

S4.E. General Form of Infinitely Divisible Distributions

S4.E.1. Levy-Khintchine Formula

Ref: See pp.310-1 of M.Loève, "Probability Theory I", 4th ed., Springer Verlag (1977).

The most general form of an infinitely divisible (or decomposable) characteristic function is

$$f_Y(k) = \exp \left[i k \alpha + \int \left(e^{i k x} - 1 - \frac{i k x}{1+x^2} \right) \frac{1+x^2}{x^2} dG(x) \right] \quad (4.166)$$

where α is a real constant, and G is a real, bounded, and nondecreasing function with $G(-\infty) = 0$. Indeed,

$$P_X(x) = \frac{dF_X(x)}{dx} = \frac{1+x^2}{Nx^2} \frac{dG(x)}{dx} \quad [N \rightarrow \infty] \quad (4.166a)$$

is the probability density function for X [c.f. (4.144)].

Thus, (4.166) is a generalization of the Kolmogorov formula by removing the requirements of finite variance and zero mean [see (4.146) & §S4.E.2]. It is called the **Levy-Khintchine formula**.

The proof that (4.166) is infinitely divisible goes as follows. Writing the integral

$$\begin{aligned} \mathcal{I}(k) &\equiv \int \left(e^{i k x} - 1 - \frac{i k x}{1+x^2} \right) \frac{1+x^2}{x^2} dG(x) \\ &= \int_{-\infty}^{\infty} dx \left(e^{i k x} - 1 - \frac{i k x}{1+x^2} \right) \frac{1+x^2}{x^2} \frac{dG(x)}{dx} \end{aligned}$$

as a Riemann sum, we have

$$\begin{aligned} \mathcal{I}(k) &= \lim_{M \rightarrow \infty} \sum_{m=-M}^M \Delta x \left(e^{i k x_m} - 1 - i k x_m \right) \frac{1+x_m^2}{x_m^2} \frac{G(x_m) - G(x_{m-1})}{\Delta x} \\ &= \lim_{M \rightarrow \infty} \sum_{m=-M}^M \left(e^{i k x_m} - 1 - i k x_m \right) \frac{1+x_m^2}{x_m^2} [G(x_m) - G(x_{m-1})] \\ &= \lim_{M \rightarrow \infty} \sum_{m=-M}^M \left(e^{i k x_m} - 1 - i k x_m \right) \lambda_m \quad \lambda_m = \frac{1+x_m^2}{x_m^2} [G(x_m) - G(x_{m-1})] \\ &= \lim_{M \rightarrow \infty} \sum_{m=-M}^M \left[i k a_m + \lambda_m (e^{i k x_m} - 1) \right] \quad a_m = \lambda_m x_m \end{aligned} \quad (4.169)$$

(4.166) then becomes

$$\begin{aligned} f_Y(k) &= \exp \left\{ i k \alpha + \lim_{M \rightarrow \infty} \sum_{m=-M}^M \left[i k a_m + \lambda_m (e^{i k x_m} - 1) \right] \right\} \\ &= \lim_{M \rightarrow \infty} \prod_{m=-M}^M \exp \left[i k \tilde{a}_m + \lambda_m (e^{i k x_m} - 1) \right] \quad \tilde{a}_m = \frac{\alpha}{2M+1} + a_m \end{aligned} \quad (4.169b)$$

which is a product of Poisson characteristic functions [see (4.126)]. Therefore, f_Y is infinitely divisible.

Again, we emphasize the fact that, unlike the Kolmogorov Formula, the Levy-Khintchine formula does not require G to have finite variance.

