

5.C.1. Spectral Properties

For a time independent \mathbf{Q} , we have

$$Q_{mn}(s) = P_{1|1}(m, s | n, s+1) = P_{1|1}(m, 0 | n, 1) = Q_{mn} \quad \forall s \quad (5.23)$$

or

$$\mathbf{Q}(s) = \mathbf{Q} \quad \forall s \quad (5.23a)$$

(5.21) can be solved by iteration starting with the self-evident "initial condition"

$$P_{1|1}(m, s | n, s) = \delta_{mn} \text{ or } \mathbf{Q}(s | s) = \mathbf{I} \quad (5.23b)$$

For the 1st time step (or iteration), (5.21) becomes

$$\begin{aligned} P_{1|1}(m, s | n, s+1) &= \sum_k P_{1|1}(m, s | k, s) P_{1|1}(k, s | n, s+1) \\ &= \sum_k \delta_{mk} P_{1|1}(k, s | n, s+1) \\ &= P_{1|1}(m, s | n, s+1) \\ &= Q_{mn} \end{aligned}$$

or

$$\mathbf{Q}(s | s+1) = \mathbf{I} \cdot \mathbf{Q}(s | s+1) = \mathbf{Q}$$

For the 2nd time step,

$$\begin{aligned} P_{1|1}(m, s | n, s+2) &= \sum_k P_{1|1}(m, s | k, s+1) P_{1|1}(k, s+1 | n, s+2) \\ &= \sum_k Q_{mk} Q_{kn} \\ &= (\mathbf{Q}^2)_{mn} \end{aligned}$$

or

$$\mathbf{Q}(s | s+2) = \mathbf{Q}(s | s+1) \cdot \mathbf{Q}(s+1 | s+2) = \mathbf{Q}^2$$

For the 3rd time step,

$$\begin{aligned} P_{1|1}(m, s | n, s+3) &= \sum_k P_{1|1}(m, s | k, s+2) P_{1|1}(k, s+2 | n, s+3) \\ &= \sum_k (\mathbf{Q}^2)_{mk} Q_{kn} \\ &= (\mathbf{Q}^3)_{mn} \end{aligned}$$

or

$$\mathbf{Q}(s | s+3) = \mathbf{Q}(s | s+2) \cdot \mathbf{Q}(s+2 | s+3) = \mathbf{Q}^3$$

In general,

$$P_{1|1}(m, s_0 | n, s) = (\mathbf{Q}^{s-s_0})_{mn} \quad (5.24)$$

or

$$\mathbf{Q}(s_0 | s) = \mathbf{Q}^{s-s_0} \quad (5.24a)$$

Similarly, (5.20) gives, by iteration,

$$\begin{aligned} P(n, s+1) &= \sum_m P(m, s) Q_{mn} \\ &= \sum_{m,k} P(k, s-1) Q_{km} Q_{mn} = \sum_k P(k, s-1) (\mathbf{Q}^2)_{kn} \\ &\vdots \\ &= \sum_k P(k, 0) (\mathbf{Q}^{s+1})_{kn} \end{aligned} \quad (5.25)$$

Setting $\mathbf{p}(s)$ to be the column vector with elements $p_n(s) = P(n, s)$, we can write (5.25) as

$$\begin{aligned} \mathbf{p}^T(s+1) &= \mathbf{p}^T(s) \cdot \mathbf{Q} \\ &= \mathbf{p}^T(s-1) \cdot \mathbf{Q}^2 \\ &\vdots \\ &= \mathbf{p}^T(0) \cdot \mathbf{Q}^{s+1} \end{aligned} \tag{5.25a}$$

where T denotes the transpose operation.

As in quantum mechanics, the Dirac notation greatly clarifies the mathematics.

Reminder: in the Heisenberg picture, the Dirac notation is equivalent to the matrix form.

The state of the system at time s is represented by the **probability vector** $|\rho(s)\rangle$ in the sample space \mathbb{S} , which is spanned by the eigenvectors $|n\rangle$ of the stochastic variable operator Y with

$$Y |n\rangle = y(n) |n\rangle \quad n = 1, \dots, M \tag{a1}$$

Since \mathbb{S} is a real space, we define the inner product of two vectors ψ & ϕ in it as

$$\langle \psi | \phi \rangle = \langle \phi | \psi \rangle$$

which means the dual vectors are related by a transpose operation

$$\langle \psi | = | \psi \rangle^T \quad \forall | \psi \rangle \in \mathbb{S} \tag{a2}$$

