

S10.A. Onsager's Relations When a Magnetic Field is Present

Examples of variables that change sign under time reversal are

$$\begin{aligned}
 \mathbf{v} = \frac{d\mathbf{x}}{dt} & \quad \rightarrow \quad \mathbf{v} \xrightarrow{t \rightarrow -t} -\mathbf{v} \\
 \mathbf{p} = m\mathbf{v} & \quad \rightarrow \quad \mathbf{p} \xrightarrow{t \rightarrow -t} -\mathbf{p} \\
 \nabla \times \mathbf{B} = \frac{4\pi}{c} \mathbf{J} = \frac{4\pi}{c} \rho \mathbf{v} & \rightarrow \quad \mathbf{B} \xrightarrow{t \rightarrow -t} -\mathbf{B} \\
 \mathbf{L} = \mathbf{r} \times \mathbf{p} = m \mathbf{r} \times \mathbf{v} & \rightarrow \quad \mathbf{L} \xrightarrow{t \rightarrow -t} -\mathbf{L} \\
 \boldsymbol{\omega} = \frac{d\theta}{dt} & \rightarrow \quad \boldsymbol{\omega} \xrightarrow{t \rightarrow -t} -\boldsymbol{\omega}
 \end{aligned} \tag{10.203a}$$

Proof that these conclusions hold for more general forms of the variables is straightforward and can be dealt with on a case by case basis.

Let $\boldsymbol{\alpha} = (\alpha_1, \dots, \alpha_n)$ & $\boldsymbol{\beta} = (\beta_1, \dots, \beta_m)$ be fluctuations such that

$$\boldsymbol{\alpha} \xrightarrow{t \rightarrow -t} \boldsymbol{\alpha} \quad \boldsymbol{\beta} \xrightarrow{t \rightarrow -t} -\boldsymbol{\beta} \tag{10.203b}$$

The change in entropy due to these fluctuations can be written as

$$\begin{aligned}
 \Delta S &= -\frac{1}{2} (\boldsymbol{\alpha}^T, \boldsymbol{\beta}^T) \mathbf{G} \begin{pmatrix} \boldsymbol{\alpha} \\ \boldsymbol{\beta} \end{pmatrix} \\
 &= -\frac{1}{2} (\boldsymbol{\alpha}^T, \boldsymbol{\beta}^T) \begin{pmatrix} \mathbf{g} & \mathbf{n} \\ \mathbf{m} & \mathbf{h} \end{pmatrix} \begin{pmatrix} \boldsymbol{\alpha} \\ \boldsymbol{\beta} \end{pmatrix} \quad \mathbf{G} = \begin{pmatrix} \mathbf{g} & \mathbf{n} \\ \mathbf{m} & \mathbf{h} \end{pmatrix} \\
 &= -\frac{1}{2} (\boldsymbol{\alpha}^T \mathbf{g} \boldsymbol{\alpha} + \boldsymbol{\alpha}^T \mathbf{n} \boldsymbol{\beta} + \boldsymbol{\beta}^T \mathbf{m} \boldsymbol{\alpha} + \boldsymbol{\beta}^T \mathbf{h} \boldsymbol{\beta}) \tag{10.203c} \\
 &= -\frac{1}{2} (\mathbf{g} : \boldsymbol{\alpha} \boldsymbol{\alpha} + \mathbf{n} : \boldsymbol{\alpha} \boldsymbol{\beta} + \mathbf{m} : \boldsymbol{\beta} \boldsymbol{\alpha} + \mathbf{h} : \boldsymbol{\beta} \boldsymbol{\beta}) \tag{10.203}
 \end{aligned}$$

which is the natural generalization of (7.22) in §7.C.1. The square matrices \mathbf{G} , \mathbf{g} & \mathbf{h} are assumed to be regular, i.e., their inverses exist.

Note that as matrices, \mathbf{m} is $m \times n$ and \mathbf{n} is $n \times m$. Since (10.203) is a scalar equation, taking its transpose does not alter its value. Hence,

$$\begin{aligned}
 \mathbf{G}^T &= \mathbf{G} \\
 \rightarrow \begin{pmatrix} \mathbf{g}^T & \mathbf{m}^T \\ \mathbf{n}^T & \mathbf{h}^T \end{pmatrix} &= \begin{pmatrix} \mathbf{g} & \mathbf{n} \\ \mathbf{m} & \mathbf{h} \end{pmatrix} \\
 \mathbf{g}^T = \mathbf{g} \quad \mathbf{h}^T = \mathbf{h} \quad \mathbf{m} = \mathbf{n}^T & \tag{10.202}
 \end{aligned}$$