Using (4.17) of §4.D.3, we have

$$\langle y \rangle = -i \left. \frac{d f_Y(k)}{d k} \right|_{k=0} = -i \left\{ \left[i \alpha + \int_{-\infty}^{\infty} \left(i x e^{i k x} - i \frac{x}{1+x^2} \right) \frac{1+x^2}{x^2} dG(x) \right] f_Y(k) \right\} \Big|_{k=0}$$

$$\begin{aligned}
 &= \left[\alpha + \int \left(x - \frac{x}{1+x^2} \right) \frac{1+x^2}{x^2} dG(x) \right] f_Y(0) \\
 &= \alpha + \int x dG(x)
 \end{aligned} \tag{4.169c}$$

where setting $k = 0$ to (4.166) gives

$$f_Y(0) = 1$$

Hence,

$$\begin{aligned}
 \alpha &= \langle Y \rangle - \int x dG(x) \\
 &= \langle Y \rangle - N \int \frac{x^3}{1+x^2} dF_X(x) \quad [(4.166a) \text{ used. }] \quad [N \rightarrow \infty]
 \end{aligned} \tag{4.171}$$

For the Gaussian distribution, (4.123) gives

$$\begin{aligned}
 f_X(k) &= \exp \left(i \frac{a}{N} k - \frac{1}{2N} \sigma_Y^2 k^2 \right) \\
 &\approx 1 + i \frac{a}{N} k - \frac{1}{2N} \sigma_Y^2 k^2 \quad \text{for } N \rightarrow \infty \\
 &= 1 + \ln f_X(k) \quad [\ln f_X(k) \text{ calculated from the 1st equality. }] \\
 &= 1 + \lim_{N \rightarrow \infty} \ln f_Y(k)^{1/N} \quad [(4.121) \text{ used with } Y = \lim_{N \rightarrow \infty} \sum_{i=1}^N X_i]
 \end{aligned}$$

$$\rightarrow \ln f_Y(k) = \lim_{N \rightarrow \infty} N [f_X(k) - 1] \tag{4.170}$$

Putting (4.124) into (4.166a) gives

$$\begin{aligned}
 dG(x) &= \lim_{N \rightarrow \infty} N \frac{x^2}{1+x^2} P_X(x) dx \\
 &= \lim_{N \rightarrow \infty} N \frac{x^2}{1+x^2} \sqrt{\frac{N}{2\pi\sigma_Y^2}} \exp \left[-\frac{\left(x - \frac{a}{N}\right)^2}{2\frac{\sigma_Y^2}{N}} \right] dx \\
 \rightarrow G(x) &= \lim_{N \rightarrow \infty} N \sqrt{\frac{N}{2\pi\sigma_Y^2}} \int_{-\infty}^x \frac{u^2}{1+u^2} \exp \left[-\frac{\left(u - \frac{a}{N}\right)^2}{2\frac{\sigma_Y^2}{N}} \right] du
 \end{aligned} \tag{4.171a}$$

As its width narrows to zero, the Gaussian becomes a delta function:

$$\lim_{N \rightarrow \infty} \sqrt{\frac{N}{2\pi\sigma_Y^2}} \exp \left[-\frac{\left(u - \frac{a}{N}\right)^2}{2\frac{\sigma_Y^2}{N}} \right] = \delta(u) \tag{4.171b}$$

Unfortunately, (4.171b) cannot be used directly in (4.171a) owing to the extra factor N . Still, we can treat the Gaussian in (4.171a) as sharply peaked at $u = 0$ with a vanishing variance $\frac{\sigma_Y^2}{N}$. The integral therefore vanishes for $x < 0$, while, for $x > 0$, we can set $x = \infty$. Hence,

$$\begin{aligned}
G(x) &= \lim_{N \rightarrow \infty} N \sqrt{\frac{N}{2\pi\sigma_Y^2}} \int_{-\infty}^x u^2 (1 - u^2 + \dots) \exp\left[-\frac{\left(u - \frac{a}{N}\right)^2}{2\frac{\sigma_Y^2}{N}}\right] du \\
&= \lim_{N \rightarrow \infty} \begin{cases} N[\langle u^2 \rangle - \langle u^4 \rangle + \dots] & \text{for } x > 0 \\ 0 & \text{for } x < 0 \end{cases} \\
&= \lim_{N \rightarrow \infty} \begin{cases} N\left[\frac{\sigma_Y^2}{N} - \frac{3\sigma_Y^4}{N^2} + \dots\right] & \text{for } x > 0 \\ 0 & \text{for } x < 0 \end{cases} \quad [\text{See Mathematica code below.}] \\
&= \begin{cases} \sigma_Y^2 & \text{for } x > 0 \\ 0 & \text{for } x < 0 \end{cases} \\
&= \sigma_Y^2 \Theta(x) \tag{4.171c}
\end{aligned}$$