Using the completeness condition

$$\sum_{n=1}^M |n\rangle \langle n| = I \tag{a3}$$

we can write any $|\psi\rangle \in \mathbb{S}$ as

$$|\psi\rangle = \sum_{n=1}^M |n\rangle \langle n | \psi \rangle = \sum_{n=1}^M |n\rangle \psi(n) \tag{a4}$$

where

$$\psi(n) \equiv \langle n | \psi \rangle \tag{a5}$$

Similarly, the inner product of ψ & ϕ can be written as

$$\begin{aligned} \langle \psi | \phi \rangle &= \sum_{n=1}^M \langle \psi | n \rangle \langle n | \phi \rangle \\ &= \sum_{n=1}^M \psi(n) \phi(n) \end{aligned} \tag{a6}$$

which is sometimes used as the definition of the inner product without introducing the concept of dual space.

Probability $P(n, s)$ is then the projection (or component) of $|\rho(s)\rangle$ along $|n\rangle$:

$$P(n, s) = \langle \rho(s) | n \rangle = \langle n | \rho(s) \rangle \tag{b}$$

From (5.20), we see that the conditional probability $P_{1|1}$ is an operator that relates probability vectors of different times. Hence, we set

$$P_{1|1}(m, s_0 | n, s) = \langle m | P_{1|1}(s_0 | s) | n \rangle \tag{c}$$

so that (5.20) becomes

$$\begin{aligned} \langle \rho(s+1) | n \rangle &= \sum_{m=1}^M \langle \rho(s) | m \rangle \langle m | P_{1|1}(s | s+1) | n \rangle \\ &= \langle \rho(s) | P_{1|1}(s | s+1) | n \rangle \end{aligned} \tag{d}$$

Since (d) is true for all $|n\rangle$, we have

$$\langle \rho(s+1) | = \langle \rho(s) | P_{1|1}(s | s+1) \tag{e}$$

(5.23) then becomes

$$Q_{mn} = \langle m | P_{1|1}(s | s+1) | n \rangle \quad (f)$$

so that \mathbf{Q} is just the matrix form of the operator $P_{1|1}(s | s+1)$. This justifies the use of \mathbf{Q} to denote both the operator $P_{1|1}(s | s+1)$ and its matrix representation. (f) thus becomes

$$Q_{mn} = \langle m | \mathbf{Q} | n \rangle \quad (g)$$

Similarly, (5.24) becomes

$$\begin{aligned} \langle m | P_{1|1}(s_0 | s) | n \rangle &= (\mathbf{Q}^{s-s_0})_{mn} \\ \rightarrow P_{1|1}(s_0 | s) &= P_{1|1}(s_0 | s_0+1) \dots P_{1|1}(s-1 | s) = \mathbf{Q}^{s-s_0} \end{aligned} \quad (h)$$

The eigenvalues of \mathbf{Q} are given by the solutions of the secular equation

$$\det | \mathbf{Q} - \lambda \mathbf{I} | = 0 \quad (5.26)$$

where \mathbf{I} is the $M \times M$ unit matrix. As can be seen from (f), \mathbf{Q} is a real but non-symmetric matrix in the basis $\{ | n \rangle \}$. Therefore, the left and right eigenvectors for the same eigenvalue are not equal.

Thus,

$$\mathbf{Q} | \psi_i \rangle = \lambda_i | \psi_i \rangle \quad (i)$$

$$\langle \chi_i | \mathbf{Q} = \langle \chi_i | \lambda_i \quad (j)$$

where $| \psi_i \rangle$ & $\langle \chi_i |$ are the right & left eigenvectors of λ_i , respectively, and

$$| \chi_i \rangle \neq | \psi_i \rangle$$

Taking the transpose of (j) gives

$$\mathbf{Q}^T | \chi_i \rangle = \lambda_i | \chi_i \rangle \quad (k)$$

which is the preferred form when using computer software for calculation.

Caution: for a real but non-symmetric matrix, both its eigenvalues and eigenvectors can be complex. See comment following (5.31).

With the help of the completeness condition and

$$\psi_i(n) = \langle n | \psi_i \rangle = \langle \psi_i | n \rangle$$

(i) can be written in component form as

$$\begin{aligned} \sum_n \langle m | \mathbf{Q} | n \rangle \langle n | \psi_i \rangle &= \lambda_i \langle m | \psi_i \rangle \\ \rightarrow \sum_n Q_{mn} \psi_i(n) &= \lambda_i \psi_i(m) \end{aligned} \quad (5.28)$$

Similarly, (j) becomes

$$\begin{aligned} \sum_n \langle \chi_i | n \rangle \langle n | \mathbf{Q} | m \rangle &= \langle \chi_i | m \rangle \lambda_i \\ \rightarrow \sum_n \chi_i(n) Q_{nm} &= \chi_i(m) \lambda_i \end{aligned} \quad (5.27)$$