In the presence of a magnetic field \mathbf{B} , the distribution of these fluctuations is given by (7.20) as

$$\begin{aligned}
 f(\boldsymbol{\alpha}, \boldsymbol{\beta}, \mathbf{B}) &= C e^{\Delta S / k_B} \\
 &= C \exp \left[-\frac{1}{2 k_B} (\boldsymbol{\alpha}^T, \boldsymbol{\beta}^T) \begin{pmatrix} \mathbf{g} & \mathbf{n} \\ \mathbf{m} & \mathbf{h} \end{pmatrix} \begin{pmatrix} \boldsymbol{\alpha} \\ \boldsymbol{\beta} \end{pmatrix} \right] \\
 &= C \exp \left[-\frac{1}{2 k_B} (\boldsymbol{\alpha}^T \mathbf{g} \boldsymbol{\alpha} + \boldsymbol{\alpha}^T \mathbf{n} \boldsymbol{\beta} + \boldsymbol{\beta}^T \mathbf{m} \boldsymbol{\alpha} + \boldsymbol{\beta}^T \mathbf{h} \boldsymbol{\beta}) \right] \tag{10.199}
 \end{aligned}$$

where C is a normalization constant and the tensors \mathbf{g} , \mathbf{m} , \mathbf{n} & \mathbf{h} may be functions of \mathbf{B} .

In both classical mechanics and non-relativistic quantum mechanics, it is assumed (and proven by experiments) that time reversal symmetry holds. For simplicity, we shall assume the Hamiltonian H itself is invariant under time reversal. Note that this includes the case when an EM field is present.

Thus, $f(\alpha, \beta, B)$ must be invariant under time reversal so that

$$f(\alpha, -\beta, -B) = f(\alpha, \beta, B) \quad (10.197)$$

(10.199) then gives

$$\begin{aligned} & (\alpha^T, -\beta^T) \mathbf{G}(-B) \begin{pmatrix} \alpha \\ -\beta \end{pmatrix} = (\alpha^T, \beta^T) \mathbf{G}(B) \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ \rightarrow & (\alpha^T, -\beta^T) \begin{pmatrix} \mathbf{g}(-B) & \mathbf{n}(-B) \\ \mathbf{m}(-B) & \mathbf{h}(-B) \end{pmatrix} \begin{pmatrix} \alpha \\ -\beta \end{pmatrix} = (\alpha^T, \beta^T) \begin{pmatrix} \mathbf{g}(B) & \mathbf{n}(B) \\ \mathbf{m}(B) & \mathbf{h}(B) \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ & = (\alpha^T, -\beta^T) \begin{pmatrix} \mathbf{g}(-B) & -\mathbf{n}(-B) \\ \mathbf{m}(-B) & -\mathbf{h}(-B) \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ & = (\alpha^T, \beta^T) \begin{pmatrix} \mathbf{g}(-B) & -\mathbf{n}(-B) \\ -\mathbf{m}(-B) & \mathbf{h}(-B) \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ \therefore & \begin{pmatrix} \mathbf{g}(-B) & -\mathbf{n}(-B) \\ -\mathbf{m}(-B) & \mathbf{h}(-B) \end{pmatrix} = \begin{pmatrix} \mathbf{g}(B) & \mathbf{n}(B) \\ \mathbf{m}(B) & \mathbf{h}(B) \end{pmatrix} \quad (10.200-1) \\ \rightarrow & \mathbf{G}(-B) = \begin{pmatrix} \mathbf{g}(-B) & \mathbf{n}(-B) \\ \mathbf{m}(-B) & \mathbf{h}(-B) \end{pmatrix} = \begin{pmatrix} \mathbf{g}(B) & -\mathbf{n}(B) \\ -\mathbf{m}(B) & \mathbf{h}(B) \end{pmatrix} \quad (10.201a) \end{aligned}$$

Note that for $B = 0$,

$$\begin{pmatrix} \mathbf{g}(0) & -\mathbf{n}(0) \\ -\mathbf{m}(0) & \mathbf{h}(0) \end{pmatrix} = \begin{pmatrix} \mathbf{g}(0) & \mathbf{n}(0) \\ \mathbf{m}(0) & \mathbf{h}(0) \end{pmatrix} \quad (10.200a)$$

$$\rightarrow \mathbf{m}(0) = \mathbf{n}^T(0) = \mathbf{0} \quad (10.202a)$$

Hence, there is no mixing between α & β when $B = 0$.