`In[]:= Assuming[σ > 0 && n > 0, Table[√(n/(2πσ²)) ∫_{-∞}^{∞} x^m e^{-x²/(2σ²/n)} dx, {m, 4}]] // PowerExpand`

`Out[]:= {0, σ²/n, 0, 3σ⁴/n²}`

Similarly, (4.169c) gives

$$\begin{aligned}
\alpha &= \langle y \rangle - N \sqrt{\frac{N}{2\pi\sigma_Y^2}} \int_{-\infty}^{\infty} \frac{x}{1+x^2} \exp\left[-\frac{\left(x - \frac{a}{N}\right)^2}{2\frac{\sigma_Y^2}{N}}\right] dx \\
&= \langle y \rangle - N \sqrt{\frac{N}{2\pi\sigma_Y^2}} \int_{-\infty}^{\infty} x(1-x^2 + \dots) \exp\left[-\frac{\left(x - \frac{a}{N}\right)^2}{2\frac{\sigma_Y^2}{N}}\right] dx \\
&= \langle y \rangle = a
\end{aligned}$$

where the integrals all vanish because their integrands are odd in x .

For the Poisson distribution, (4.126) gives

$$\begin{aligned}
f_X(k) &= \exp\left[i\frac{a}{N}k + \frac{\lambda}{N}(e^{ikh} - 1)\right] \\
&\approx 1 + i\frac{a}{N}k + \frac{\lambda}{N}(e^{ikh} - 1) \quad [N \rightarrow \infty] \\
&= 1 + \ln f_X(k)
\end{aligned}$$

which again leads to (4.170).

Putting (4.127) into (4.166a) gives

$$\begin{aligned}
dG(x) &= \lim_{N \rightarrow \infty} N \frac{x^2}{1+x^2} P_X(x) dx \\
&= \lim_{N \rightarrow \infty} N \frac{x^2}{1+x^2} \sum_{m=0}^{\infty} \frac{1}{m!} \left(\frac{\lambda}{N}\right)^m e^{-\lambda/N} \delta\left(x - \frac{a}{N} - mh\right) dx \\
\rightarrow G(x) &= \lim_{N \rightarrow \infty} N \sum_{m=0}^{\infty} \frac{1}{m!} \left(\frac{\lambda}{N}\right)^m e^{-\lambda/N} \int_{-\infty}^x \frac{u^2}{1+u^2} \delta\left(u - \frac{a}{N} - mh\right) du
\end{aligned}$$

$$\begin{aligned}
 &= \lim_{N \rightarrow \infty} N \sum_{m=0}^{\infty} \frac{1}{m!} \left(\frac{\lambda}{N}\right)^m e^{-\lambda/N} \frac{\left(\frac{a}{N} + mh\right)^2}{1 + \left(\frac{a}{N} + mh\right)^2} \Theta\left(x - \frac{a}{N} - mh\right) \\
 &= \lambda \frac{h^2}{1 + h^2} \Theta(x - h)
 \end{aligned} \tag{4.171d}$$

Similarly, (4.171) gives

$$\begin{aligned}
 \alpha = \langle Y \rangle &- \lim_{N \rightarrow \infty} N \sum_{m=0}^{\infty} \frac{1}{m!} \left(\frac{\lambda}{N}\right)^m e^{-\lambda/N} \int_{-\infty}^{\infty} \frac{x}{1+x^2} \delta\left(x - \frac{a}{N} - mh\right) dx \\
 &= \langle Y \rangle - \lim_{N \rightarrow \infty} N \sum_{m=0}^{\infty} \frac{1}{m!} \left(\frac{\lambda}{N}\right)^m e^{-\lambda/N} \frac{\frac{a}{N} + mh}{1 + \left(\frac{a}{N} + mh\right)^2} \\
 &= \langle Y \rangle - \lambda \frac{h}{1 + h^2}
 \end{aligned}$$