(i) & (j) give

$$\langle \chi_i | \mathbf{Q} | \psi_j \rangle = \lambda_j \langle \chi_i | \psi_j \rangle$$

$$\langle \chi_i | \mathbf{Q} | \psi_j \rangle = \lambda_i \langle \chi_i | \psi_j \rangle$$

$$\rightarrow 0 = (\lambda_j - \lambda_i) \langle \chi_i | \psi_j \rangle$$

$$\therefore \langle \chi_i | \psi_j \rangle = 0 \quad \text{for} \quad \lambda_i \neq \lambda_j$$

i.e., left and right eigenvectors belonging to different eigenvalues are orthogonal.

For eigenvectors belonging to the same eigenvalue, we can use the Schmidt orthogonalization method to obtain a set of mutually orthogonal vectors.

Therefore, assuming all eigenvectors are properly normalized, we can always write

$$\langle \chi_i | \psi_j \rangle = \delta_{ij} \quad (5.30)$$

$$\begin{aligned} \rightarrow \quad & \sum_{n=1}^M \langle \chi_i | n \rangle \langle n | \psi_j \rangle = \delta_{ij} \\ & \sum_{n=1}^M \chi_i(n) \psi_j(n) = \delta_{ij} \end{aligned} \tag{5.30a}$$

Consider the expansion

$$\begin{aligned} \langle p | &= \sum_{i=1}^M \alpha_i \langle \chi_i | \tag{m} \\ \rightarrow \quad \langle p | \psi_j \rangle &= \sum_{i=1}^M \alpha_i \langle \chi_i | \psi_j \rangle = \sum_{i=1}^M \alpha_i \delta_{ij} = \alpha_j \end{aligned}$$

Putting this back into (m) gives

$$\langle p | = \sum_{i=1}^M \langle p | \psi_i \rangle \langle \chi_i |$$

Since this is true for arbitrary $\langle p |$, we must have

$$\sum_{i=1}^M | \psi_i \rangle \langle \chi_i | = I \tag{5.31}$$

which is the **completeness condition** for the eigenvectors.

Note that what (5.31) means is that these possibly complex eigenvectors span an M -D real space that is also spanned by the M real vectors $| n \rangle$. To be consistent, eigenvalues and eigenvectors, if complex, must appear in conjugate pairs. This is so since Q is real.

In terms of components with respect to the basis $| n \rangle$, we have

$$\begin{aligned} & \sum_{i=1}^M \langle m | \psi_i \rangle \langle \chi_i | n \rangle = \langle m | n \rangle \\ \rightarrow \quad & \sum_{i=1}^M \psi_i(m) \chi_i(n) = \delta_{mn} \end{aligned} \tag{5.31a}$$

Taking the absolute value of (5.28) gives

$$\left| \sum_{n=1}^M Q_{mn} \psi_i(n) \right| = | \lambda_i | | \psi_i(m) |$$

Since all Q_{mn} are real and non-negative,

$$| \lambda_i | | \psi_i(m) | \leq \sum_n Q_{mn} | \psi_i(n) | \tag{5.32}$$

Let

$$| \psi_i(n) | \leq C \quad \forall n$$

then

$$\sum_{n=1}^M Q_{mn} | \psi_i(n) | \leq C \sum_{n=1}^M Q_{mn} = C \tag{5.33}$$

where the discrete version of (5.13)

$$\sum_{n=1}^M P_{1|1}(m, s | n, s+1) = \sum_{n=1}^M Q_{mn} = 1 \tag{n}$$

was used.

Let

$$| \psi_i(n_0) | = C$$

then setting $m = n_0$ in (5.32) gives

$$\begin{aligned} & |\lambda_i| \leq C \\ \rightarrow & |\lambda_i| \leq 1 \quad \forall i \end{aligned} \quad (5.34)$$

Using (n), we see that (5.28) can always be satisfied by

$$\lambda_1 = 1 \quad \text{with} \quad \psi_1(n) = c \quad \forall n \quad (o)$$

where c is a normalization constant and we have set arbitrarily $i = 1$.

In other words, every \mathbf{Q} has an eigenvalue of 1 with left eigenvector given by (o).

