

2.B. State Variables and Exact Differentials

Read Reichl's Chapter 1, §2.A, and the opening paragraphs of §2.B.

For convenience, definitions of some basic terms are listed below.

Equilibrium state \equiv Time-independent macroscopic state characterized by a few **state variables**.

Extensive variables are state variables that are proportional to the volume V or particle number N of the system. They take the role of generalized displacements in thermodynamics and will be denoted by $\mathbf{X} = \{X_j\}$.

Intensive variables are state variables that are independent of V or N . They take the role of generalized forces in thermodynamics and will be denoted by $\mathbf{Y} = \{Y_j\}$.

Intensive & extensive variables form conjugate pairs so that the **work** done by the system takes the form

$$\bar{d}W = -\mathbf{Y} \cdot \bar{d}\mathbf{X} = -\sum_j Y_j \bar{d}X_j \quad (2.1a)$$

where \bar{d} is an inexact differential [to be defined later].

(**Equilibrium**) **thermodynamics** studies the dynamics of a system on a time-scale that is long enough so that the system can be described by an equilibrium state at every "instance".

Each equilibrium state of the system is denoted by a point in the **phase space** $\{\mathbf{X}, \mathbf{Y}\}$ and the evolution of the system by a path in it. **Inexact differentials** are differentials whose integration is path dependent.

Consider the **differential form** ϕ of two variables

$$\phi = c_1(x_1, x_2) dx_1 + c_2(x_1, x_2) dx_2 \quad (2.1b)$$

where c_1 & c_2 are some functions of x_1 & x_2 .

Let

$$\left(\frac{\partial f}{\partial x}\right)_y = \text{derivative of } f \text{ with respect to } x \text{ while } y \text{ is held constant.} \quad (2.1c)$$

If

$$\left(\frac{\partial c_1}{\partial x_2}\right)_{x_1} = \left(\frac{\partial c_2}{\partial x_1}\right)_{x_2} \quad (2.1d)$$

then there exists a twice differentiable function $F(x_1, x_2)$ such that

$$c_1 = \left(\frac{\partial F}{\partial x_1}\right)_{x_2} \quad c_2 = \left(\frac{\partial F}{\partial x_2}\right)_{x_1} \quad (2.1e)$$

ϕ is then an **exact differential** commonly denoted by dF so that (2.1b) becomes

$$dF = \left(\frac{\partial F}{\partial x_1}\right)_{x_2} dx_1 + \left(\frac{\partial F}{\partial x_2}\right)_{x_1} dx_2 \quad (2.1)$$

while (2.1d) now reads

$$\left(\frac{\partial}{\partial x_2} \left(\frac{\partial F}{\partial x_1}\right)_{x_2}\right)_{x_1} = \left(\frac{\partial}{\partial x_1} \left(\frac{\partial F}{\partial x_2}\right)_{x_1}\right)_{x_2} \quad (2.2)$$

By (2.1), the differential of any given function is always exact.

On the other hand, if (2.1d) is not satisfied, ϕ is called an **inexact differential** and shall be denoted by $\bar{d}F$ so that (2.1b) becomes

$$\overline{dF} = \left(\frac{\partial F}{\partial x_1} \right)_{x_2} dx_1 + \left(\frac{\partial F}{\partial x_2} \right)_{x_1} dx_2 \tag{2.2a}$$

The violation of (2.1d) means that the entity F obtained by integrating (2.2a) is path dependent. Hence, F cannot be a function, even though we shall continue to interpret $\left(\frac{\partial F}{\partial x_j} \right)_{x_k}$ via (2.1c).

In contrast, the significance of (2.2) is that it guarantees that dF is independent of directions. In other words, dF is a scalar and (2.1) can be written as

$$dF = \nabla F \cdot d\mathbf{r}$$

where ∇F and $d\mathbf{r}$ are vectors in the (x_1, x_2) space with components

$$\nabla F = \left(\left(\frac{\partial F}{\partial x_1} \right)_{x_2}, \left(\frac{\partial F}{\partial x_2} \right)_{x_1} \right) \quad d\mathbf{r} = (dx_1, dx_2)$$

Thus, the line integral

$$\int_A^B dF = F(B) - F(A) \tag{2.2a}$$

is path independent and vanishes for any closed path

$$\oint dF = 0 \tag{2.2b}$$

Given dF , one can always integrate to get the function (or scalar field) F to within an additive constant.