For the Cauchy distribution, (4.130) gives

$$\begin{aligned}
 f_X(k) &= \exp\left(i \frac{b}{N} k + \left| k \right| \frac{a}{N}\right) \\
 &\approx 1 + i \frac{a}{N} k + \left| k \right| \frac{a}{N} \quad [N \rightarrow \infty] \\
 &= 1 + \ln f_X(k)
 \end{aligned}$$

which also leads to (4.170).

Putting (4.131) into (4.166a) gives

$$\begin{aligned}
 dG(x) &= \lim_{N \rightarrow \infty} N \frac{x^2}{1+x^2} P_X(x) dx \\
 &= \lim_{N \rightarrow \infty} N \frac{x^2}{1+x^2} \frac{1}{\pi} \frac{\frac{a}{N}}{\left(\frac{a}{N}\right)^2 + \left(x - \frac{b}{N}\right)^2} dx \\
 \rightarrow G(x) &= \lim_{N \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^x \frac{u^2}{1+u^2} \frac{a}{\left(\frac{a}{N}\right)^2 + \left(u - \frac{b}{N}\right)^2} du \\
 &= \frac{a}{\pi} \int_{-\infty}^x \frac{1}{1+u^2} du \\
 &= \frac{a}{\pi} \left(\tan^{-1} x + \frac{\pi}{2} \right)
 \end{aligned}$$

$$\text{In[*]} := \left\{ \int \frac{1}{1+u^2} du, \int \frac{1}{(1+u^2)u} du \right\}$$

$$\text{Out[*]} := \left\{ \text{ArcTan}[u], \text{Log}[u] - \frac{1}{2} \text{Log}[1+u^2] \right\}$$

Similarly, (4.171b) gives

$$\alpha = \langle Y \rangle - \lim_{N \rightarrow \infty} N \int_{-\infty}^{\infty} \frac{x}{1+x^2} \frac{1}{\pi} \frac{\frac{a}{N}}{\left(\frac{a}{N}\right)^2 + \left(x - \frac{b}{N}\right)^2} dx$$

$$\begin{aligned}
&= \langle y \rangle - \lim_{N \rightarrow \infty} \frac{a}{\pi} \int_{-\infty}^{\infty} \frac{1}{(1+x^2)x} dx \\
&= \langle y \rangle = b \qquad \qquad \qquad \text{[Integrand odd in } x. \text{]}
\end{aligned}$$

Ex.4.15.

Show that the characteristic function

$$f(k) = \frac{1-b}{1+a} \frac{1+ae^{ik}}{1-be^{-ik}} \qquad 0 < a \leq b < 1 \qquad (1)$$

is not infinitely divisible.

Answer

Taking the logarithm of (1) gives

$$\begin{aligned}
\ln f(k) &= \ln(1-b) - \ln(1+a) + \ln(1+ae^{ik}) - \ln(1-be^{-ik}) \\
&= \sum_{n=1}^{\infty} \left[-\frac{b^n}{n} - (-)^{n+1} \frac{a^n}{n} + (-)^{n+1} \frac{a^n e^{ink}}{n} + \frac{b^n e^{-ink}}{n} \right] \\
&= \sum_{n=1}^{\infty} \left[\frac{b^n}{n} (e^{-ink} - 1) + (-)^{n+1} \frac{a^n}{n} (e^{ink} - 1) \right] \qquad (2)
\end{aligned}$$

Meanwhile, (4.166) gives

$$\begin{aligned}
\ln f(k) &= ik\alpha + \int \left(e^{ikx} - 1 - \frac{ikx}{1+x^2} \right) \frac{1+x^2}{x^2} dG(x) \\
&= ik\alpha + \int_{-\infty}^{\infty} (e^{ikx} - 1) \frac{1+x^2}{x^2} \frac{dG(x)}{dx} dx - ik \int_{-\infty}^{\infty} \frac{1}{x} \frac{dG(x)}{dx} dx \qquad (3)
\end{aligned}$$

