

7.6.c. Skyrmions

Consider the $O(N)$ nonlinear sigma model (see 4.4._SigmaModel.pdf)

$$\begin{aligned} \mathcal{L} &= \frac{1}{2} f \partial_\mu \phi \cdot \partial^\mu \phi & (\text{Ezawa: } f = J_s \text{ \& } \eta_{\mu\nu} = -g_{\mu\nu}^{\text{Ezawa}}) \\ &= \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 - \partial_i \phi \cdot \partial_i \phi \right) \end{aligned}$$

where

$$\phi = (\phi_1, \dots, \phi_N)$$

with

$$\sum_{a=1}^N \phi_a^2 = \phi \cdot \phi = \phi^2 = 1$$

$$\begin{aligned} \therefore \mathcal{H} &= \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \cdot \dot{\phi} - \mathcal{L} \\ &= \frac{1}{2} f \left(\frac{1}{c^2} \dot{\phi}^2 + \partial_i \phi \cdot \partial_i \phi \right) \end{aligned}$$

For static fields

$$H = \frac{1}{2} f \int d^m r \partial_i \phi \cdot \partial_i \phi$$

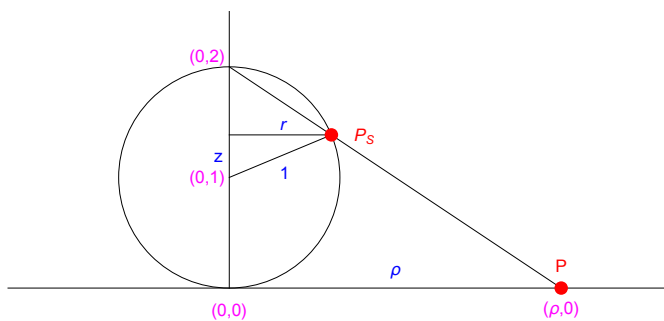
The integral is finite only if the integrand goes to 0 as $r \rightarrow \infty$.

This is guaranteed if

$$\lim_{r \rightarrow \infty} \phi(r) = \phi^\infty$$

where ϕ^∞ is a constant vector. Thus, to each ϕ_a , all points at infinity are the same, i.e., the configuration space \mathbb{R}^m is compactified into a spherical surface S^m .

For example, the xy plane \mathbb{R}^2 is compactified into the unit spherical surface S^2 so that all points at infinity are mapped to a single point, say, the north pole, on S^2 . (This is analogous to the use of the Riemann sphere to compactify the complex plane.)



With reference to the graph above, a point $P = (\rho, 0)$ on the xy -plane is mapped to a point

$P_S = (r, 1+z)$ on the sphere. Using

$$\frac{\rho}{2} = \frac{r}{1-z} \quad \& \quad z^2 + r^2 = 1$$

we get

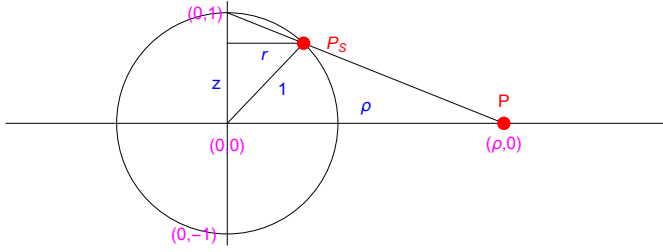
$$r = \sqrt{1-z^2} \quad \rightarrow \quad \frac{\rho}{2} = \sqrt{\frac{1+z}{1-z}}$$

$$\begin{aligned} \therefore z &= \frac{\rho^2 - 4}{\rho^2 + 4} \\ \rightarrow 1 + z &= \frac{2\rho^2}{\rho^2 + 4} \quad \& \quad r = \frac{4\rho}{\rho^2 + 4} \end{aligned}$$

i.e., the mapping is

$$(\rho, 0) \rightarrow \left(\sqrt{1 - z^2}, 1 + z \right) = \frac{2\rho}{\rho^2 + 4} (2, \rho)$$

Alternatively, we can place the sphere center at (0, 0) so that points with $\rho > 1$ is mapped to the upper hemisphere ($z > 0$), & points with $\rho < 1$ to the lower one ($z < 0$).



Setting $P_S = (r, z)$, we have

$$\rho = \frac{r}{1 - z} \quad \& \quad z^2 + r^2 = 1$$

so that

$$r = \sqrt{1 - z^2} \quad \rightarrow \quad \rho = \sqrt{\frac{1 + z}{1 - z}}$$

$$\therefore z = \frac{\rho^2 - 1}{\rho^2 + 1} \quad \& \quad r = \frac{2\rho}{\rho^2 + 1}$$

i.e., the mapping is

$$(\rho, 0) \rightarrow \left(\sqrt{1 - z^2}, z \right) = \frac{1}{\rho^2 + 1} (2\rho, \rho^2 - 1)$$

To get Ezawa's results (with allowance for error), set $\rho = r_{\text{Ezawa}}$.

Now, ϕ^∞ is a vector on S^{N-1} . The boundary of an m -D system is S^m .

The boundary condition $\lim_{r \rightarrow \infty} \phi(r) = \phi^\infty$ thus maps S^m to S^{N-1} .

The homotopy group is therefore $\pi_m(S^{N-1})$.

If it's nontrivial, i.e., $\pi_m(S^{N-1}) \neq 0$, there'll be topological solitons called skyrmions.

Note that the argument also apply to the anisotropic nonlinear sigma model

$$\mathcal{L} = \frac{1}{2} f \sum_a c_a \partial_\mu \phi_a \cdot \partial^\mu \phi_a$$