There are now two types of generalized forces [c.f. (10.79)],

$$\begin{aligned} \mathbf{F} &= -T \frac{\partial \Delta S}{\partial \alpha} = \frac{1}{2} T \left[\mathbf{g} \alpha + (\alpha^T \mathbf{g})^T + \mathbf{n} \beta + (\beta^T \mathbf{m})^T \right] \\ &= T (\mathbf{g} \alpha + \mathbf{m}^T \beta) \quad [(10.202) \text{ used. }] \\ &= T (\mathbf{g} \alpha + \mathbf{n} \beta) \quad (10.204) \\ &\xrightarrow[t \rightarrow -t]{} \mathbf{F} \end{aligned}$$

and

$$\begin{aligned} \mathbf{G} &= -T \frac{\partial \Delta S}{\partial \beta} = \frac{1}{2} T \left[(\alpha^T \mathbf{n})^T + \mathbf{m} \alpha + \mathbf{h} \beta + (\beta^T \mathbf{h})^T \right] \\ &= T (\mathbf{n}^T \alpha + \mathbf{h} \beta) = T (\mathbf{m} \alpha + \mathbf{h} \beta) \quad (10.205) \\ &\xrightarrow[t \rightarrow -t]{} -\mathbf{G} \end{aligned}$$

In order to evaluate the correlation functions, we retrace the steps leading from (7.24) to (7.27) in in §7.C.1. To begin,

$$\begin{aligned} \mathcal{I} &= \int d\alpha \int d\beta \exp \left[-\frac{1}{2k_B} (\alpha^T, \beta^T) \mathbf{G} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \right] \\ &= \int d\alpha \int d\beta \exp \left[-\frac{1}{2k_B} (\alpha^T \mathbf{g} \alpha + \alpha^T \mathbf{n} \beta + \beta^T \mathbf{m} \alpha + \beta^T \mathbf{h} \beta) \right] \end{aligned}$$

$$\begin{aligned}
&= \int d\boldsymbol{\alpha} \int d\boldsymbol{\beta} \exp\left[-\frac{1}{2k_B} \left\{ (\boldsymbol{\alpha} + \mathbf{g}^{-1} \mathbf{n} \boldsymbol{\beta})^T \mathbf{g} (\boldsymbol{\alpha} + \mathbf{g}^{-1} \mathbf{n} \boldsymbol{\beta}) + \boldsymbol{\beta}^T (\mathbf{h} - \mathbf{m} \mathbf{g}^{-1} \mathbf{n}) \boldsymbol{\beta} \right\}\right] \\
&= \int d\tilde{\boldsymbol{\alpha}} \int d\boldsymbol{\beta} \exp\left[-\frac{1}{2k_B} (\tilde{\boldsymbol{\alpha}}^T \mathbf{g} \tilde{\boldsymbol{\alpha}} + \boldsymbol{\beta}^T \tilde{\mathbf{h}} \boldsymbol{\beta})\right] \quad (10.205a)
\end{aligned}$$

where

$$\tilde{\boldsymbol{\alpha}} = \boldsymbol{\alpha} + \mathbf{g}^{-1} \mathbf{n} \boldsymbol{\beta} \quad \tilde{\mathbf{h}} = \mathbf{h} - \mathbf{m} \mathbf{g}^{-1} \mathbf{n} \quad (10.205b)$$

Carrying out the Gaussian integrals gives

$$I = \sqrt{\frac{(2\pi k_B)^n}{\det \mathbf{g}}} \sqrt{\frac{(2\pi k_B)^m}{\det \tilde{\mathbf{h}}}} = \sqrt{\frac{(2\pi k_B)^{n+m}}{\det \mathbf{G}}} \quad (10.205c)$$

where

$$\det \mathbf{G} = \det \mathbf{g} \cdot \det \tilde{\mathbf{h}}$$

is easily proved using the Gaussian elimination method to turn \mathbf{G} into an upper-block-triangular, or block-diagonal, matrix.

Comparing with (10.199) gives

$$C = \sqrt{\frac{\det \mathbf{G}}{(2\pi k_B)^{n+m}}} \quad (10.205d)$$

so that

$$\int d\boldsymbol{\alpha} \int d\boldsymbol{\beta} f(\boldsymbol{\alpha}, \boldsymbol{\beta}, \mathbf{B}) = 1 \quad (10.205e)$$