For the right eigenvector, setting $\lambda_1 = 1$ in (5.27) gives

$$\chi_1(m) = \sum_{n=1}^M \chi_1(n) Q_{nm} \quad (5.35)$$

Setting

$$\chi_1(m) = \langle p(s) | m \rangle = P(m, s)$$

(5.35) becomes

$$\begin{aligned} \langle p(s) | m \rangle &= \sum_{n=1}^M \langle p(s) | n \rangle \langle n | P_{1|1}(s | s+1) | m \rangle \\ &= \langle p(s+1) | m \rangle \quad [(5.20) \text{ used. }] \end{aligned}$$

$$\rightarrow \quad | p(s) \rangle = | p(s+1) \rangle$$

χ_1 is therefore a stationary probability vector. We denote this as

$$\chi_1(n) = P_{\text{ST}}(n) \quad (p)$$

The constant c in (o) can be determined from the normalization

$$\begin{aligned} \langle \chi_1 | \psi_1 \rangle &= 1 \\ \rightarrow \quad \sum_{n=1}^M \chi_1(n) \psi_1(n) &= c \sum_{n=1}^M \chi_1(n) = 1 \end{aligned}$$

Using the sum rule for probabilities

$$\sum_{n=1}^M \chi_1(n) = \sum_{n=1}^M P_{\text{ST}}(n) = 1$$

we have

$$c = 1 \quad \rightarrow \quad \psi_1(n) = 1 \quad \forall n \quad (o1)$$

To summarize, every \mathbf{Q} has an eigenvalue of 1 whose left eigenvector has every component equals to 1, and right eigenvector that is stationary with the sum of its components equal to 1.

On the other hand, $\text{eq}(j) \times | \psi_j \rangle$ gives

$$\begin{aligned} \langle \chi_i | \mathbf{Q} | \psi_j \rangle &= \lambda_i \langle \chi_i | \psi_j \rangle \\ &= \lambda_i \delta_{ij} \quad [(5.30) \text{ used. }] \end{aligned} \quad (q)$$

$$\rightarrow \quad \sum_{i=1}^M \langle \chi_i | \mathbf{Q} | \psi_j \rangle = \lambda_j$$

which implies

$$\mathbf{Q} = \sum_{i=1}^M \lambda_i | \psi_i \rangle \langle \chi_i | \quad (r)$$

Proof: (r) gives

$$\langle \chi_i | \mathbf{Q} | \psi_j \rangle = \sum_{k=1}^M \lambda_k \langle \chi_i | \psi_k \rangle \langle \chi_k | \psi_j \rangle$$

$$= \sum_{k=1}^M \lambda_k \delta_{ik} \delta_{kj} = \lambda_i \delta_{ij}$$

which agrees with (q). QED.

$\langle m | \text{eq}(r) | n \rangle$ gives

$$\begin{aligned} Q_{mn} &= \sum_{i=1}^M \lambda_i \langle m | \psi_i \rangle \langle \chi_i | n \rangle \\ &= \sum_{i=1}^M \lambda_i \psi_i(m) \chi_i(n) \\ &= P_{1|1}(m, s | n, s+1) \quad \forall s \quad [\text{see (5.23).}] \end{aligned} \tag{5.36}$$

From (5.24), we have

$$\begin{aligned} P_{1|1}(m, s | n, s+2) &= (\mathbf{Q}^2)_{mn} \\ &= \sum_{k=1}^M Q_{mk} Q_{kn} \\ &= \sum_{i,j,k=1}^M \lambda_i \lambda_j \langle m | \psi_i \rangle \langle \chi_i | k \rangle \langle k | \psi_j \rangle \langle \chi_j | n \rangle \\ &= \sum_{i,j=1}^M \lambda_i \lambda_j \langle m | \psi_i \rangle \langle \chi_i | \psi_j \rangle \langle \chi_j | n \rangle \\ &= \sum_{i,j=1}^M \lambda_i \lambda_j \langle m | \psi_i \rangle \delta_{ij} \langle \chi_j | n \rangle \\ &= \sum_{i=1}^M \lambda_i^2 \langle m | \psi_i \rangle \langle \chi_i | n \rangle \\ &= \sum_{i=1}^M \lambda_i^2 \psi_i(m) \chi_i(n) \end{aligned}$$

Alternatively,

$$\begin{aligned} (\mathbf{Q}^2)_{mn} &= \langle m | \mathbf{Q}^2 | n \rangle \\ &= \sum_{i=1}^M \langle m | \mathbf{Q}^2 | \psi_i \rangle \langle \chi_i | n \rangle \\ &= \sum_{i=1}^M \langle m | \lambda_i^2 | \psi_i \rangle \langle \chi_i | n \rangle \\ &= \sum_{i=1}^M \lambda_i^2 \psi_i(m) \chi_i(n) \end{aligned}$$

In general,

$$P_{1|1}(m, s_0 | n, s) = (\mathbf{Q}^{s-s_0})_{mn} = \sum_{i=1}^M \lambda_i^{s-s_0} \psi_i(m) \chi_i(n) \tag{5.37}$$

which, together with the stationary state [see (p)]

$$P_{\text{ST}}(n) = \chi_1(n)$$

provide the general solution to the Markov chain.

