

### 10.3.a. Angular Momentum

From 10.2.\_CyclotronMotion.pdf ,

$$a = \sqrt{\frac{M}{2\hbar\omega_c}} (v_1 + i\mathbf{e}v_2)$$

$$[a, a^\dagger] = 1$$

$$\mathbf{R} = \mathbf{r} - \frac{e}{\omega_c} \hat{\mathbf{z}} \times \mathbf{v}$$

$$\mathbf{r}^\dagger = \mathbf{r} \quad \mathbf{v}^\dagger = \mathbf{v}$$

$$\rightarrow \mathbf{R}^\dagger = \mathbf{R}$$

In order to find  $b$  such that

$$[b, b^\dagger] = 1$$

we set

$$b = \alpha (R_1 - i\mathbf{e}R_2) \quad \rightarrow \quad b^\dagger = \alpha (R_1 + i\mathbf{e}R_2)$$

$$[R_i, R_j] = -i\mathbf{e}l_B^2 \varepsilon_{ij3}$$

$$\begin{aligned} \rightarrow [b, b^\dagger] &= \alpha^2 [R_1 - i\mathbf{e}R_2, R_1 + i\mathbf{e}R_2] \\ &= i\mathbf{e}\alpha^2 ([R_1, R_2] - [R_2, R_1]) \\ &= 2\alpha^2 l_B^2 \\ &\equiv 1 \end{aligned}$$

$$\rightarrow \alpha = \frac{1}{\sqrt{2} l_B}$$

$$b = \frac{1}{\sqrt{2} l_B} (R_1 - i\mathbf{e}R_2)$$

$$b^\dagger = \frac{1}{\sqrt{2} l_B} (R_1 + i\mathbf{e}R_2)$$

$$bb^\dagger = \frac{1}{2l_B^2} (R_1^2 + R_2^2 + i\mathbf{e}[R_1, R_2])$$

$$b^\dagger b = \frac{1}{2l_B^2} (R_1^2 + R_2^2 - i\mathbf{e}[R_1, R_2])$$

$$\rightarrow bb^\dagger + b^\dagger b = \frac{1}{l_B^2} \mathbf{R}^2 = 2b^\dagger b + 1$$

$$[v_i, R_j] = 0$$

$$\rightarrow [a, b] = [a, b^\dagger] = 0$$

$$H = \hbar\omega_c \left( a^\dagger a + \frac{1}{2} \right)$$

$$\rightarrow \dot{b} = \frac{1}{i\hbar} [b, H] = 0 = \dot{b}^\dagger$$

Thus, the stationary states of the system can be written as  $|n, m\rangle$  with

$$a^\dagger a |n, m\rangle = n |n, m\rangle \quad \& \quad b^\dagger b |n, m\rangle = m |n, m\rangle$$

where  $n$  is the label of the Landau level. The meaning of  $m$  is discussed in 10.3.\_Symmetric-Gauge.pdf.

$$\begin{aligned}
\mathbf{R}^2 &= \left( r_i + \frac{e}{\omega_c} \varepsilon_{ij3} v_j \right) \left( r_i + \frac{e}{\omega_c} \varepsilon_{ik3} v_k \right) \\
&= r^2 + \frac{1}{\omega_c^2} \varepsilon_{ij3} \varepsilon_{ik3} v_j v_k + \frac{e}{\omega_c} (\varepsilon_{ik3} r_i v_k + \varepsilon_{ij3} v_j r_i) \\
&= r^2 + \frac{1}{\omega_c^2} \mathbf{v}^2 + \frac{e}{\omega_c} (\varepsilon_{ik3} r_i v_k + \varepsilon_{ij3} v_j r_i) \\
[r_i, v_j] &= i \frac{\hbar}{M} \delta_{ij} \\
\rightarrow \quad \varepsilon_{ij3} v_j r_i &= \varepsilon_{ij3} r_i v_j \\
\therefore \quad \mathbf{R}^2 &= r^2 + \frac{1}{\omega_c^2} \mathbf{v}^2 + \frac{2e}{\omega_c} (\mathbf{r} \times \mathbf{v})_3
\end{aligned}$$

The angular momentum about the gliding center is

$$\begin{aligned}
\mathcal{L}_3 &= M (\mathbf{X} \times \mathbf{v})_3 = M \varepsilon_{ij3} X_i v_j \\
X_i &= -\frac{e}{\omega_c} \varepsilon_{ij3} v_j \\
\rightarrow \quad \mathcal{L}_3 &= -\frac{Me}{\omega_c} \varepsilon_{ij3} \varepsilon_{ik3} v_k v_j = -\frac{Me}{\omega_c} \mathbf{v}^2 = -2 \frac{H e}{\omega_c} = -e \hbar (2 a^+ a + 1) \\
\therefore \quad H &= -e \frac{1}{2} \omega_c \mathcal{L}_3 = \frac{1}{2} \omega_c \left| \mathcal{L}_3 \right|
\end{aligned}$$

i.e., the energy is entirely rotational, as expected.

The angular momentum with respect to the origin is defined as

$$\begin{aligned}
L_3 &= M (\mathbf{r} \times \mathbf{v})_3 = \frac{1}{2} e M \left( \omega_c \mathbf{R}^2 - \frac{1}{\omega_c} \mathbf{v}^2 - \omega_c r^2 \right) \\
&= \frac{1}{2} \left( \frac{eB}{c} \mathbf{R}^2 - \frac{cM^2}{eB} \mathbf{v}^2 - \frac{eB}{c} r^2 \right) \quad \left( l_B^2 = \frac{\hbar c}{|e| B} \right) \\
&= \frac{1}{2} \left( e \hbar (2 b^+ b + 1) - e \hbar (2 a^+ a + 1) - \frac{eB}{c} r^2 \right) \\
&= e \hbar (b^+ b - a^+ a) - \frac{eB}{2c} r^2 \\
&= e \left( \hbar (b^+ b - a^+ a) - \frac{1}{2} M \omega_c r^2 \right)
\end{aligned}$$

In the symmetric gauge (see 10.3.\_SymmetricGauge.pdf), it reduces to

$$L_3 = e \hbar (b^+ b - a^+ a)$$