

2.F.5. Grand Potential

The **grand potential** Ω is defined by the Legendre transform

$$\begin{aligned}\Omega &= U - ST - \sum_j \mu_j' N_j \\ &= \mathbf{X} \cdot \mathbf{Y} \quad [(2.61) \text{ used. }]\end{aligned}\quad (2.122)$$

so that

$$\Omega = \Omega(\xi_\Omega) \quad \text{with} \quad \xi_\Omega = (T, \mathbf{X}, \{\mu_j'\}) \quad (2.122a)$$

Following closely the procedure described in §2.F.2, we have

$$\begin{aligned}d\Omega &= -S dT + \mathbf{Y} \cdot d\mathbf{X} - \sum_j N_j d\mu_j' \\ &= \Psi_\Omega \cdot d\xi_\Omega \quad \text{with} \quad \Psi_\Omega = (-S, \mathbf{Y}, \{-N_j\}) \\ &= \frac{\partial \Omega}{\partial \xi_\Omega} \cdot d\xi_\Omega\end{aligned}\quad (2.123)$$

Thus,

$$\Psi_\Omega = \frac{\partial \Omega}{\partial \xi_\Omega} \quad (2.122b)$$

or, more specifically,

$$-S = \left(\frac{\partial \Omega}{\partial T} \right)_{\mathbf{X}, \{\mu_j'\}} \quad (2.124)$$

$$\mathbf{Y} = \left(\frac{\partial \Omega}{\partial \mathbf{X}} \right)_{T, \{\mu_j'\}} \quad \text{or} \quad Y_k = \left(\frac{\partial \Omega}{\partial X_k} \right)_{T, \{X_{i \neq k}\}, \{\mu_j'\}} \quad (2.125)$$

$$N_j = - \left(\frac{\partial \Omega}{\partial \mu_j'} \right)_{T, \mathbf{X}, \{\mu_{k \neq j}'\}} \quad (2.126)$$

The Maxwell equations are

$$\frac{\partial \Psi_{\Omega k}}{\partial \xi_{\Omega j}} = \frac{\partial \Psi_{\Omega j}}{\partial \xi_{\Omega k}} \quad (2.126a)$$

or

$$\begin{aligned}\frac{\partial}{\partial \mathbf{X}} \left(\frac{\partial \Omega}{\partial T} \right) &= \frac{\partial}{\partial T} \left(\frac{\partial \Omega}{\partial \mathbf{X}} \right) \\ \rightarrow - \left(\frac{\partial S}{\partial X_k} \right)_{T, \{X_{i \neq k}\}, \{\mu_j'\}} &= \left(\frac{\partial Y_k}{\partial T} \right)_{\mathbf{X}, \{\mu_j'\}}\end{aligned}\quad (2.127)$$

$$\begin{aligned}\frac{\partial}{\partial \mu_k'} \left(\frac{\partial \Omega}{\partial T} \right) &= \frac{\partial}{\partial T} \left(\frac{\partial \Omega}{\partial \mu_k'} \right) \\ \rightarrow - \left(\frac{\partial S}{\partial \mu_k'} \right)_{T, \mathbf{X}, \{\mu_{j \neq k}'\}} &= - \left(\frac{\partial N_k}{\partial T} \right)_{\mathbf{X}, \{\mu_j'\}}\end{aligned}\quad (2.128)$$

$$\begin{aligned}\frac{\partial}{\partial \mu_k'} \left(\frac{\partial \Omega}{\partial X_i} \right) &= \frac{\partial}{\partial X_i} \left(\frac{\partial \Omega}{\partial \mu_k'} \right) \\ \rightarrow \left(\frac{\partial Y_i}{\partial \mu_k'} \right)_{T, \mathbf{X}, \{\mu_{j \neq k}'\}} &= - \left(\frac{\partial N_k}{\partial X_i} \right)_{T, \{X_{n \neq i}\}, \{\mu_j'\}}\end{aligned}\quad (2.129)$$

$$\begin{aligned} \frac{\partial}{\partial \mu_k'} \left(\frac{\partial \Omega}{\partial \mu_i'} \right) &= \frac{\partial}{\partial \mu_i'} \left(\frac{\partial \Omega}{\partial \mu_k'} \right) \\ \rightarrow - \left(\frac{\partial N_i}{\partial \mu_k'} \right)_{T, \mathcal{X}, \{\mu_{j \neq k}'\}} &= - \left(\frac{\partial N_k}{\partial \mu_i'} \right)_{T, \mathcal{X}, \{\mu_{j \neq i}'\}} \end{aligned} \quad (2.130)$$

(2.83d) becomes

$$d\Omega = -\bar{d}W_{\text{free}} \Big|_{(T, \mathcal{X}, \{\mu_j'\}) = \text{const}} \rightarrow \Delta\Omega = -W_{\text{free}} \Big|_{(T, \mathcal{X}, \{\mu_j'\}) = \text{const}} \quad (2.132)$$

so that at constant $(T, \mathcal{X}, \{\mu_j'\})$, the reversible work $-W_{\text{free}}$ done on the system is converted entirely to $\Delta\Omega$. Hence [c.f. (2.83e)],

$$\Delta\Omega < 0 \quad \text{for all spontaneous processes at constant } (T, \mathcal{X}, \{\mu_j'\}) \quad (2.133)$$

which means the equilibrium state for processes at constant $(T, \mathcal{X}, \{\mu_j'\})$ has the minimal grand potential.