\mathbf{Q} is called **regular** if

$$\mathbf{Q}^N \neq \mathbf{0} \quad \forall N \tag{5.37a}$$

This means, according to (5.37), any state n can be reached by state m after a time $s - s_0$, provided

there are no other prohibitive factors. The necessary condition for a Markov chain to be **ergodic** is therefore a regular \mathbf{Q} .

A regular \mathbf{Q} can be non-ergodic if it is block-diagonal since states in different blocks cannot be bridged by \mathbf{Q}^N .

Taking the determinant of (5.37a) gives

$$\det \mathbf{Q}^N = (\det \mathbf{Q})^N = \left(\prod_i \lambda_i \right)^N = \prod_i \lambda_i^N \neq 0$$

Hence, the necessary and sufficient condition for \mathbf{Q} to be regular is

$$\lambda_i \neq 0 \quad \forall i \tag{5.37b}$$

Since $|\lambda_i| \leq 1$ with $\lambda_1 = 1$, we write (5.37) as

$$\begin{aligned} P_{1|1}(m, s_0 | n, s) &= \psi_1(m) \chi_1(n) + \sum_{i=2}^M \lambda_i^{s-s_0} \psi_i(m) \chi_i(n) \\ &= P_{ST}(n) + \sum_{i=2}^M \lambda_i^{s-s_0} \psi_i(m) \chi_i(n) \end{aligned} \tag{5.38}$$

where (o), (o1) & (p) were used.

Now,

$$|\lambda_i| < 1 \quad \rightarrow \quad \lim_{s \rightarrow \infty} |\lambda_i|^{s-s_0} = 0 \quad \forall i \geq 2$$

$$\therefore \lim_{s \rightarrow \infty} \lambda_i^{s-s_0} = 0 \quad \forall i \geq 2$$

and (5.38) gives

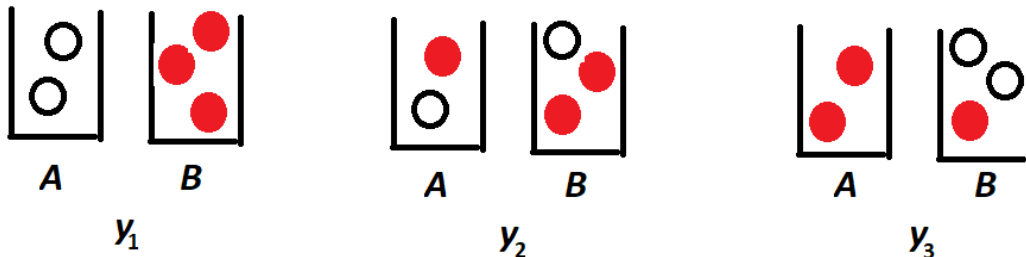
$$\lim_{s \rightarrow \infty} P_{1|1}(m, s_0 | n, s) = P_{ST}(n) \tag{5.39}$$

Similarly, (5.25) gives

$$\begin{aligned} \lim_{s \rightarrow \infty} P(n, s) &= \sum_{m=1}^M P(m, 0) P_{ST}(n) \\ &= P_{ST}(n) \quad [(5.6) \text{ used. }] \end{aligned} \tag{5.40}$$

Exercise 5.1.

Consider 2 pots, A & B , with 3 red & 2 white balls distributed between them so that A & B always have 2 & 3 balls, respectively. There are 3 different configurations as shown in the figure below.



At each time step, transition between configurations is attained by picking up randomly 1 ball each from A & B and interchanging them.

(a) Compute the transition matrix \mathbf{Q} and the conditional probability $P_{1|1}(m, s_0 | n, s)$.

(b) If initially $P(1, 0) = 1$, $P(2, 0) = 0$ and $P(3, 0) = 0$, compute the probabilities $P(n, s)$ for $n = 1, 2, 3$ at time s .

(c) Assume the realization $y(n) = n$. Compute the first moment, $\langle y(s) \rangle$, and the autocorrelation function, $\langle y(0)y(s) \rangle$, for the same initial conditions as (b).

Answer (a)

Starting with $Y(s) = y_1$, one can only reach $Y(s + 1) = y_2$. Therefore

$$Q_{11} = 0 \quad Q_{12} = 1 \quad Q_{13} = 0$$

Starting with $Y(s) = y_2$, one can reach $Y(s + 1) = y_1$ by exchanging a red ball in A with a white ball in

B with a probability $\frac{1}{2} \times \frac{1}{3} = \frac{1}{6}$. Therefore $Q_{21} = \frac{1}{6}$.

