

## 10.E.4. Power Absorption

Consider a system with extensive variable  $\mathbf{A} = \mathbf{A}_0 + \boldsymbol{\alpha}$  coupled to an external force (or intensive variable)  $\mathbf{F}$ . The work done on the system by  $\mathbf{F}$  is [ see §2.D.2 or Exercise 10.5 ]

$$-dW = \mathbf{F} \cdot d\boldsymbol{\alpha} \quad (10.155)$$

where  $dW$  is the work done by the system. The power absorption of the system is therefore

$$\begin{aligned} P(t) &= -\frac{dW}{dt} = \mathbf{F}(t) \cdot \dot{\boldsymbol{\alpha}}_F(t) \quad [\text{Dot product between vectors.}] \\ &= \mathbf{F}^T(t) \cdot \dot{\boldsymbol{\alpha}}_F(t) \quad [\text{Matrix multiplication between row \& column matrices.}] \\ &= \mathbf{F}^T(t) \cdot \frac{d}{dt} \int_{-\infty}^{\infty} dt' \mathbf{K}(t-t') \cdot \mathbf{F}(t') \quad [(10.113) \text{ used.}] \quad (10.156) \\ &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} \tilde{\mathbf{F}}^T(\omega) \cdot \frac{d}{dt} \int_{-\infty}^{\infty} dt' \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} e^{-i\omega'(t-t')} \boldsymbol{\chi}(\omega') \cdot \int_{-\infty}^{\infty} \frac{d\omega''}{2\pi} e^{-i\omega'' t'} \tilde{\mathbf{F}}(\omega'') \\ &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} dt' \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega''}{2\pi} (-i\omega') e^{-i(\omega+\omega')t} e^{-i(\omega''-\omega')t'} \tilde{\mathbf{F}}^T(\omega) \cdot \boldsymbol{\chi}(\omega') \cdot \tilde{\mathbf{F}}(\omega'') \\ &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \int_{-\infty}^{\infty} d\omega'' (-i\omega') \delta(\omega'' - \omega') e^{-i(\omega+\omega')t} \tilde{\mathbf{F}}^T(\omega) \cdot \boldsymbol{\chi}(\omega') \cdot \tilde{\mathbf{F}}(\omega'') \\ &= -i \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \omega' e^{-i(\omega+\omega')t} \tilde{\mathbf{F}}^T(\omega) \cdot \boldsymbol{\chi}(\omega') \cdot \tilde{\mathbf{F}}(\omega') \quad (10.157) \end{aligned}$$

### 10.E.4.1. Delta Function Force

Let

$$\mathbf{F}(t) = \mathbf{F} \delta(t) \quad \rightarrow \quad \tilde{\mathbf{F}}(\omega) = \int_{-\infty}^{\infty} dt e^{i\omega t} \mathbf{F} \delta(t) = \mathbf{F} \quad (10.158)$$

(10.157) becomes

$$P(t) = -i \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \omega' e^{-i(\omega+\omega')t} \mathbf{F}^T \cdot \boldsymbol{\chi}(\omega') \cdot \mathbf{F} \quad (10.159)$$

$$= -i \delta(t) \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \omega' e^{-i\omega' t} \mathbf{F}^T \cdot \boldsymbol{\chi}(\omega') \cdot \mathbf{F}$$

$$= -i \delta(t) \mathbf{F}^T \cdot \left[ \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega \boldsymbol{\chi}(\omega) \right] \cdot \mathbf{F}$$

$$= \delta(t) \mathbf{F}^T \cdot \left[ \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega \boldsymbol{\chi}''(\omega) \right] \cdot \mathbf{F} \quad [\boldsymbol{\chi}'(-\omega) = \boldsymbol{\chi}'(\omega)] \quad (10.159a)$$

$$= \delta(t) \mathbf{F}^T \cdot \left[ \int_0^{\infty} \frac{d\omega}{\pi} \omega \boldsymbol{\chi}''(\omega) \right] \cdot \mathbf{F} \quad [\boldsymbol{\chi}''(-\omega) = -\boldsymbol{\chi}''(\omega)] \quad (10.159b)$$