The foregoing can easily be generalized to the case of n variables $\mathbf{x} = \{x_1, \dots, x_n\}$. Thus, the differential form

$$\phi = \sum_{j=1}^n c_j dx_j \tag{2.3}$$

is exact if and only if

$$\left(\frac{\partial}{\partial x_j} \left(\frac{\partial F}{\partial x_k} \right)_{\{x_m \neq x_k\}} \right)_{\{x_m \neq x_j\}} = \left(\frac{\partial}{\partial x_k} \left(\frac{\partial F}{\partial x_j} \right)_{\{x_m \neq x_j\}} \right)_{\{x_m \neq x_k\}} \tag{2.4}$$

for all $j, k = 1, \dots, n$.

An example for $n = 3$ is

$$\phi = c_1 dx + c_2 dy + c_3 dz$$

which is exact if and only if

$$\left(\frac{\partial c_1}{\partial y} \right)_{x,z} = \left(\frac{\partial c_2}{\partial x} \right)_{y,z} \quad \left(\frac{\partial c_1}{\partial z} \right)_{x,y} = \left(\frac{\partial c_3}{\partial x} \right)_{z,y} \quad \left(\frac{\partial c_2}{\partial z} \right)_{y,x} = \left(\frac{\partial c_3}{\partial y} \right)_{z,x}$$

By definition, the differential of any state variable is exact.

Let x, y & z be 3 state variables satisfying

$$F(x, y, z) = 0 \tag{2.5a}$$

which means only 2 of them are independent. Also, let w be a function of any two of x, y & z .

We shall prove the following useful identities known as the **Maxwell relations**.

$$\left(\frac{\partial x}{\partial y} \right)_z = \frac{1}{\left(\frac{\partial y}{\partial x} \right)_z} \tag{2.5}$$

$$\left(\frac{\partial x}{\partial y}\right)_z \left(\frac{\partial y}{\partial z}\right)_x \left(\frac{\partial z}{\partial x}\right)_y = -1 \quad (2.6)$$

$$\left(\frac{\partial x}{\partial w}\right)_z = \left(\frac{\partial x}{\partial y}\right)_z \left(\frac{\partial y}{\partial w}\right)_z \quad (2.7)$$

$$\left(\frac{\partial x}{\partial y}\right)_z = \left(\frac{\partial x}{\partial y}\right)_w + \left(\frac{\partial x}{\partial w}\right)_y \left(\frac{\partial w}{\partial y}\right)_z \quad (2.8)$$

To begin, (2.5a) implies $x = x(y, z)$ and $y = y(x, z)$. Hence,

$$\left(\frac{\partial x}{\partial y}\right)_z = \frac{(dx)_z}{(dy)_z} = \left[\frac{(dy)_z}{(dx)_z}\right]^{-1} = \left[\left(\frac{\partial y}{\partial x}\right)_z\right]^{-1}$$

which proves (2.5).

(2.5a) also implies $z = z(x, y)$ so that

$$dz = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x dy$$

Keeping z constant then gives

$$0 = \left(\frac{\partial z}{\partial x}\right)_y (dx)_z + \left(\frac{\partial z}{\partial y}\right)_x (dy)_z$$

$$\rightarrow \left(\frac{\partial x}{\partial y}\right)_z = \frac{(dx)_z}{(dy)_z} = -\frac{\left(\frac{\partial z}{\partial y}\right)_x}{\left(\frac{\partial z}{\partial x}\right)_y}$$

$$\therefore \left(\frac{\partial x}{\partial y}\right)_z \left[\left(\frac{\partial z}{\partial y}\right)_x\right]^{-1} \left(\frac{\partial z}{\partial x}\right)_y = -1$$

which is just (2.6) since

$$\left[\left(\frac{\partial z}{\partial y}\right)_x\right]^{-1} = \left(\frac{\partial y}{\partial z}\right)_x \quad [(2.5) \text{ used. }]$$