There are two ways to remain in y_2 . The first is to pick up a white ball from both pots with probability

$\frac{1}{2} \times \frac{1}{3} = \frac{1}{6}$. The second is to pick a red ball from both pots with probability $\frac{1}{2} \times \frac{2}{3} = \frac{1}{3}$. Therefore,

$$Q_{22} = \frac{1}{6} + \frac{1}{3} = \frac{1}{2}$$

In order to change to y_3 , one must pick a white ball from A and a red ball from B with probability

$\frac{1}{2} \times \frac{2}{3} = \frac{1}{3}$. Therefore $Q_{23} = \frac{1}{3}$.

Starting with $Y(s) = y_3$, there is no way to reach $Y(s + 1) = y_1$. Therefore $Q_{31} = 0$.

One can reach $Y(s + 1) = y_2$ by exchanging a red ball in A with a white ball in B with a probability

$1 \times \frac{2}{3} = \frac{2}{3}$. Therefore $Q_{32} = \frac{2}{3}$.

In order to remain in y_3 , one must pick a red ball from both A and B with probability $1 \times \frac{1}{3} = \frac{1}{3}$. There-

fore $Q_{33} = \frac{1}{3}$.

Thus,

$$Q = \begin{pmatrix} 0 & 1 & 0 \\ \frac{1}{6} & \frac{1}{2} & \frac{1}{3} \\ 0 & \frac{2}{3} & \frac{1}{3} \end{pmatrix} \quad Q^2 = \begin{pmatrix} \frac{1}{6} & \frac{1}{2} & \frac{1}{3} \\ \frac{1}{12} & \frac{23}{36} & \frac{5}{18} \\ \frac{1}{9} & \frac{5}{9} & \frac{1}{3} \end{pmatrix}$$

Using the *Mathematica* code in §Code, we get

$$\begin{aligned} \lambda_1 = 1 & \quad | \psi_1 \rangle = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} & \quad \langle X_1 | = \begin{pmatrix} \frac{1}{3} & 2 & 1 \end{pmatrix} \\ \lambda_2 = -\frac{1}{3} & \quad | \psi_2 \rangle = \begin{pmatrix} 3 \\ -1 \\ 1 \end{pmatrix} & \quad \langle X_2 | = (1 \quad -2 \quad 1) \\ \lambda_3 = \frac{1}{6} & \quad | \psi_3 \rangle = \begin{pmatrix} -\frac{3}{2} \\ -\frac{1}{4} \\ 1 \end{pmatrix} & \quad \langle X_3 | = \begin{pmatrix} -\frac{1}{2} & -\frac{1}{2} & 1 \end{pmatrix} \end{aligned}$$

The eigenvectors above are mutually orthogonal but not normalized. Keeping $| \psi_i \rangle$ as is, the normalized $\langle X_i |$ are easily found to be

$$\langle X_1 | = \frac{3}{10} \begin{pmatrix} \frac{1}{3} & 2 & 1 \end{pmatrix} \quad \langle X_2 | = \frac{1}{6} (1 \quad -2 \quad 1) \quad \langle X_3 | = \frac{8}{15} \begin{pmatrix} -\frac{1}{2} & -\frac{1}{2} & 1 \end{pmatrix}$$

The conditional probability matrix is therefore

$$\begin{aligned}
 P_{1|1}(s_0 | s) &= \sum_{i=1}^3 \lambda_i^{s-s_0} | \psi_i \rangle \langle \chi_i | \\
 &= \begin{pmatrix} \frac{1}{10} & \frac{3}{5} & \frac{3}{10} \\ \frac{1}{10} & \frac{3}{5} & \frac{3}{10} \\ \frac{1}{10} & \frac{3}{5} & \frac{3}{10} \end{pmatrix} + \left(-\frac{1}{3}\right)^{s-s_0} \begin{pmatrix} \frac{1}{2} & -1 & \frac{1}{2} \\ -\frac{1}{6} & \frac{1}{3} & -\frac{1}{6} \\ \frac{1}{6} & -\frac{1}{3} & \frac{1}{6} \end{pmatrix} + \left(\frac{1}{6}\right)^{s-s_0} \begin{pmatrix} \frac{2}{5} & \frac{2}{5} & -\frac{4}{5} \\ \frac{1}{15} & \frac{1}{15} & -\frac{2}{15} \\ -\frac{4}{15} & -\frac{4}{15} & \frac{8}{15} \end{pmatrix}
 \end{aligned}$$

where $| \psi_i \rangle \langle \chi_i |$ was calculated as "rIProd" in the section Code.