The total energy absorbed is therefore

$$\begin{aligned} W_{\text{abs}} &= \int_{-\infty}^{\infty} dt P(t) = \mathbf{F}^T \cdot \left[ \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega \boldsymbol{\chi}''(\omega) \right] \cdot \mathbf{F} \quad (10.160) \\ &= \mathbf{F}^T \cdot \left[ \int_0^{\infty} \frac{d\omega}{\pi} \omega \boldsymbol{\chi}''(\omega) \right] \cdot \mathbf{F} \end{aligned}$$

For a dissipative system [see Exercise 10.6],

$$\omega \boldsymbol{\chi}''(\omega) \geq 0 \quad \forall \omega$$

$$\boldsymbol{\chi}''(\omega) \geq 0 \quad \forall \omega \geq 0$$

therefore,

$$W_{\text{abs}} > 0$$

so that the system absorbs energy, part of which will turn into heat.

### 10.E.4.2. Oscillating Force

Let

$$\mathbf{F}(t) = \mathbf{F} \cos \omega_0 t = \frac{1}{2} \mathbf{F} (e^{i\omega_0 t} + e^{-i\omega_0 t}) \quad (10.161)$$

$$\begin{aligned} \rightarrow \tilde{\mathbf{F}}(\omega) &= \frac{1}{2} \mathbf{F} \int_{-\infty}^{\infty} dt e^{i\omega t} (e^{i\omega_0 t} + e^{-i\omega_0 t}) \\ &= \pi \mathbf{F} \left[ \delta(\omega + \omega_0) + \delta(\omega - \omega_0) \right] \end{aligned} \quad (10.162)$$

(10.157) becomes

$$\begin{aligned} P(t) &= -\frac{i}{4} \int_{-\infty}^{\infty} d\omega \int_{-\infty}^{\infty} d\omega' \omega' e^{-i(\omega + \omega')t} \mathbf{F}^T \cdot \chi(\omega') \cdot \mathbf{F} \\ &\quad \times \left[ \delta(\omega + \omega_0) + \delta(\omega - \omega_0) \right] \left[ \delta(\omega' + \omega_0) + \delta(\omega' - \omega_0) \right] \\ &= -\frac{i}{4} \int_{-\infty}^{\infty} d\omega' \omega' \left( e^{-i(-\omega_0 + \omega')t} + e^{-i(\omega_0 + \omega')t} \right) \mathbf{F}^T \cdot \chi(\omega') \cdot \mathbf{F} \left[ \delta(\omega' + \omega_0) + \delta(\omega' - \omega_0) \right] \\ &= -\frac{i}{4} \mathbf{F}^T \cdot \left[ -\omega_0 (e^{2i\omega_0 t} + 1) \chi(-\omega_0) + \omega_0 (1 + e^{-2i\omega_0 t}) \chi(\omega_0) \right] \cdot \mathbf{F} \quad (10.163) \\ &= \frac{1}{2} \mathbf{F}^T \cdot \text{Im} \left[ \omega_0 (1 + e^{-2i\omega_0 t}) \chi(\omega_0) \right] \cdot \mathbf{F} \quad (10.163a) \end{aligned}$$

Averaging over one period of oscillation, we have

$$\begin{aligned} \langle P(t) \rangle &= \frac{\omega_0}{2\pi} \int_0^{2\pi/\omega_0} dt \frac{1}{2} \mathbf{F}^T \cdot \text{Im} \left[ \omega_0 (1 + e^{-2i\omega_0 t}) \chi(\omega_0) \right] \cdot \mathbf{F} \\ &= \frac{1}{2} \mathbf{F}^T \cdot \text{Im} \left[ \omega_0 \chi(\omega_0) \right] \cdot \mathbf{F} \\ &= \frac{1}{2} \omega_0 \mathbf{F}^T \cdot \chi''(\omega_0) \cdot \mathbf{F} \end{aligned} \quad (10.164)$$

Similar to the case of the delta function force,  $\langle P(t) \rangle > 0$  for a dissipative system so that it absorbs power.