Alternatively, we can write (10.205a) as

$$\begin{aligned}
I &= \int d\boldsymbol{\alpha} \int d\boldsymbol{\beta} \exp\left[-\frac{1}{2k_B} \left\{ \boldsymbol{\alpha}^T (\mathbf{g} - \mathbf{n} \mathbf{h}^{-1} \mathbf{m}) \boldsymbol{\alpha} + (\boldsymbol{\beta} + \mathbf{h}^{-1} \mathbf{m} \boldsymbol{\alpha})^T \mathbf{h} (\boldsymbol{\beta} + \mathbf{h}^{-1} \mathbf{m} \boldsymbol{\alpha}) \right\}\right] \\
&= \int d\boldsymbol{\alpha} \int d\tilde{\boldsymbol{\beta}} \exp\left[-\frac{1}{2k_B} (\boldsymbol{\alpha}^T \tilde{\mathbf{g}} \boldsymbol{\alpha} + \tilde{\boldsymbol{\beta}}^T \mathbf{h} \tilde{\boldsymbol{\beta}})\right] \quad (10.205f)
\end{aligned}$$

where

$$\tilde{\boldsymbol{\beta}} = \boldsymbol{\beta} + \mathbf{h}^{-1} \mathbf{m} \boldsymbol{\alpha} \quad \tilde{\mathbf{g}} = \mathbf{g} - \mathbf{n} \mathbf{h}^{-1} \mathbf{m} \quad (10.205g)$$

The correlation functions are calculated using the generalized form of (7.26).

$$\begin{aligned}
I(\mathbf{h}, \mathbf{k}) &= C \int d\tilde{\boldsymbol{\alpha}} \int d\boldsymbol{\beta} \exp\left[-\frac{1}{2k_B} (\tilde{\boldsymbol{\alpha}}^T \mathbf{g} \tilde{\boldsymbol{\alpha}} + \boldsymbol{\beta}^T \tilde{\mathbf{h}} \boldsymbol{\beta}) + \mathbf{h}^T \tilde{\boldsymbol{\alpha}} + \mathbf{k}^T \boldsymbol{\beta}\right] \\
&= C \int d\bar{\boldsymbol{\alpha}} \int d\bar{\boldsymbol{\beta}} \exp\left[-\frac{1}{2k_B} (\bar{\boldsymbol{\alpha}}^T \mathbf{g} \bar{\boldsymbol{\alpha}} + \bar{\boldsymbol{\beta}}^T \tilde{\mathbf{h}} \bar{\boldsymbol{\beta}}) + \frac{1}{2} k_B (\mathbf{h}^T \mathbf{g}^{-1} \mathbf{h} + \mathbf{k}^T \tilde{\mathbf{h}}^{-1} \mathbf{k})\right] \quad (10.205h)
\end{aligned}$$

where

$$\bar{\boldsymbol{\alpha}} = \tilde{\boldsymbol{\alpha}} - k_B \mathbf{g}^{-1} \mathbf{h} \quad \bar{\boldsymbol{\beta}} = \boldsymbol{\beta} - k_B \tilde{\mathbf{h}}^{-1} \mathbf{k}$$

Doing the Gaussian integrals gives

$$I(\mathbf{h}, \mathbf{k}) = \exp\left[\frac{1}{2} k_B (\mathbf{h}^T \mathbf{g}^{-1} \mathbf{h} + \mathbf{k}^T \tilde{\mathbf{h}}^{-1} \mathbf{k})\right] \quad (10.205i)$$

Hence,

$$\langle \tilde{\alpha}_i \tilde{\alpha}_j \rangle = \int d\tilde{\boldsymbol{\alpha}} \int d\boldsymbol{\beta} f(\tilde{\boldsymbol{\alpha}}, \boldsymbol{\beta}) \tilde{\alpha}_i \tilde{\alpha}_j$$

$$\begin{aligned}
 &= \frac{\partial^2 \mathcal{I}(\mathbf{h}, \mathbf{k})}{\partial h_i \partial h_j} \Big|_{\mathbf{h}=\mathbf{k}=0} && \text{[(10.205f) used.]} \\
 &= \frac{\partial}{\partial h_i} \left[\frac{1}{2} k_B (\delta_{kj} g_{km}^{-1} h_m + h_k g_{km}^{-1} \delta_{mj}) \mathcal{I}(\mathbf{h}, \mathbf{k}) \right]_{\mathbf{h}=\mathbf{k}=0} \\
 &= \frac{\partial}{\partial h_i} \left[\frac{1}{2} k_B (g_{jm}^{-1} h_m + h_k g_{kj}^{-1}) \mathcal{I}(\mathbf{h}, \mathbf{k}) \right]_{\mathbf{h}=\mathbf{k}=0} \\
 &= \left[\frac{1}{2} k_B (g_{jm}^{-1} \delta_{mi} + \delta_{ki} g_{kj}^{-1}) \mathcal{I}(\mathbf{h}, \mathbf{k}) \right. \\
 &\quad \left. + \frac{1}{4} k_B^2 (g_{jm}^{-1} h_m + h_k g_{kj}^{-1}) (g_{im}^{-1} h_{m'} + h_{k'} g_{k'i}^{-1}) \mathcal{I}(\mathbf{h}, \mathbf{k}) \right]_{\mathbf{h}=\mathbf{k}=0} \\
 &= k_B g_{ij}^{-1}
 \end{aligned}$$

or

$$\langle \tilde{\alpha} \tilde{\alpha}^T \rangle = k_B \mathbf{g}^{-1} \quad (10.206a)$$