Answer (b)

The initial condition can be written as

$$\langle p(0) | = (1 \ 0 \ 0)$$

Using

$$\langle p(s) | = \langle p(0) | P_{1|1}(0 | s)$$

we have [see Code]

$$\langle p(s) | = \left(\frac{1}{10} + \frac{1}{2} \left(-\frac{1}{3}\right)^s + \frac{1}{5} \times 2^{1-s} \times 3^{-s}, \frac{3}{5} + (-1)^{1+s} 3^{-s} + \frac{1}{5} \times 2^{1-s} \times 3^{-s}, \frac{3}{10} + \frac{1}{2} \left(-\frac{1}{3}\right)^s - \frac{1}{5} \times 2^{2-s} \times 3^{-s} \right)$$

Note that

$$\begin{aligned}
 P(n, s) &= \langle p(s) | n \rangle \\
 &= \sum_m \langle p(0) | m \rangle \langle m | P_{1|1}(0 | s) | n \rangle \\
 &= \sum_m P(m, 0) \langle m | P_{1|1}(0 | s) | n \rangle \\
 &= \langle 1 | P_{1|1}(0 | s) | n \rangle
 \end{aligned}$$

Answer (c)

According to (5.7), the 1st moment is

$$\begin{aligned}
 \langle y(s) \rangle &= \sum_{n=1}^3 y(n) P(n, s) \\
 &= \sum_{n=1}^3 n \langle p(s) | n \rangle \\
 &= \frac{11}{5} - \frac{6^{1-s}}{5} \quad [\text{ see Code. }]
 \end{aligned}$$

For the autocorrelation,

$$\begin{aligned}
 \langle y(0) y(s) \rangle &= \sum_{m, n=1}^3 y(m) y(n) P_2(m, 0; n, s) \\
 &= \sum_{m, n=1}^3 y(m) y(n) P(m, 0) P_{1|1}(m, 0 | n, s) \\
 &= \frac{11}{5} - \frac{6^{1-s}}{5} \quad [\text{ see Code. }] \\
 &= \langle y(s) \rangle
 \end{aligned}$$

This coincidence is due to the special initial conditions & realization we used:

$$\begin{aligned}
 \sum_{m,n=1}^3 y(m) y(n) P(m, 0) P_{1|1}(m, 0 | n, s) \\
 &= \sum_{n=1}^3 y(n) P_{1|1}(1, 0 | n, s) \\
 &= \sum_{n=1}^3 n \langle 1 | P_{1|1}(0 | s) | n \rangle \\
 &= \sum_{n=1}^3 n P(n, s) \\
 &= \langle y(s) \rangle
 \end{aligned}$$

Code

$$\mathbf{Q} = \begin{pmatrix} 0 & 1 & 0 \\ \frac{1}{6} & \frac{1}{2} & \frac{1}{3} \\ 0 & \frac{2}{3} & \frac{1}{3} \end{pmatrix};$$

(* eigenvalues & right eigenvectors *)

```
{λ, rev} = Eigensystem[Q]
```

```
{{1, -1/3, 1/6}, {{1, 1, 1}, {3, -1, 1}, {-3/2, -1/4, 1}}}
```

(* Checking *)

```
Table[Q.rev[[i]] == λ[[i]] rev[[i]], {i, 3}]
```

```
{True, True, True}
```

(* eigenvalues & left eigenvectors *)

```
{λ, lev} = Eigensystem[Q^T]
```

```
{{1, -1/3, 1/6}, {{1/3, 2, 1}, {1, -2, 1}, {-1/2, -1/2, 1}}}
```

(* Checking *)

```
Table[lev[[i]].Q == λ[[i]] lev[[i]], {i, 3}]
```

```
{True, True, True}
```

(* normalization constants *)

```
nconst = MapThread[Dot, {lev, rev}]
```

```
{10/3, 6, 15/8}
```

(* Normalized left eigenvectors; right eigenvectors unchanged *)

```
lev = lev/nconst
```

```
{{1/10, 3/5, 3/10}, {1/6, -1/3, 1/6}, {-4/15, -4/15, 8/15}}
```

(* Orthonormality: $\langle \chi_i | \psi_j \rangle = \delta_{ij}$ *)

Table[lev[[i]].rev[[j]], {i, 3}, {j, 3}] // MatrixForm

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

(* $|\psi_i\rangle\langle\chi_i|$ *)

r1Prod = Table[Outer[Times, rev[[i], lev[[i]]], {i, 3}];