### Exercise 10.6

An 1-D Brownian particle of mass  $m$  is attached to a harmonic spring of force constant  $k$ , and is driven by an external force  $F(t)$ . The Langevin equation is

$$m \frac{d^2 x(t)}{dt^2} + \gamma \frac{dx(t)}{dt} + m \omega_0^2 x(t) = \xi(t) + F(t) \quad (1a)$$

where  $\omega_0 = \sqrt{\frac{k}{m}}$ ,  $\gamma$  is the friction constant, and  $\xi(t)$  is a Gaussian white noise with zero mean.

Hence,

$$m \frac{d^2 \langle x(t) \rangle_F}{dt^2} + \gamma \frac{d \langle x(t) \rangle_F}{dt} + m \omega_0^2 \langle x(t) \rangle_F = F(t) \quad (2)$$

(a) Compute and plot the linear response function  $K(t)$ , where

$$\langle x(t) \rangle_F = \int_{-\infty}^{\infty} dt' K(t-t') F(t') \quad (1b)$$

(b) Compute the total energy absorbed for the case  $F(t) = F_0 \delta(t)$ .

### Answer (a)

By causality,

$$K(t) = 0 \quad \forall t < 0 \quad (1c)$$

Since  $K$  is independent of  $F(t)$ , we can set  $F(t) = F_0 \delta(t)$  so that (1c) becomes

$$\begin{aligned} \langle x(t) \rangle_F &= \int_{-\infty}^t dt' K(t-t') F_0 \delta(t') & \forall t > 0 \\ &= K(t) F_0 \end{aligned} \quad (1)$$

Putting (1) into (2) gives

$$m \frac{d^2 K(t)}{dt^2} + \gamma \frac{dK(t)}{dt} + m \omega_0^2 K(t) = \delta(t) \quad (3)$$

which is often used as the defining equation of  $K(t)$ .

Taking the Fourier transform, (3) becomes

$$-m \omega^2 \chi(\omega) - i \omega \gamma \chi(\omega) + m \omega_0^2 \chi(\omega) = 1 \quad (4)$$

$$\rightarrow \chi(\omega) = \frac{1}{m(\omega_0^2 - \omega^2) - i \gamma \omega} \quad (5)$$

with

$$\chi'(\omega) = \frac{m(\omega_0^2 - \omega^2)}{m^2(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \quad (7)$$

$$\chi''(\omega) = \frac{\gamma \omega}{m^2(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \quad (8)$$

The inverse transform of (5) gives

$$K(t) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{e^{-i\omega t}}{m(\omega_0^2 - \omega^2) - i\gamma\omega} \quad (9)$$

where the poles of the integrand are at

$$\omega_{\pm} = \frac{1}{2m} \left( -i\gamma \pm \sqrt{-\gamma^2 + 4m^2\omega_0^2} \right) \quad (9a)$$

Using

$$\text{Res} \left[ \frac{e^{-i\omega t}}{m(\omega_0^2 - \omega^2) - i\gamma\omega}; \omega_{\pm} \right] = \frac{e^{-i\omega t}}{-2m\omega - i\gamma} \Big|_{\omega=\omega_{\pm}} = \mp \frac{e^{-i\omega_{\pm} t}}{\sqrt{-\gamma^2 + 4m^2\omega_0^2}}$$