Similarly,

$$\langle \tilde{\alpha} \tilde{\beta}^T \rangle = 0 \quad (10.207a)$$

$$\langle \tilde{\beta} \tilde{\alpha}^T \rangle = 0 \quad (10.207b)$$

$$\langle \tilde{\beta} \tilde{\beta}^T \rangle = k_B \tilde{\mathbf{h}}^{-1} = k_B (\mathbf{h} - \mathbf{m} \mathbf{g}^{-1} \mathbf{n})^{-1} \quad \text{[(10.205b) used.]} \quad (10.209)$$

Using (10.205b) on (10.207a) gives

$$\langle (\boldsymbol{\alpha} + \mathbf{g}^{-1} \mathbf{n} \boldsymbol{\beta}) \boldsymbol{\beta}^T \rangle = 0$$

$$\rightarrow \langle \boldsymbol{\alpha} \boldsymbol{\beta}^T \rangle = -\mathbf{g}^{-1} \mathbf{n} \langle \tilde{\beta} \tilde{\beta}^T \rangle = -k_B \mathbf{g}^{-1} \mathbf{n} (\mathbf{h} - \mathbf{m} \mathbf{g}^{-1} \mathbf{n})^{-1} \quad (10.207c)$$

Taking the transpose gives

$$\langle \tilde{\beta} \boldsymbol{\alpha}^T \rangle = -k_B (\mathbf{h} - \mathbf{m} \mathbf{g}^{-1} \mathbf{n})^{-1} \mathbf{m} \mathbf{g}^{-1} \quad (10.208)$$

Alternatively, we can use (10.205a) to write the generator of the correlation functions as

$$\begin{aligned}
 \mathcal{J}(\mathbf{h}, \mathbf{k}) &= \int d\boldsymbol{\alpha} \int d\tilde{\boldsymbol{\beta}} \exp \left[-\frac{1}{2k_B} (\boldsymbol{\alpha}^T \tilde{\mathbf{g}} \boldsymbol{\alpha} + \tilde{\boldsymbol{\beta}}^T \tilde{\mathbf{h}} \tilde{\boldsymbol{\beta}}) + \mathbf{h}^T \boldsymbol{\alpha} + \mathbf{k}^T \tilde{\boldsymbol{\beta}} \right] \\
 &= C \int d\hat{\boldsymbol{\alpha}} \int d\hat{\boldsymbol{\beta}} \exp \left[-\frac{1}{2k_B} (\hat{\boldsymbol{\alpha}}^T \tilde{\mathbf{g}} \hat{\boldsymbol{\alpha}} + \hat{\boldsymbol{\beta}}^T \tilde{\mathbf{h}} \hat{\boldsymbol{\beta}}) + \frac{1}{2} k_B (\mathbf{h}^T \tilde{\mathbf{g}}^{-1} \mathbf{h} + \mathbf{k}^T \tilde{\mathbf{h}}^{-1} \mathbf{k}) \right] \\
 &= \exp \left[\frac{1}{2} k_B (\mathbf{h}^T \tilde{\mathbf{g}}^{-1} \mathbf{h} + \mathbf{k}^T \tilde{\mathbf{h}}^{-1} \mathbf{k}) \right]
 \end{aligned} \quad (10.208a)$$

where

$$\hat{\boldsymbol{\alpha}} = \boldsymbol{\alpha} - k_B \tilde{\mathbf{g}}^{-1} \mathbf{h} \quad \hat{\boldsymbol{\beta}} = \tilde{\boldsymbol{\beta}} - k_B \tilde{\mathbf{h}}^{-1} \mathbf{k}$$