From $x = x(y, z)$, we have

$$dx = \left(\frac{\partial x}{\partial y}\right)_z dy + \left(\frac{\partial x}{\partial z}\right)_y dz$$

$$\rightarrow \frac{dx}{dw} = \left(\frac{\partial x}{\partial y}\right)_z \frac{dy}{dw} + \left(\frac{\partial x}{\partial z}\right)_y \frac{dz}{dw} \quad (2.8a)$$

where dw is any exact differential. Keeping z constant and setting $w = w(x, y)$, (2.8a) becomes

$$\left(\frac{\partial x}{\partial w}\right)_z = \left(\frac{\partial x}{\partial y}\right)_z \left(\frac{\partial y}{\partial w}\right)_z$$

which proves (2.7).

Finally, consider $x = x(y, w)$ so that

$$dx = \left(\frac{\partial x}{\partial y}\right)_w dy + \left(\frac{\partial x}{\partial w}\right)_y dw \quad (2.8b)$$

Setting $w = w(y, z)$ gives

$$dw = \left(\frac{\partial w}{\partial y}\right)_z dy + \left(\frac{\partial w}{\partial z}\right)_y dz$$

so that (2.8b) becomes

$$dx = \left(\frac{\partial x}{\partial y}\right)_w dy + \left(\frac{\partial x}{\partial w}\right)_y \left[\left(\frac{\partial w}{\partial y}\right)_z dy + \left(\frac{\partial w}{\partial z}\right)_y dz \right]$$

Setting $z = \text{const}$ then gives

$$(dx)_z = \left(\frac{\partial x}{\partial y}\right)_w (dy)_z + \left(\frac{\partial x}{\partial w}\right)_y \left(\frac{\partial w}{\partial y}\right)_z (dy)_z$$

$$\rightarrow \left(\frac{\partial x}{\partial y}\right)_z = \left(\frac{\partial x}{\partial y}\right)_w + \left(\frac{\partial x}{\partial w}\right)_y \left(\frac{\partial w}{\partial y}\right)_z$$

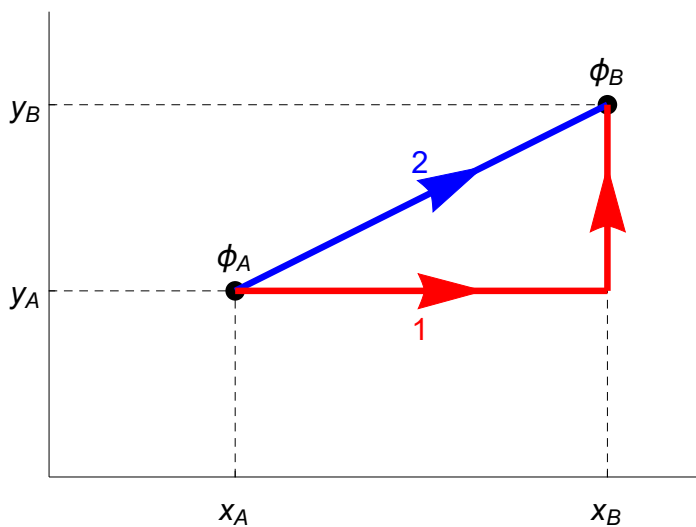
which proves (2.8)

Ex.2.1.

Consider the differential

$$d\phi = (x^2 + y) dx + x dy \quad (1a)$$

- Show that it is an exact differential.
- Integrate $d\phi$ between the points A & B in the figure below, along the paths, 1 & 2.
- Integrate $d\phi$ between the points A & B using indefinite integrals.