Hence, for  $t > 0$ ,

$$\begin{aligned} K(t) &= -i \left\{ \text{Res} \left[ \frac{e^{-i\omega t}}{m(\omega_0^2 - \omega^2) - i\gamma\omega}; \omega_{+} \right] + \text{Res} \left[ \frac{e^{-i\omega t}}{m(\omega_0^2 - \omega^2) - i\gamma\omega}; \omega_{-} \right] \right\} \\ &= -i \frac{-e^{-i\omega_{+} t} + e^{-i\omega_{-} t}}{\sqrt{-\gamma^2 + 4m^2\omega_0^2}} \\ &= \frac{2 \sin \sqrt{-\frac{\gamma^2}{4m^2} + \omega_0^2} t}{\sqrt{-\gamma^2 + 4m^2\omega_0^2}} e^{-\gamma t/2m} \end{aligned} \quad (10)$$

Combining with (1c), we have

$$K(t) = \theta(t) \frac{\sin \Omega t}{m \Omega} e^{-\Gamma t} \quad (10a)$$

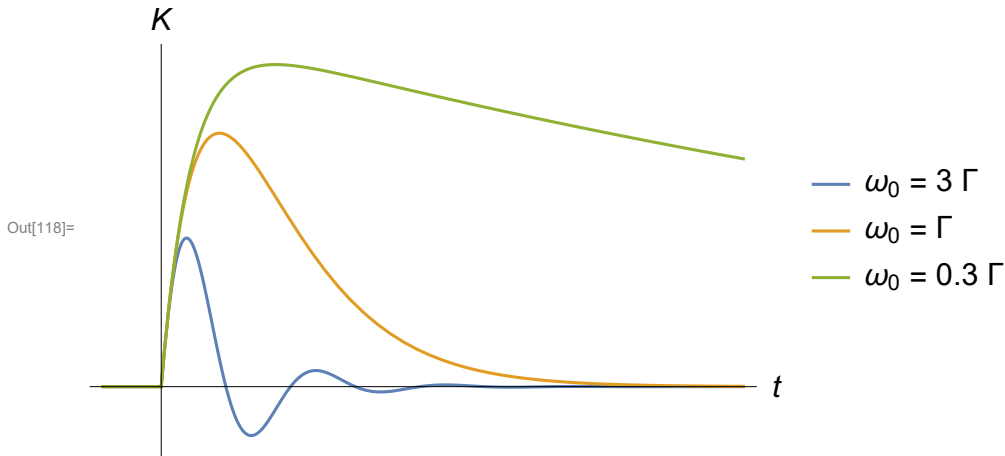
where

$$\Omega = \sqrt{-\frac{\gamma^2}{4m^2} + \omega_0^2} = \sqrt{-\Gamma^2 + \omega_0^2} \qquad \Gamma = \frac{\gamma}{2m} \qquad (10b)$$

Thus,

$$\begin{aligned} \omega_0 > \Gamma &\rightarrow \Omega \text{ is real} && (\text{system is underdamped.}) \\ \omega_0 = \Gamma &\rightarrow \Omega = 0, \frac{\sin \Omega t}{\Omega} = t && (\text{system is critically damped.}) \\ \omega_0 < \Gamma &\rightarrow \Omega \text{ is imaginary} && (\text{system is overdamped.}) \end{aligned}$$

The plots [see §Code] for these cases, with  $m = 1$ , are shown below.



### Answer (b)

Putting (8) into (10.160) gives

$$\begin{aligned} W_{\text{abs}} &= \int_0^\infty \frac{d\omega}{\pi} \frac{\gamma \omega^2}{m^2 (\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} F_0^2 \\ &= \frac{\gamma}{m \sqrt{2} \left( \sqrt{A+B-\gamma\sqrt{A}} + \sqrt{A+B+\gamma\sqrt{A}} \right)} F_0^2 \end{aligned}$$