Following the same route as used for $\mathcal{I}(\mathbf{h}, \mathbf{k})$, we get

$$\langle \boldsymbol{\alpha} \boldsymbol{\alpha}^T \rangle = k_B \tilde{\mathbf{g}}^{-1} = k_B (\mathbf{g} - \mathbf{n} \mathbf{h}^{-1} \mathbf{m})^{-1} \quad (10.206)$$

$$\langle \boldsymbol{\alpha} \tilde{\boldsymbol{\beta}}^T \rangle = \langle \boldsymbol{\alpha} (\boldsymbol{\beta} + \mathbf{h}^{-1} \mathbf{m} \boldsymbol{\alpha})^T \rangle = 0$$

$$\rightarrow \langle \boldsymbol{\alpha} \boldsymbol{\beta}^T \rangle = -\langle \boldsymbol{\alpha} \boldsymbol{\alpha}^T \rangle \mathbf{n} \mathbf{h}^{-1} = -k_B (\mathbf{g} - \mathbf{n} \mathbf{h}^{-1} \mathbf{m})^{-1} \mathbf{n} \mathbf{h}^{-1} \quad (10.207)$$

$$\langle \boldsymbol{\beta} \boldsymbol{\alpha}^T \rangle = \langle \boldsymbol{\alpha} \boldsymbol{\beta}^T \rangle^T = -k_B \mathbf{h}^{-1} \mathbf{m} (\mathbf{g} - \mathbf{n} \mathbf{h}^{-1} \mathbf{m})^{-1} \quad (10.207d)$$

$$\langle \tilde{\boldsymbol{\beta}} \tilde{\boldsymbol{\beta}}^T \rangle = k_B \tilde{\mathbf{h}}^{-1}$$

Using (10.201a) on

$$\mathcal{G}^{-1}(\mathbf{B}) \equiv \begin{pmatrix} \langle \boldsymbol{\alpha} \boldsymbol{\alpha}^T; \mathbf{B} \rangle & \langle \boldsymbol{\alpha} \boldsymbol{\beta}^T; \mathbf{B} \rangle \\ \langle \boldsymbol{\beta} \boldsymbol{\alpha}^T; \mathbf{B} \rangle & \langle \boldsymbol{\beta} \boldsymbol{\beta}^T; \mathbf{B} \rangle \end{pmatrix}$$

$$= k_B \begin{pmatrix} [\mathbf{g}(\mathbf{B}) - \mathbf{n}(\mathbf{B}) \mathbf{h}^{-1}(\mathbf{B}) \mathbf{m}(\mathbf{B})]^{-1} & -[\mathbf{g}(\mathbf{B}) - \mathbf{n}(\mathbf{B}) \mathbf{h}^{-1}(\mathbf{B}) \mathbf{m}(\mathbf{B})]^{-1} \mathbf{n}(\mathbf{B}) \mathbf{h}^{-1}(\mathbf{B}) \\ -[\mathbf{h}(\mathbf{B}) - \mathbf{m}(\mathbf{B}) \mathbf{g}^{-1}(\mathbf{B}) \mathbf{n}(\mathbf{B})]^{-1} \mathbf{m}(\mathbf{B}) \mathbf{g}(\mathbf{B})^{-1} & [\mathbf{h}(\mathbf{B}) - \mathbf{m}(\mathbf{B}) \mathbf{g}^{-1}(\mathbf{B}) \mathbf{n}(\mathbf{B})]^{-1} \end{pmatrix}$$

gives

$$\begin{aligned} \mathcal{G}^{-1}(-\mathbf{B}) &= \begin{pmatrix} \langle \alpha \alpha^T; -\mathbf{B} \rangle & \langle \alpha \beta^T; -\mathbf{B} \rangle \\ \langle \beta \alpha^T; -\mathbf{B} \rangle & \langle \beta \beta^T; -\mathbf{B} \rangle \end{pmatrix} \\ &= \begin{pmatrix} \langle \alpha \alpha^T; \mathbf{B} \rangle & -\langle \alpha \beta^T; \mathbf{B} \rangle \\ -\langle \beta \alpha^T; \mathbf{B} \rangle & \langle \beta \beta^T; \mathbf{B} \rangle \end{pmatrix} \end{aligned} \quad (10.209a)$$

Assuming time-reversal symmetry, the condition, (10.84) in §10.D.1, on the conditional probability generalizes, in the presence of a magnetic field, to [see (10.203a-b)]

$$P(\mathbf{q}^N, \mathbf{p}^N \mid \mathbf{q}'^N, \mathbf{p}'^N; \tau, \mathbf{B}) = P(\mathbf{q}'^N, -\mathbf{p}'^N \mid \mathbf{q}^N, -\mathbf{p}^N; \tau, -\mathbf{B}) \quad (10.210)$$

while (10.85) becomes

$$f(\alpha, \beta, \mathbf{B}) P(\alpha, \beta \mid \alpha', \beta'; \tau, \mathbf{B}) = f(\alpha', -\beta', -\mathbf{B}) P(\alpha', -\beta' \mid \alpha, -\beta; \tau, -\mathbf{B}) \quad (10.211)$$