MatrixForm/@r1Prod

$$\left\{ \begin{pmatrix} \frac{1}{10} & \frac{3}{5} & \frac{3}{10} \\ \frac{1}{10} & \frac{3}{5} & \frac{3}{10} \\ \frac{1}{10} & \frac{3}{5} & \frac{3}{10} \end{pmatrix}, \begin{pmatrix} \frac{1}{2} & -1 & \frac{1}{2} \\ -\frac{1}{6} & \frac{1}{3} & -\frac{1}{6} \\ \frac{1}{6} & -\frac{1}{3} & \frac{1}{6} \end{pmatrix}, \begin{pmatrix} \frac{2}{5} & \frac{2}{5} & -\frac{4}{5} \\ \frac{1}{15} & \frac{1}{15} & -\frac{2}{15} \\ -\frac{4}{15} & -\frac{4}{15} & \frac{8}{15} \end{pmatrix} \right\}$$

(* Completeness: $\sum_{i=1}^3 |\psi_i\rangle\langle\chi_i| = \mathbf{I}$ *)

$\sum_{i=1}^3$ r1Prod[[i]] // MatrixForm

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

(* $P_{1|1}(s_0|s)$ *)

P11[s0_, s_] := Sum[ev[[i]]^{s-s0} r1Prod[[i]], {i, 3}]

P11[0, s] // MatrixForm

$$\begin{pmatrix} \frac{1}{10} + \frac{1}{2} \left(-\frac{1}{3}\right)^s + \frac{1}{5} \times 2^{1-s} \times 3^{-s} & \frac{3}{5} + (-1)^{1+s} 3^{-s} + \frac{1}{5} \times 2^{1-s} \times 3^{-s} & \frac{3}{10} + \frac{1}{2} \left(-\frac{1}{3}\right)^s - \frac{1}{5} \times 2^{2-s} \times 3^{-s} \\ \frac{1}{10} - \frac{1}{2} \left(-1\right)^s 3^{-1-s} + \frac{1}{5} \times 2^{-s} \times 3^{-1-s} & \frac{3}{5} + (-1)^s 3^{-1-s} + \frac{1}{5} \times 2^{-s} \times 3^{-1-s} & \frac{3}{10} - \frac{1}{2} \left(-1\right)^s 3^{-1-s} - \frac{1}{5} \times 2^{1-s} \times 3^{-s} \\ \frac{1}{10} + \frac{1}{2} \left(-1\right)^s 3^{-1-s} - \frac{1}{5} \times 2^{2-s} \times 3^{-1-s} & \frac{3}{5} + \left(-\frac{1}{3}\right)^{1+s} - \frac{1}{5} \times 2^{2-s} \times 3^{-1-s} & \frac{3}{10} + \frac{1}{2} \left(-1\right)^s 3^{-1-s} + \frac{1}{5} \times 2^{3-s} \times 3^{-s} \end{pmatrix}$$

(* Q *)

P11[s, s+1] // MatrixForm

$$\begin{pmatrix} 0 & 1 & 0 \\ \frac{1}{6} & \frac{1}{2} & \frac{1}{3} \\ 0 & \frac{2}{3} & \frac{1}{3} \end{pmatrix}$$

(* Q^2 *)

P11[s, s+2] // MatrixForm

$$\begin{pmatrix} \frac{1}{6} & \frac{1}{2} & \frac{1}{3} \\ \frac{1}{12} & \frac{23}{36} & \frac{5}{18} \\ \frac{1}{9} & \frac{5}{9} & \frac{1}{3} \end{pmatrix}$$

(* Answer to (b) *)

p0 = {1, 0, 0};

ps = p0 . P11[0, s]

$$\left\{ \frac{1}{10} + \frac{1}{2} \left(-\frac{1}{3}\right)^s + \frac{1}{5} \times 2^{1-s} \times 3^{-s}, \frac{3}{5} + (-1)^{1+s} 3^{-s} + \frac{1}{5} \times 2^{1-s} \times 3^{-s}, \frac{3}{10} + \frac{1}{2} \left(-\frac{1}{3}\right)^s - \frac{1}{5} \times 2^{2-s} \times 3^{-s} \right\}$$

(* Answer to (c) *)

(* 1st moment *)

$$yb = \sum_{n=1}^3 n ps[[n]] // Simplify$$

$$\frac{11}{5} - \frac{6^{1-s}}{5}$$

(* autocorrelation function *)

$$yyb = \sum_{m=1}^3 \sum_{n=1}^3 m n p0[[m]] P11[\theta, s][[m, n]] // Simplify$$

$$\frac{11}{5} - \frac{6^{1-s}}{5}$$