[ See §Code. ]

where

$$A = \gamma^2 - 4m^2 \omega_0^2 \qquad B = 2m^2 \omega_0^2$$

Now,

$$\begin{aligned} (\sqrt{a} + \sqrt{b})^2 &= a + b + 2\sqrt{ab} \\ \rightarrow \sqrt{a} + \sqrt{b} &= \sqrt{a + b + 2\sqrt{ab}} \\ \therefore \sqrt{A+B-\gamma\sqrt{A}} + \sqrt{A+B+\gamma\sqrt{A}} &= \sqrt{2(A+B) + 2\sqrt{(A+B)^2 - \gamma^2 A}} \\ &= \sqrt{2} \sqrt{\left( \gamma^2 - 2m^2 \omega_0^2 + \sqrt{(\gamma^2 - 2m^2 \omega_0^2)^2 - \gamma^2 (\gamma^2 - 4m^2 \omega_0^2)} \right)} \\ &= \sqrt{2} \sqrt{\gamma^2 - 2m^2 \omega_0^2 + 2m^2 \omega_0^2} \end{aligned}$$

$$= \sqrt{2} \gamma$$

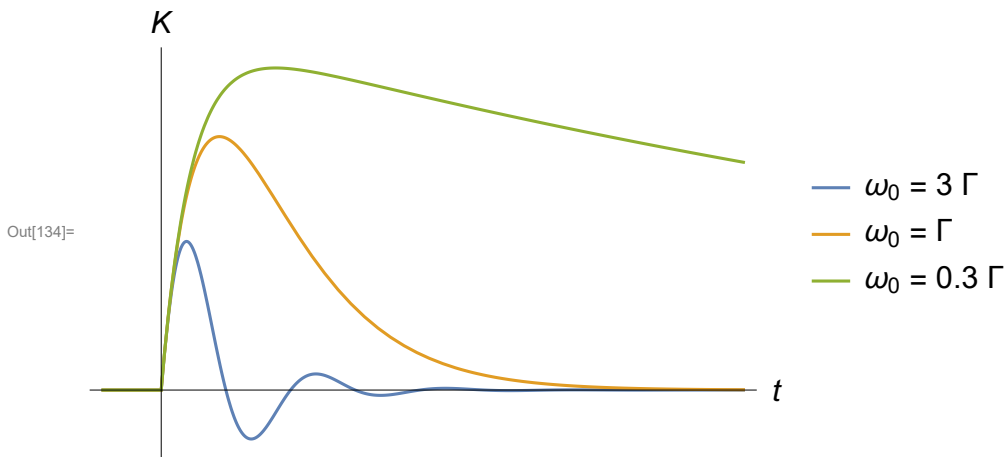
Hence,

$$W_{\text{abs}} = \frac{F_0^2}{2m} \quad (8)$$

## Code

```
In[131]= K[t_, Γ_, ω0_] := HeavisideTheta[t]  $\frac{e^{-\Gamma t} \text{Sin}[\sqrt{\omega_0^2 - \Gamma^2} t]}{\sqrt{\omega_0^2 - \Gamma^2}}$ 
```

```
In[132]= Γ = 1;
pL = {K[t, Γ, 3 Γ], K[t, Γ, (1 + 10^-8) Γ], K[t, Γ, 0.3 Γ]};
Plot[pL // Evaluate, {t, -1, 10}, PlotRange -> All,
  AxesLabel -> {"t", "K"}, Ticks -> None,
  PlotLegends -> {"ω0 = 3 Γ", "ω0 = Γ", "ω0 = 0.3 Γ"}]
```



```
In[136]= Assuming[m > 0 && γ > 0 && ω0 > 0,  $\int_0^\infty \frac{1}{\pi m^2 (\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} d\omega$ ] // Simplify
```

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Out[136]= ConditionalExpression[ $\frac{\gamma}{\left( \sqrt{2} m \left( \sqrt{\gamma^2 - 2 m^2 \omega_0^2 - \gamma \sqrt{\gamma^2 - 4 m^2 \omega_0^2}} + \sqrt{-2 m^2 \omega_0^2 + \gamma \left( \gamma + \sqrt{\gamma^2 - 4 m^2 \omega_0^2} \right)} \right) \right)}$ ,  $2 m \omega_0 \leq \gamma$ ]
```