Similarly, (10.75) generalizes to

$$\langle \alpha \alpha^T(\tau); \mathbf{B} \rangle = \langle \alpha(\tau) \alpha^T; -\mathbf{B} \rangle \quad (10.212)$$

$$\langle \alpha \beta^T(\tau); \mathbf{B} \rangle = -\langle \alpha(\tau) \beta^T; -\mathbf{B} \rangle \quad (10.213)$$

$$\langle \beta \alpha^T(\tau); \mathbf{B} \rangle = -\langle \beta(\tau) \alpha^T; -\mathbf{B} \rangle \quad (10.214)$$

$$\langle \beta \beta^T(\tau); \mathbf{B} \rangle = \langle \beta(\tau) \beta^T; -\mathbf{B} \rangle \quad (10.215)$$

or

$$\begin{pmatrix} \langle \alpha \alpha^T(\tau); \mathbf{B} \rangle & \langle \alpha \beta^T(\tau); \mathbf{B} \rangle \\ \langle \beta \alpha^T(\tau); \mathbf{B} \rangle & \langle \beta \beta^T(\tau); \mathbf{B} \rangle \end{pmatrix} = \begin{pmatrix} \langle \alpha(\tau) \alpha^T; -\mathbf{B} \rangle & -\langle \alpha \beta^T(\tau); -\mathbf{B} \rangle \\ -\langle \beta \alpha^T(\tau); -\mathbf{B} \rangle & \langle \beta \beta^T(\tau); -\mathbf{B} \rangle \end{pmatrix} \quad (10.215a)$$

Next, (10.87) needs to be upgraded to

$$\frac{d}{dt} \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix} = -\mathbf{M}(\mathbf{B}) \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix} \quad (10.216)$$

where

$$\mathbf{M}(\mathbf{B}) = \begin{pmatrix} \mathbf{M}^{\alpha\alpha}(\mathbf{B}) & \mathbf{M}^{\alpha\beta}(\mathbf{B}) \\ \mathbf{M}^{\beta\alpha}(\mathbf{B}) & \mathbf{M}^{\beta\beta}(\mathbf{B}) \end{pmatrix} \quad (10.217)$$

Finally, (10.96) generalizes to

$$\frac{d}{dt} \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix} = -\frac{1}{T} \mathbb{L}(\mathbf{B}) \begin{pmatrix} \mathbf{F} \\ \mathbf{G} \end{pmatrix} \quad (10.218)$$

with

$$\mathbb{L}(\mathbf{B}) = \begin{pmatrix} \mathbb{L}^{\alpha\alpha}(\mathbf{B}) & \mathbb{L}^{\alpha\beta}(\mathbf{B}) \\ \mathbb{L}^{\beta\alpha}(\mathbf{B}) & \mathbb{L}^{\beta\beta}(\mathbf{B}) \end{pmatrix} \quad (10.218a)$$

For short times, solution of (10.216) is

$$\begin{pmatrix} \alpha(\tau) \\ \beta(\tau) \end{pmatrix} = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} - \tau \mathbf{M}(\mathbf{B}) \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + O(\tau^2) \quad \begin{pmatrix} \alpha(0) \\ \beta(0) \end{pmatrix} = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

Hence,

$$\begin{pmatrix} \langle \alpha \alpha^T(\tau); \mathbf{B} \rangle & \langle \alpha \beta^T(\tau); \mathbf{B} \rangle \\ \langle \beta \alpha^T(\tau); \mathbf{B} \rangle & \langle \beta \beta^T(\tau); \mathbf{B} \rangle \end{pmatrix} = \int d\alpha \int d\beta f(\alpha, \beta, \mathbf{B}) \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \begin{pmatrix} \alpha^T(\tau) & \beta^T(\tau) \end{pmatrix}$$

$$\begin{aligned}
&= \int d\alpha \int d\beta f(\alpha, \beta, \mathbf{B}) \begin{pmatrix} \alpha \\ \beta \end{pmatrix} (\alpha^T, \beta^T) e^{-\mathbf{M}^T(\mathbf{B})\tau} \\
&= \left(\begin{array}{cc} \langle \alpha \alpha^T; \mathbf{B} \rangle & \langle \alpha \beta^T; \mathbf{B} \rangle \\ \langle \beta \alpha^T; \mathbf{B} \rangle & \langle \beta \beta^T; \mathbf{B} \rangle \end{array} \right) e^{-\mathbf{M}^T(\mathbf{B})\tau} \\
&= \mathcal{G}^{-1}(\mathbf{B}) e^{-\mathbf{M}^T(\mathbf{B})\tau}
\end{aligned}$$