Answer (a)

Writing (1a) as

$$d\phi = \left(\frac{\partial \phi}{\partial x}\right)_y dx + \left(\frac{\partial \phi}{\partial y}\right)_x dy$$

we have

$$\left[\frac{\partial}{\partial y} \left(\frac{\partial \phi}{\partial x} \right)_y \right]_x = \left(\frac{\partial (x^2 + y)}{\partial y} \right)_x = 1$$

$$\left[\frac{\partial}{\partial x} \left(\frac{\partial \phi}{\partial y} \right)_x \right]_y = \left(\frac{\partial x}{\partial x} \right)_y = 1$$

Hence, the criterion (2.4) is satisfied and $d\phi$ is exact.

Answer (b)

Along path 1,

$$\begin{aligned}
 \int_A^B d\phi &= \int_{x_A}^{x_B} dx \left(\frac{\partial \phi}{\partial x} \right)_y + \int_{y_A}^{y_B} dy \left(\frac{\partial \phi}{\partial y} \right)_x \\
 &= \int_{x_A}^{x_B} dx (x^2 + y)_{y=y_A} + \int_{y_A}^{y_B} dy (x)_{x=x_B} \\
 &= \int_{x_A}^{x_B} dx (x^2 + y_A) + \int_{y_A}^{y_B} dy x_B \\
 &= \frac{1}{3} (x_B^3 - x_A^3) + (x_B - x_A) y_A + (y_B - y_A) x_B \\
 &= \frac{1}{3} (x_B^3 - x_A^3) - x_A y_A + y_B x_B
 \end{aligned} \tag{1}$$

Along path 2,

$$y = y_A + \frac{y_B - y_A}{x_B - x_A} (x - x_A)$$

so that

$$dy = \frac{y_B - y_A}{x_B - x_A} dx$$

and (1a) becomes

$$\begin{aligned}
 d\phi &= \left[x^2 + y_A + \frac{y_B - y_A}{x_B - x_A} (x - x_A) \right] dx + x \frac{y_B - y_A}{x_B - x_A} dx \\
 &= \left[x^2 + y_A + \frac{y_B - y_A}{x_B - x_A} (2x - x_A) \right] dx
 \end{aligned}$$

$$\begin{aligned}
 \rightarrow \int_A^B d\phi &= \int_{x_A}^{x_B} dx \left[x^2 + y_A + \frac{y_B - y_A}{x_B - x_A} (2x - x_A) \right] \\
 &= \frac{1}{3} (x_B^3 - x_A^3) + (x_B - x_A) y_A + \frac{y_B - y_A}{x_B - x_A} [x_B^2 - x_A^2 - (x_B - x_A) x_A] \\
 &= \frac{1}{3} (x_B^3 - x_A^3) + (x_B - x_A) y_A + (y_B - y_A) x_B \\
 &= \frac{1}{3} (x_B^3 - x_A^3) - x_A y_A + y_B x_B
 \end{aligned} \tag{2}$$

which agrees with (1).

Answer (c)

Treating y as a constant, we have

$$\int dx \left(\frac{\partial \phi}{\partial x} \right)_y = \int dx (x^2 + y) = \frac{1}{3} x^3 + xy + K_1(y) \tag{3}$$

where $K_1(y)$ is some function of y .

Treating x as a constant, we have

$$\int dy \left(\frac{\partial \phi}{\partial y} \right)_x = \int dy x = xy + K_2(x) \tag{4}$$

where $K_2(x)$ is some function of x .

(3) & (4) are the same if we set

$$K_1(y) = C \quad K_2(x) = \frac{1}{3}x^3 + C$$

where C is a constant. Hence,

$$\phi = \int d\phi = \frac{1}{3}x^3 + xy + C$$

so that

$$\int_A^B d\phi = \phi_B - \phi_A = \frac{1}{3}(x_B^3 - x_A^3) + x_B y_B - x_A y_A$$

in agreement with (1).

Ex.2.1.a

Consider the differential

$$d\phi = 2(x^2 + y)dx + xdy \quad (5)$$

(a) Show that it is an inexact differential.