Comparing (3) with (2), we get

$$\frac{dG(x)}{dx} = \sum_{n=1}^{\infty} \left[\delta(x+n) \frac{n^2}{1+n^2} \frac{b^n}{n} + \delta(x-n) (-)^{n+1} \frac{n^2}{1+n^2} \frac{a^n}{n} \right] \qquad (4)$$

and

$$\alpha = \int_{-\infty}^{\infty} \frac{1}{x} \frac{dG(x)}{dx} dx \qquad (5)$$

Putting (4) into (5) gives

$$\begin{aligned}
\alpha &= \sum_{n=1}^{\infty} \left[-\frac{1}{1+n^2} b^n + (-)^{n+1} \frac{1}{1+n^2} a^n \right] \\
&= -\sum_{n=1}^{\infty} \frac{b^n + (-)^n a^n}{1+n^2} \qquad (6)
\end{aligned}$$

With $G(-\infty) = 0$, (4) can be integrated to give

$$G(x) = \sum_{n=1}^{\infty} \left[\Theta(x+n) \frac{nb^n}{1+n^2} + \Theta(x-n) (-)^{n+1} \frac{na^n}{1+n^2} \right] \qquad (7)$$

For x equal to a positive even integer m , (4) gives

$$\left. \frac{dG(x)}{dx} \right|_{x=m} = (-)^{m+1} \frac{ma^m}{1+m^2} < 0$$

Hence, $G(x)$ is not a monotonically increasing function of x , i.e., it cannot be a distribution function and $\frac{dG(x)}{dx}$ is not a probability density. $f(k)$ is therefore not infinitely divisible.

S4.E.2. Kolmogorov Formula

If f_Y has finite second moment, i.e.,

$$N \int x^2 dF_X(x) = \int (1+x^2) dG \quad \text{is finite for } N \rightarrow \infty$$

we can define

$$\begin{aligned} dK(x) &= (1+x^2) dG(x) \\ \rightarrow K(x) &= \int_{-\infty}^x (1+u^2) \frac{dG(u)}{du} du \quad K(-\infty) = 0 \end{aligned} \quad (4.172a)$$

so that $K(x)$ is also a (unnormalized) distribution function.

(4.166) then becomes

$$\begin{aligned} f_Y(k) &= \exp \left[ik\alpha + \int \left(e^{ikx} - 1 - \frac{ikx}{1+x^2} \right) \frac{1}{x^2} dK(x) \right] \\ &= \exp \left\{ ik \left[\alpha + \int \frac{x}{1+x^2} dK(x) \right] + \int (e^{ikx} - 1 - ikx) \frac{1}{x^2} dK(x) \right\} \quad \frac{x}{1+x^2} = x - \frac{x^3}{1+x^2} \\ &= \exp \left\{ ik\gamma + \int (e^{ikx} - 1 - ikx) \frac{1}{x^2} dK(x) \right\} \end{aligned} \quad (4.172)$$

where

$$\gamma = \alpha + \int \frac{x}{1+x^2} dK(x) = \alpha + \int x dG(x) = \alpha + \int_{-\infty}^{\infty} x \frac{dG(x)}{dx} dx \quad (4.172b)$$

(4.172) is called the **Kolmogorov formula**.

Using (4.17) of §4.D.3, we have

$$\langle Y \rangle = -i \left. \frac{df_Y(k)}{dk} \right|_{k=0} = -i \left\{ \left[i\gamma + \int_{-\infty}^{\infty} (ix e^{ikx} - ix) \frac{1}{x^2} dK(x) \right] f_Y(k) \right\} \Big|_{k=0} = \gamma f_Y(0) = \gamma \quad (4.172c)$$

where (4.172) gives

$$f_Y(0) = 1$$

Similarly,

$$\begin{aligned} \langle Y^2 \rangle &= (-i)^2 \left. \frac{d^2 f_Y(k)}{dk^2} \right|_{k=0} \\ &= - \left\{ - \int_{-\infty}^{\infty} e^{ikx} dK(x) + \left[i\gamma + \int_{-\infty}^{\infty} (ix e^{ikx} - ix) \frac{1}{x^2} dK(x) \right]^2 \right\} f_Y(k) \Big|_{k=0} \\ &= K(\infty) + \gamma^2 \quad [(4.172a) \text{ used. }] \end{aligned}$$