Similarly,

$$\begin{aligned}
\left(\begin{array}{cc} \langle \alpha(\tau) \alpha^T; -\mathbf{B} \rangle & -\langle \alpha \beta^T(\tau); -\mathbf{B} \rangle \\ -\langle \beta \alpha^T(\tau); -\mathbf{B} \rangle & \langle \beta \beta^T(\tau); -\mathbf{B} \rangle \end{array} \right) &= \int d\alpha \int d\beta f(\alpha, \beta, -\mathbf{B}) \begin{pmatrix} \alpha(\tau) \\ -\beta(\tau) \end{pmatrix} (\alpha^T, -\beta^T) \\
&= \int d\alpha \int d\beta f(\alpha, \beta, -\mathbf{B}) e^{-\mathbf{M}(-\mathbf{B})\tau} \begin{pmatrix} \alpha \\ -\beta \end{pmatrix} (\alpha^T, -\beta^T) \\
&= e^{-\mathbf{M}(-\mathbf{B})\tau} \left(\begin{array}{cc} \langle \alpha \alpha^T; -\mathbf{B} \rangle & -\langle \alpha \beta^T; -\mathbf{B} \rangle \\ -\langle \beta \alpha^T; -\mathbf{B} \rangle & \langle \beta \beta^T; -\mathbf{B} \rangle \end{array} \right) \\
&= e^{-\mathbf{M}(-\mathbf{B})\tau} \left(\begin{array}{cc} \langle \alpha \alpha^T; \mathbf{B} \rangle & \langle \alpha \beta^T; \mathbf{B} \rangle \\ \langle \beta \alpha^T; \mathbf{B} \rangle & \langle \beta \beta^T; \mathbf{B} \rangle \end{array} \right) \\
&= e^{-\mathbf{M}(-\mathbf{B})\tau} \mathcal{G}^{-1}(\mathbf{B})
\end{aligned}$$

Note that in general

$$f(\alpha, \beta, -\mathbf{B}) \neq f(\alpha, \beta, \mathbf{B})$$

even though by time reversal invariant

$$f(\alpha, -\beta, -\mathbf{B}) = f(\alpha, \beta, \mathbf{B})$$

The time reversal invariant relations (10.215a) becomes

$$\mathcal{G}^{-1}(\mathbf{B}) e^{-\mathbf{M}^T(\mathbf{B})\tau} = e^{-\mathbf{M}(-\mathbf{B})\tau} \mathcal{G}^{-1}(\mathbf{B})$$

which, for short times, gives [c.f. (10.93)]

$$\mathcal{G}^{-1}(\mathbf{B}) \mathbf{M}^T(\mathbf{B}) = \mathbf{M}(-\mathbf{B}) \mathcal{G}^{-1}(\mathbf{B}) \quad (10.218b)$$

Setting [c.f. (10.94)]

$$\mathbb{L}(\mathbf{B}) = \mathbf{M}(\mathbf{B}) \mathcal{G}^{-1}(\mathbf{B}) \quad (10.218c)$$

turns (10.216) into

$$\frac{d}{dt} \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix} = -\mathbb{L}(\mathbf{B}) \mathcal{G}(\mathbf{B}) \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix} = -\frac{1}{T} \mathbb{L}(\mathbf{B}) \begin{pmatrix} F \\ G \end{pmatrix}$$

which is the generalization of the generalized Ohm's law (10.96).

Using (10.218c) to eliminate \mathbf{M} in (10.218b) gives

$$\mathbb{L}^T(\mathbf{B}) = \mathbb{L}(-\mathbf{B}) \mathcal{G}(-\mathbf{B}) \mathcal{G}^{-1}(\mathbf{B})$$

or

$$[\mathcal{G}(\mathbf{B}) \mathbb{L}(\mathbf{B})]^T = \mathbb{L}(-\mathbf{B}) \mathcal{G}(-\mathbf{B}) \quad (10.219)$$

which is the Onsager's relations in the presence of a magnetic field.

For $\mathbf{B} = 0$, (10.219) reduces to

$$\mathbb{L}^T(0) = \mathbb{L}(0)$$

in agreement with (10.95).

Also, if only one type of fluctuations (either α or β) is involved, then

$$\mathcal{G}(-\mathbf{B}) = \mathcal{G}(\mathbf{B})$$

and (10.219) reduces to a more pleasant form

$$\mathbb{L}^T(\mathbf{B}) = \mathbb{L}(-\mathbf{B}) \quad (10.219a)$$