(b) Integrate $d\phi$ between the points A & B in the figure above, along the paths, 1 & 2.

Answer (a)

Writing (5) as

$$d\phi = \left(\frac{\partial \phi}{\partial x}\right)_y dx + \left(\frac{\partial \phi}{\partial y}\right)_x dy$$

we have

$$\left[\frac{\partial}{\partial y}\left(\frac{\partial \phi}{\partial x}\right)_y\right]_x = \left(\frac{\partial 2(x^2 + y)}{\partial y}\right)_x = 2$$

$$\left[\frac{\partial}{\partial x}\left(\frac{\partial \phi}{\partial y}\right)_x\right]_y = \left(\frac{\partial x}{\partial x}\right)_y = 1$$

Hence, the criterion (2.4) is not satisfied and $d\phi$ is inexact.

Answer (b)

Along path 1,

$$\begin{aligned} \int_A^B d\phi &= \int_{x_A}^{x_B} dx \left(\frac{\partial \phi}{\partial x}\right)_y + \int_{y_A}^{y_B} dy \left(\frac{\partial \phi}{\partial y}\right)_x \\ &= \int_{x_A}^{x_B} dx 2(x^2 + y)_{y=y_A} + \int_{y_A}^{y_B} dy (x)_{x=x_B} \\ &= \int_{x_A}^{x_B} dx 2(x^2 + y_A) + \int_{y_A}^{y_B} dy x_B \\ &= 2 \left[\frac{1}{3}(x_B^3 - x_A^3) + (x_B - x_A)y_A \right] + (y_B - y_A)x_B \\ &= \frac{2}{3}(x_B^3 - x_A^3) - 2x_A y_A + (y_B + y_A)x_B \end{aligned} \quad (6)$$

Along path 2,

$$y = y_A + \frac{y_B - y_A}{x_B - x_A} (x - x_A)$$

so that

$$dy = \frac{y_B - y_A}{x_B - x_A} dx$$

and (5) becomes

$$\begin{aligned} d\phi &= 2 \left[x^2 + y_A + \frac{y_B - y_A}{x_B - x_A} (x - x_A) \right] dx + x \frac{y_B - y_A}{x_B - x_A} dx \\ &= \left[2x^2 + 2y_A + \frac{y_B - y_A}{x_B - x_A} (3x - 2x_A) \right] dx \end{aligned}$$

$$\begin{aligned} \rightarrow \int_A^B d\phi &= \int_{x_A}^{x_B} dx \left[2x^2 + 2y_A + \frac{y_B - y_A}{x_B - x_A} (3x - 2x_A) \right] \\ &= \frac{2}{3} (x_B^3 - x_A^3) + 2(x_B - x_A)y_A + \frac{y_B - y_A}{x_B - x_A} \left[\frac{3}{2} (x_B^2 - x_A^2) - 2(x_B - x_A)x_A \right] \\ &= \frac{2}{3} (x_B^3 - x_A^3) + 2(x_B - x_A)y_A + \frac{1}{2} (y_B - y_A) (3x_B - x_A) \end{aligned}$$

which is not the same as (6).

Code

```
In[130]:= s = .5; off = {0, .2};
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Graphics[{ {Dashed, Line[{1, 0}, A]}, Line[{3, 0}, F]},
          Line[{0, 1}, A]}, Line[{0, 2}, B]}],
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PointSize[.03], Point[{A, B}],
Thickness[.01], Arrowheads[{0, 0, .1, 0}],
Blue, Arrow[{A, B}], Text["2",  $\frac{1}{2} (A + B) + \text{off}$ ],
Red, Arrow[{A, F}], Arrow[{F, B}], Text["1",  $\frac{1}{2} (A + F) - \text{off}$ ],
},
Axes → True, AxesOrigin → {0, 0},
PlotRange → {{0, 3.5}, {0, 2.5}},
Ticks → {{1, "xA"}, {3, "xB"}, {{1, "yA"}, {2, "yB"}}
```