Using (4.172b), we have

$$\text{var}(Y) = \sigma_Y^2 = \langle Y^2 \rangle - \langle Y \rangle^2 = K(\infty) \quad (4.172c)$$

For the Gaussian distribution, putting (4.171c) into (4.172a) gives

$$K(x) = \int_{-\infty}^x (1+u^2) \sigma_Y^2 \delta(u) du = \sigma_Y^2 \Theta(x)$$

For the Poisson distribution, putting (4.171d) into (4.172a) gives

$$K(x) = \int_{-\infty}^x (1+u^2) \lambda \frac{h^2}{1+h^2} \Theta(u-h) du = \lambda h^2 \Theta(x-h)$$

$$\rightarrow dK(x) = \lambda h^2 \delta(x-h) dx$$

Ex. 4.16.

Consider the characteristic function

$$f(k) = \frac{1-b}{1-b e^{ik}} \quad 0 < b < 1 \quad (1)$$

Find the Kolmogorov function $K(u)$.

Answer

$$\begin{aligned} \ln f(k) &= \ln(1-b) - \ln(1-b e^{ik}) \\ &= \sum_{n=1}^{\infty} \left(-\frac{b^n}{n} + \frac{b^n e^{ink}}{n} \right) \\ &= \sum_{n=1}^{\infty} \frac{b^n}{n} (e^{ink} - 1) \end{aligned} \quad (2)$$

Meanwhile, (4.172) gives

$$\ln f(k) = ik\gamma + \int_{-\infty}^{\infty} dx (e^{ikx} - 1 - ikx) \frac{1}{x^2} \frac{dK(x)}{dx} \quad (3)$$

Comparing (2) & (3) gives

$$\frac{dK(x)}{dx} = \sum_{n=1}^{\infty} \delta(x-n) b^n n \quad (4)$$

$$\& \quad \gamma = \int_{-\infty}^{\infty} dx \frac{1}{x} \frac{dK(x)}{dx} \quad (5)$$

With $K(-\infty) = 0$, (4) can be integrated to give

$$K(x) = \sum_{n=1}^{\infty} \theta(x-n) b^n n \quad (6)$$

Putting (4) into (5) gives

$$\gamma = \int_{-\infty}^{\infty} dx \frac{1}{x} \sum_{n=1}^{\infty} \delta(x-n) b^n n = \sum_{n=1}^{\infty} b^n = \frac{b}{1-b} \quad [b < 1 \text{ used.}]$$

(4) gives

$$\frac{dK(x)}{dx} > 0 \quad \forall x$$

while so that $\frac{1}{x^2} \frac{dK(x)}{dx}$ can be used as a probability function with moments

$$\langle x \rangle = \int_{-\infty}^{\infty} dx \sum_{n=1}^{\infty} \delta(x-n) b^n n \frac{1}{x} = \sum_{n=1}^{\infty} b^n = \frac{b}{1-b} = \gamma$$

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} dx \sum_{n=1}^{\infty} \delta(x-n) b^n n = \sum_{n=1}^{\infty} b^n n = b \frac{\partial}{\partial b} \sum_{n=1}^{\infty} b^n = b \frac{\partial}{\partial b} \frac{1}{1-b} = \frac{b}{(1-b)^2}$$

$$\rightarrow \sigma^2 = \frac{b}{(1-b)^2} - \left(\frac{b}{1-b}\right)^2 = \frac{b}{1-b} \text{ is finite.}$$

Thus, (4.172) is applicable to $f(k)$, which is therefore infinitely divisible.