while the system is left with

$$\mathbf{P} = 0 - (-\hbar \mathbf{k}) = \hbar \mathbf{k} \quad , \quad \mathbf{S} = 0 - \left(-\frac{1}{2} \sigma \right) = \frac{1}{2} \sigma$$

&
$$E - E_0 = -(-\hbar \omega_k) = \hbar \omega_k$$

One can assign these attributes either to the hole or to an anti-particle so that $d_{\sigma}^{+}(\mathbf{k})$ creates. The advantage of the anti-particle interpretation is that we can invoke the renormalization scheme & drop the 'zero-point energy' E_0 . In doing so, we've also replaced the Fermi sea with the vacuum.